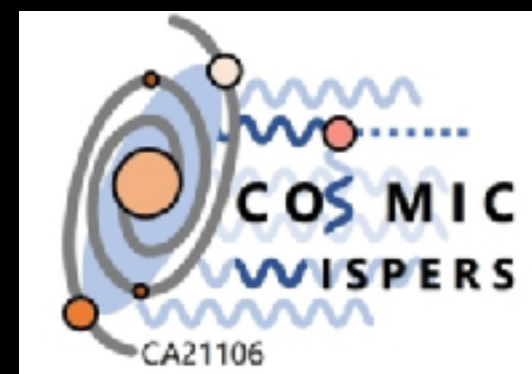


A Solar Halo of Ultralight Dark Matter

Joshua Eby
Oskar Klein Centre
Stockholm University

New Horizons for Psi
IST Lisbon
2024/07/04



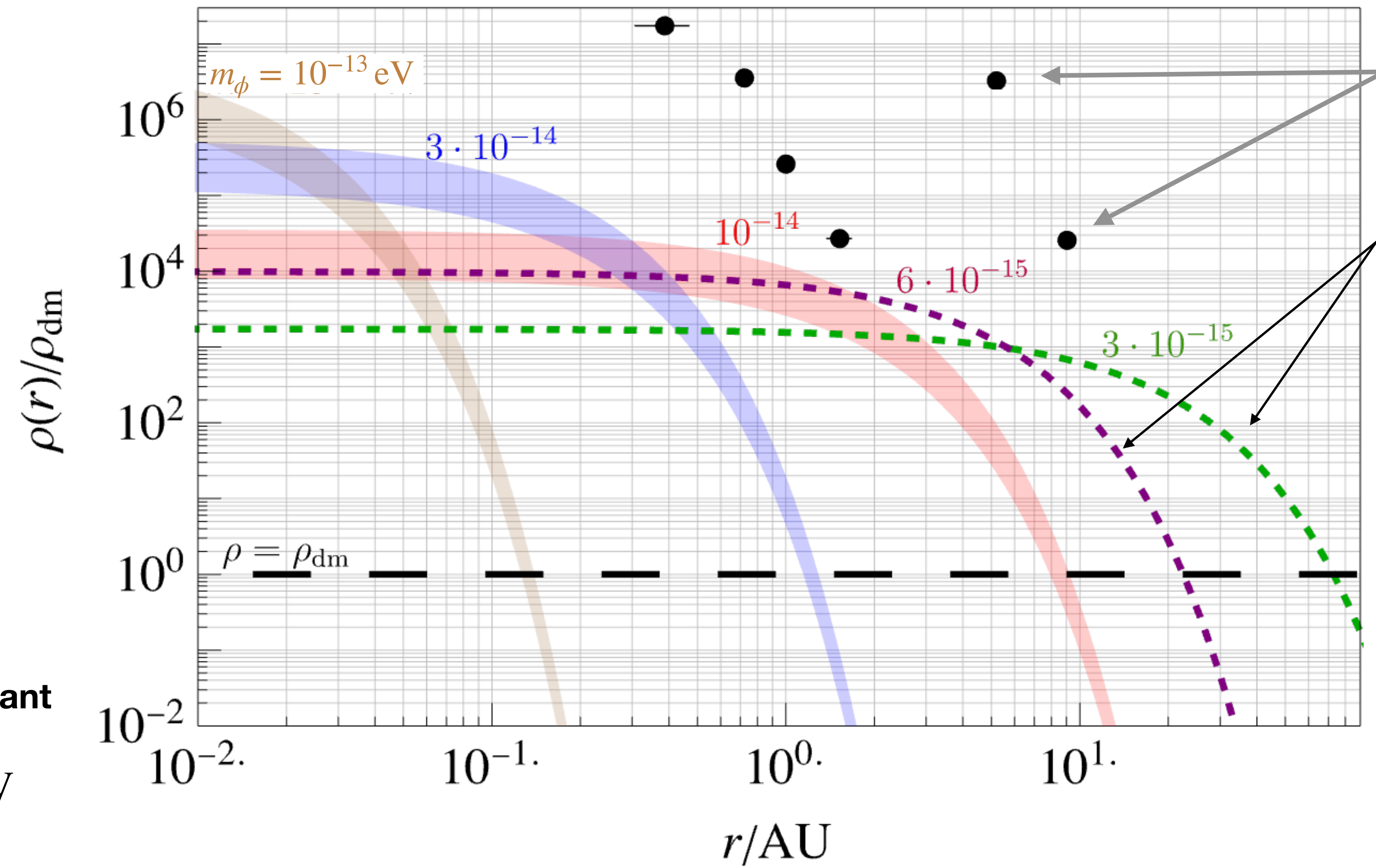
Stockholms
universitet

Conclusion: Solar Halo Density

Large modifications to the **very local density** from axions captured over **5Gyr (solar lifetime)**
(i.e. in our solar system)

Budker, **JE**, Gorghetto, Jiang, Perez (2306.12477)

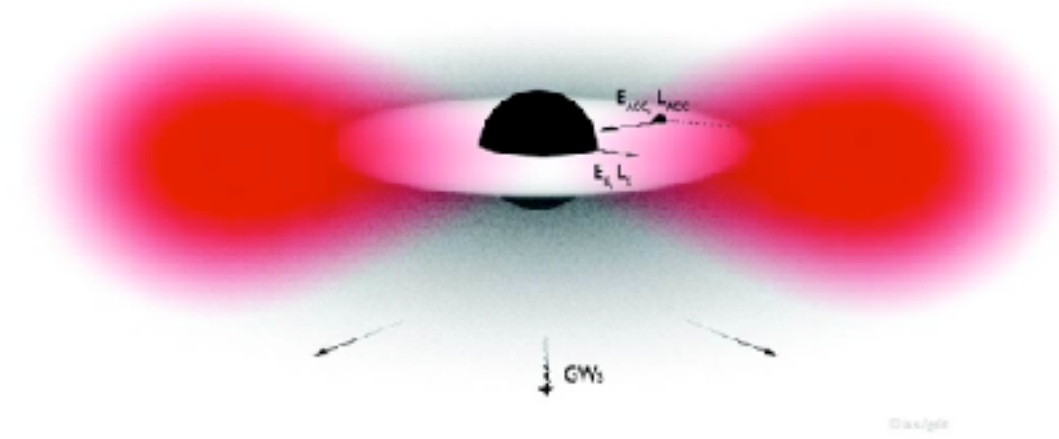
for axion decay constant
 $f_a \simeq 10^7 - 10^8 \text{ GeV}$



Direct constraints:
See H. Kim talk (Wednesday)
Exponentially-growing in low-velocity environment
 $v_{dm} \ll 200 \text{ km/s}$

assuming $\lambda < 0$;
density could be larger for $\lambda > 0$

Gravitational Atoms



States populated

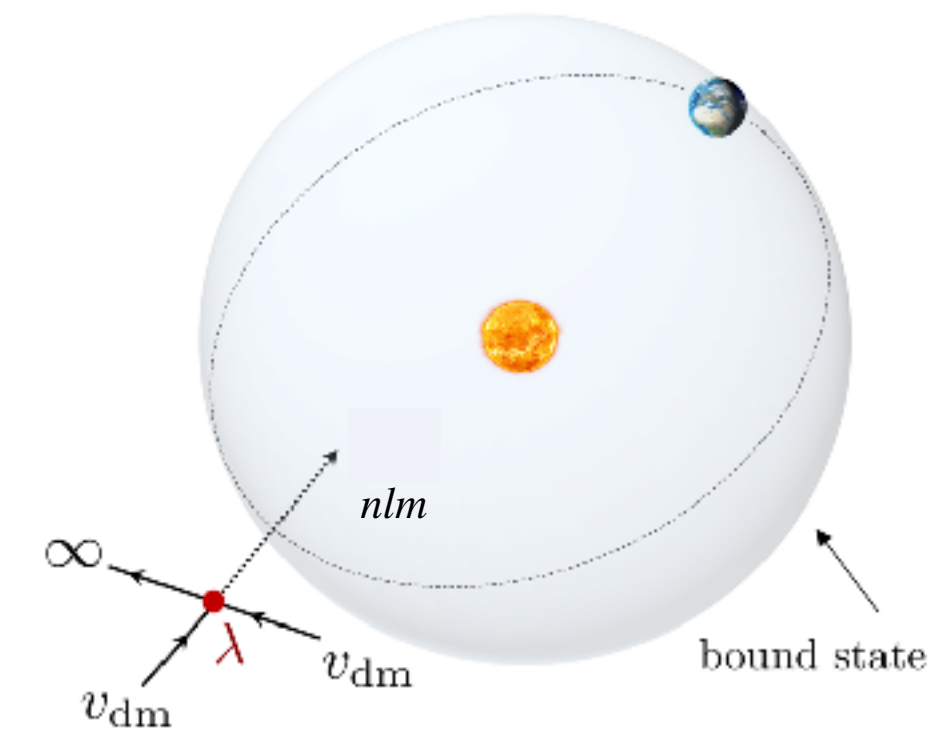
Superradiance

Many talks so far!

only $\ell \neq 0$
 primarily $\ell = m$ small

Ultralight dark matter (ULDM) capture

all $n\ell m$
 primarily $n\ell m = 100$



Efficient production on...

with $M \sim \frac{1}{Gm_\phi}$
 when $\omega_{n\ell m} \lesssim m\Omega_H$

Also stars (?) See
 J. McDonald talk (Wednesday)

with $M \gtrsim \frac{v_{dm}}{2\pi Gm_\phi}$

when $\lambda\phi^4$ is “strong enough”

Gravitational coupling

$$\alpha_g \equiv GMm_\phi \sim \mathcal{O}(1)$$

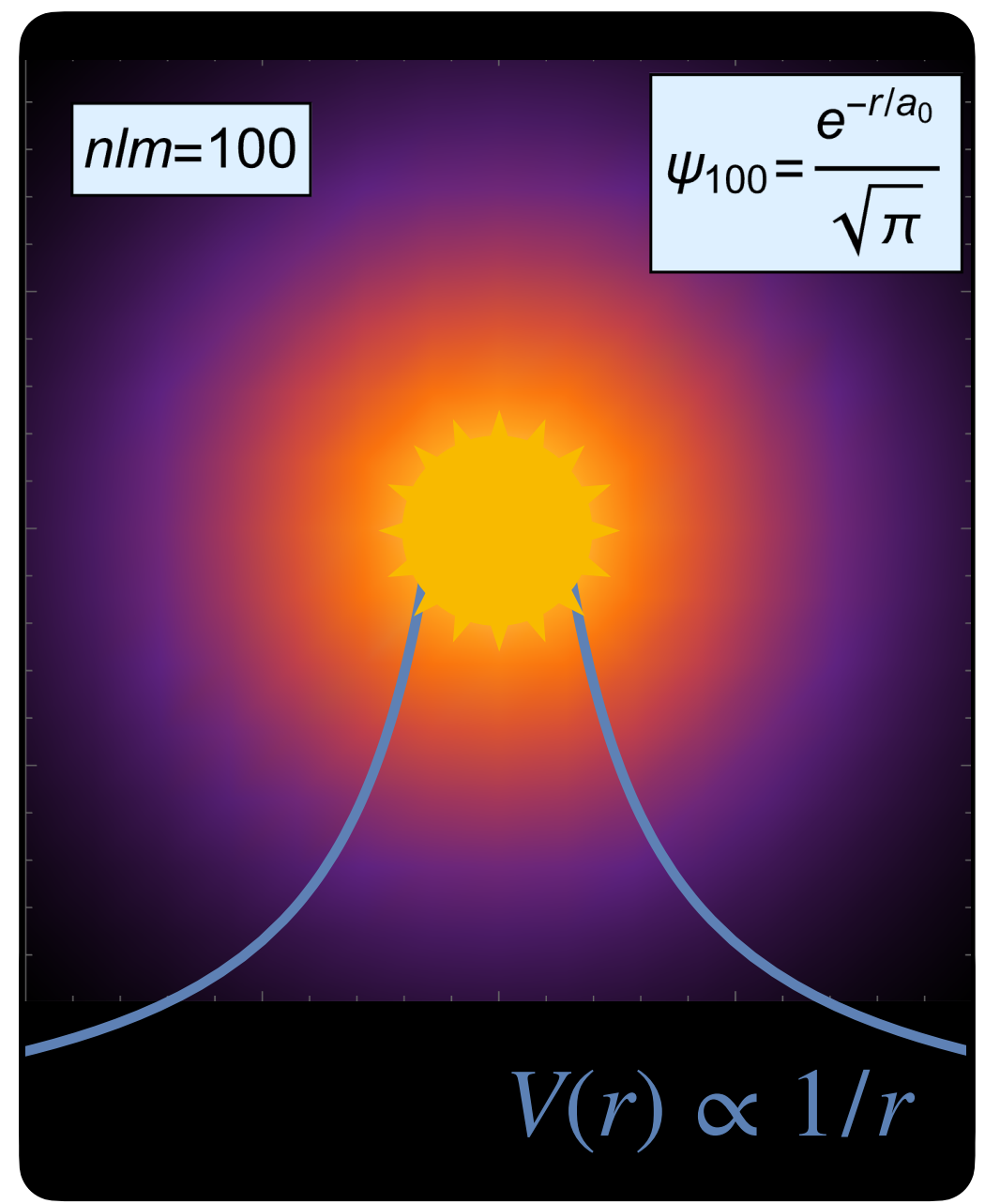
$$\alpha_g \gtrsim \frac{v_{dm}}{2\pi} \sim 10^{-4}$$

Energy density from...

BH mass / spin

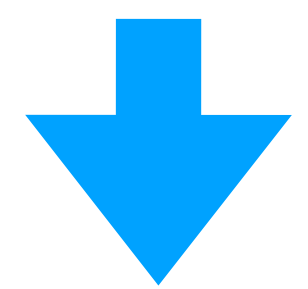
DM background

Gravitational Atoms

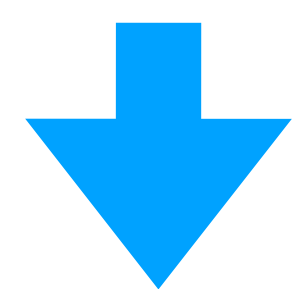


$$i \frac{\partial \psi}{\partial t} = \left[-\frac{\nabla^2}{2m_\phi} + \cancel{V_g(|\psi|^2)} + V_{g,ext}(r) + \cancel{\frac{\lambda}{8m_\phi^2} |\psi|^2} \right] \psi$$

low density



$$i \frac{\partial \psi}{\partial t} = \left[-\frac{\nabla^2}{2m_\phi} + \frac{\alpha_g}{r} \right] \psi$$



solutions

$$\psi = \psi_w + \psi_b$$

√[number density]

DM 'waves'
(scattering states)

$\sum_{nlm} \psi_{nlm}$
(bound states)

Hydrogen atom E.o.M. with

Fine structure constant

α

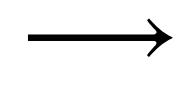


$$\alpha_g \equiv GMm_\phi$$

Gravitational coupling

Bohr radius

$$a_0 = \frac{1}{m_e \alpha}$$



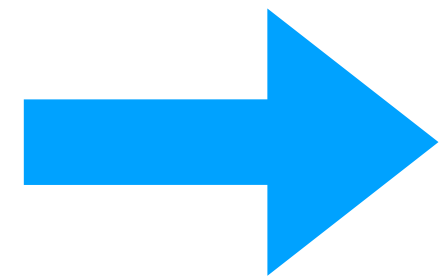
$$R_\star \equiv \frac{1}{m_\phi \alpha_g}$$

Gravitational 'Bohr radius'

DM Waves and Bound States

$$i \frac{\partial \psi}{\partial t} - \left[-\frac{\nabla^2}{2m_\phi} + \frac{\alpha_g}{r} \right] \psi = 0$$

Hydrogenic E.o.M.



$$\psi = \psi_w + \psi_b$$

Scattering states
(DM 'waves')

Solutions to
Hydrogenic E.o.M.

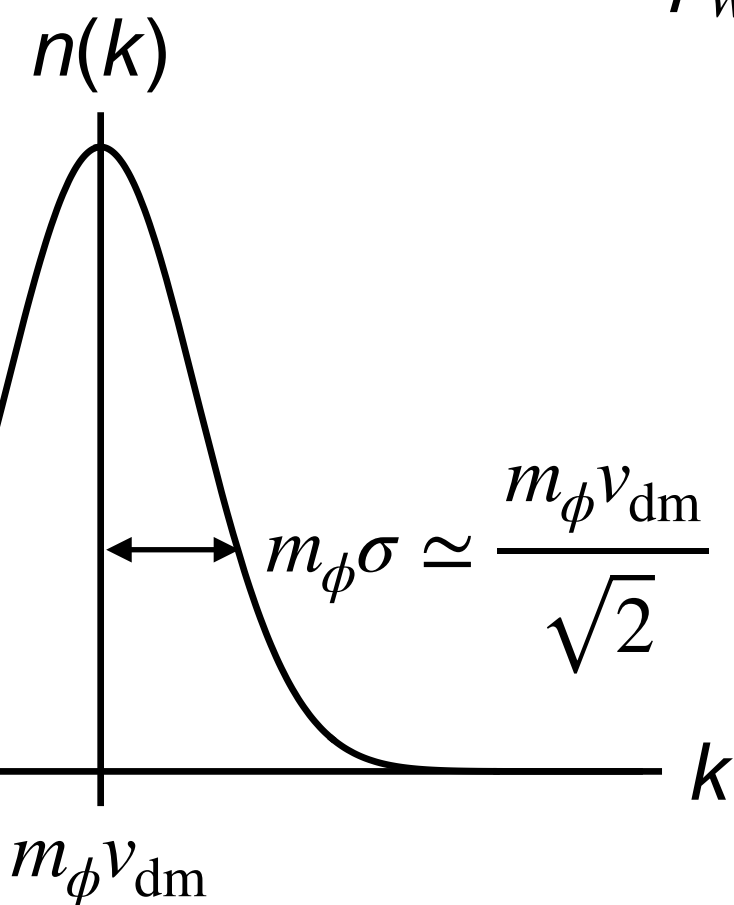
Bound states
of gravitational atom

$$\psi_b(t, x) = \sum_{nlm} e^{-i\omega_n t} \psi_{nlm}(x)$$

$$\propto L_{n-l-1}^{2l+1}(r) Y_l^m(\theta, \varphi)$$

energy
 $\omega_n < m_\phi$
quantum
numbers
 nlm

momentum
distribution



$$\psi_w(t, x) = \int \frac{d^3k}{(2\pi)^3} f(k) e^{-i\omega_k t} \psi_k(x)$$

Statistical sampling
of DM momenta in halo

energy
 $\omega_k > m_\phi$

$$\langle f^*(k) f(k') \rangle = (2\pi)^3 n(k) \delta^3(k - k')$$

momentum distribution:
Maxwell-Boltzmann

Solutions: Coulomb Scattering States

Simple solutions in the limits

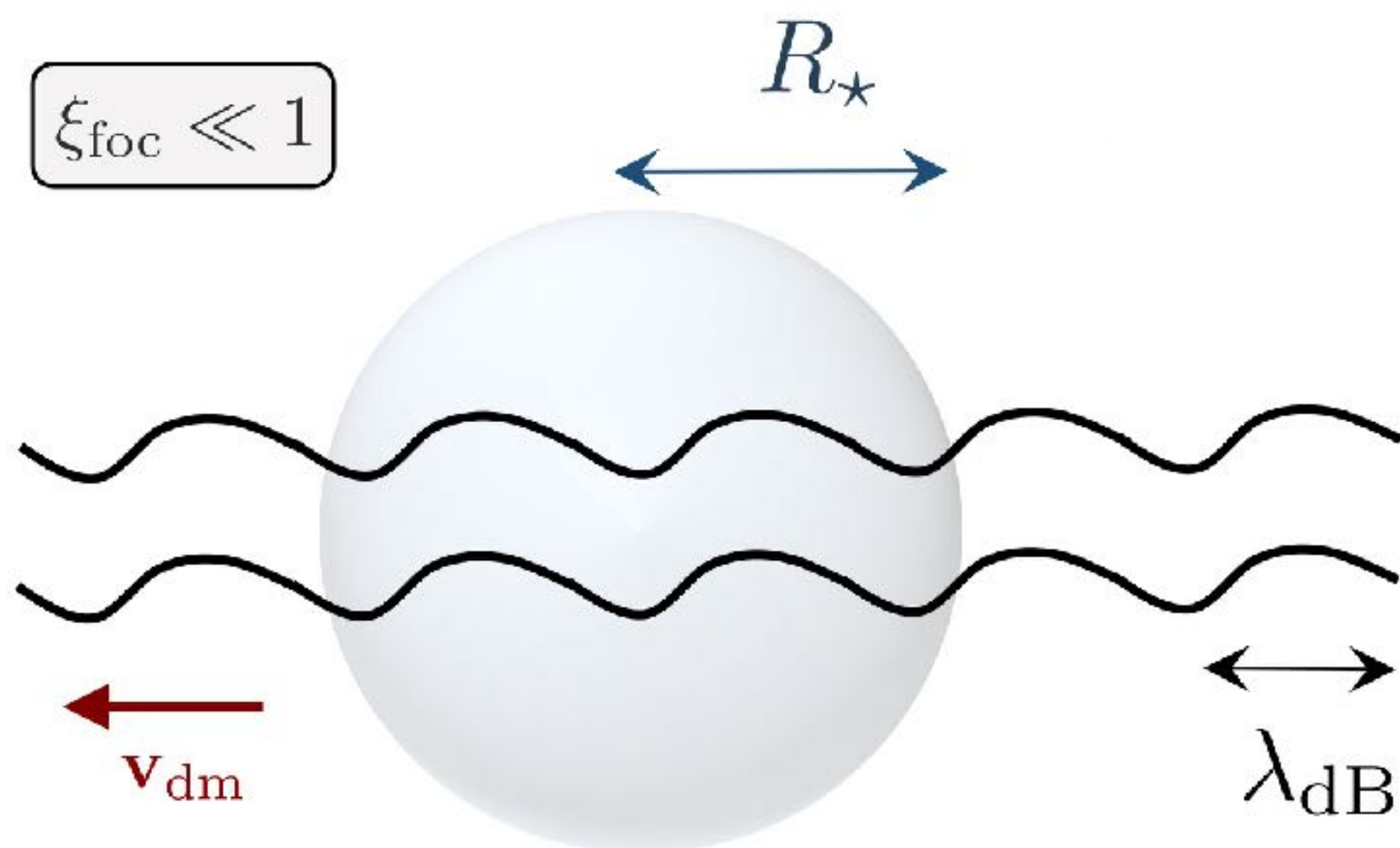
$$kR_\star \gg (2\pi)^{-1} \text{ and } kR_\star \ll (2\pi)^{-1}$$

We use $\xi_{\text{foc}} \equiv \frac{2\pi}{kR_\star} = \frac{\lambda_{\text{dB}}}{R_\star}$

DM Waves: Focusing

$$i \frac{\partial \psi_k}{\partial t} - \left[-\frac{\nabla^2}{2m_\phi} + \frac{\alpha_g}{r} \right] \psi_k = 0$$

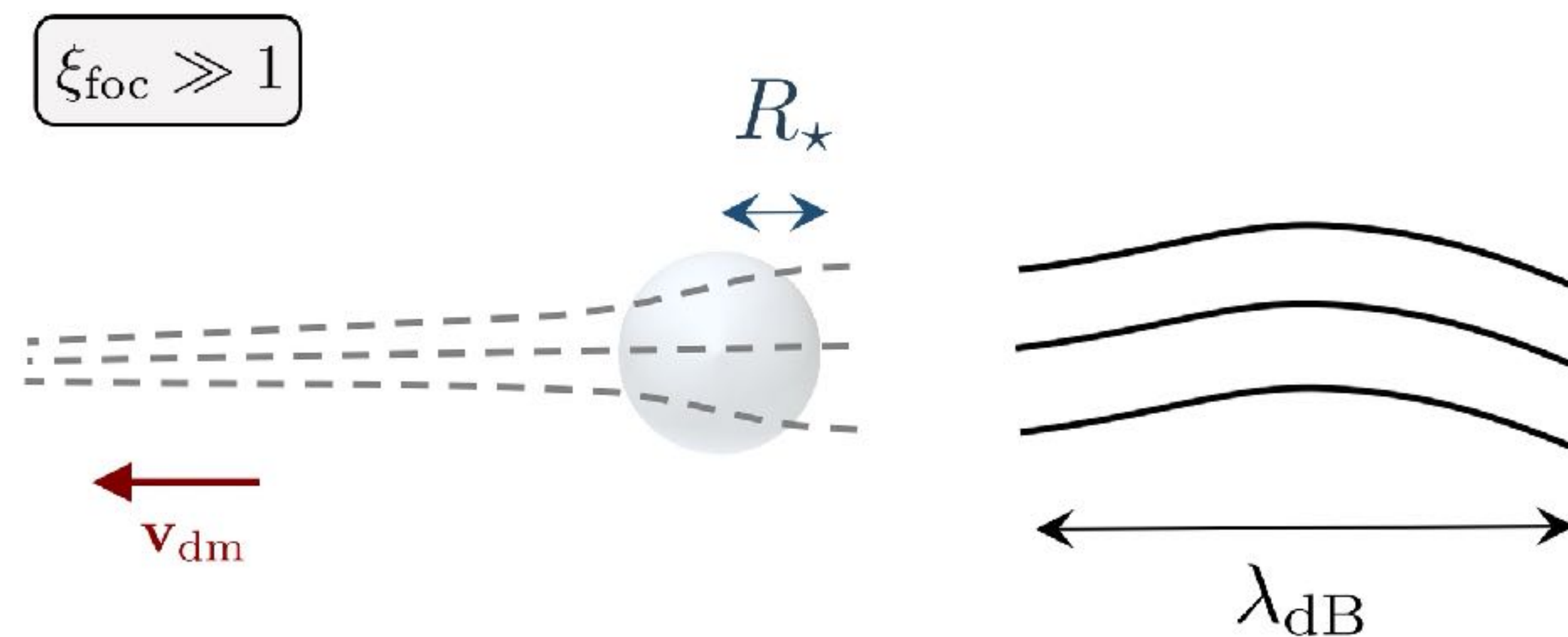
$$\xi_{\text{foc}} \equiv \frac{\lambda_{\text{dB}}}{R_\star} = \frac{2\pi\alpha_g}{v_{\text{dm}}} \quad (\text{dictates effectiveness of 'gravitational focusing'})$$



$$\psi_k \longrightarrow e^{ik \cdot x},$$

$$\text{typical energy } \omega_k \gg \left| \omega_n \right|$$

$\frac{m_\phi v_{\text{dm}}^2}{2} \quad \frac{m_\phi \alpha_g^2}{2}$

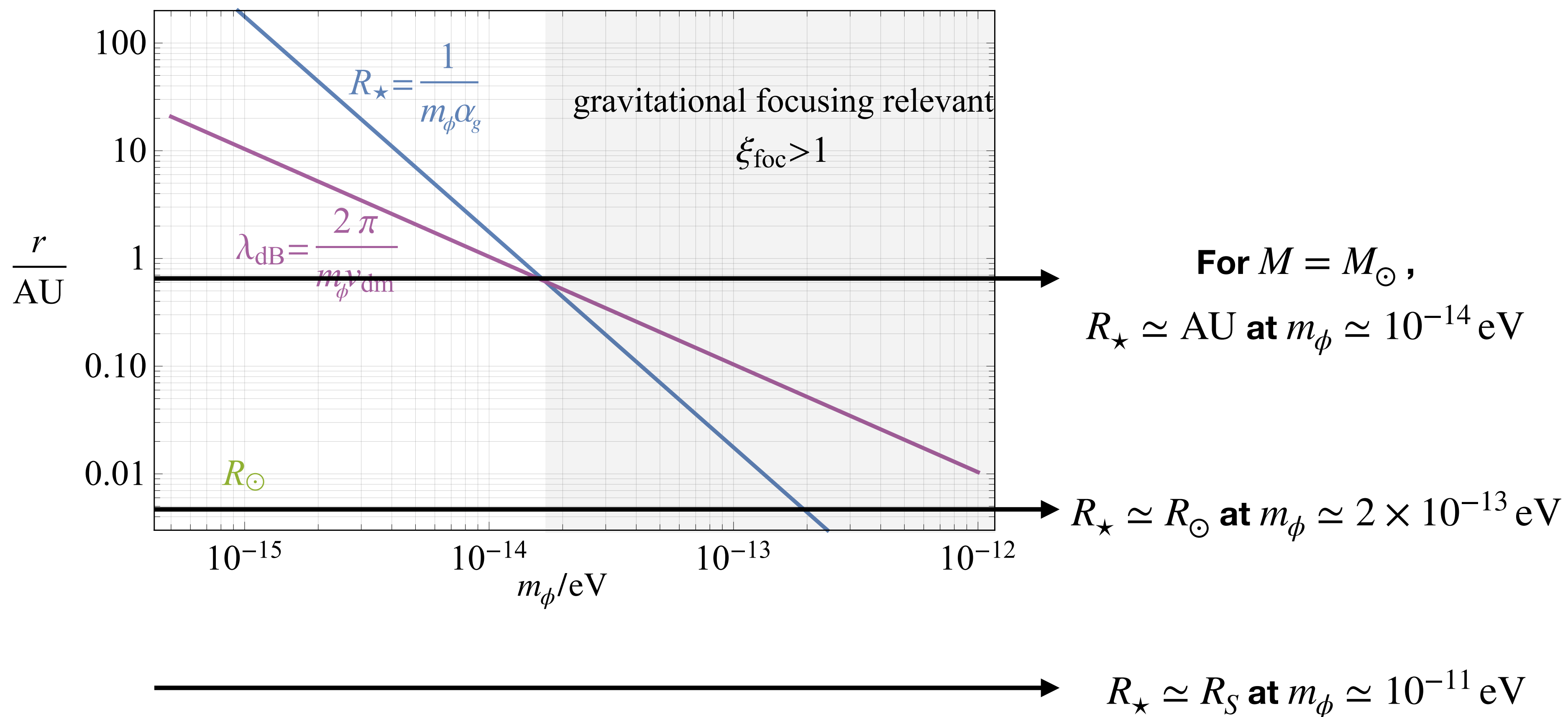


ψ_k focused onto region of size $\simeq R_\star$,
typical energy $\omega_k \ll |\omega_n|$

DM Waves in the Solar System

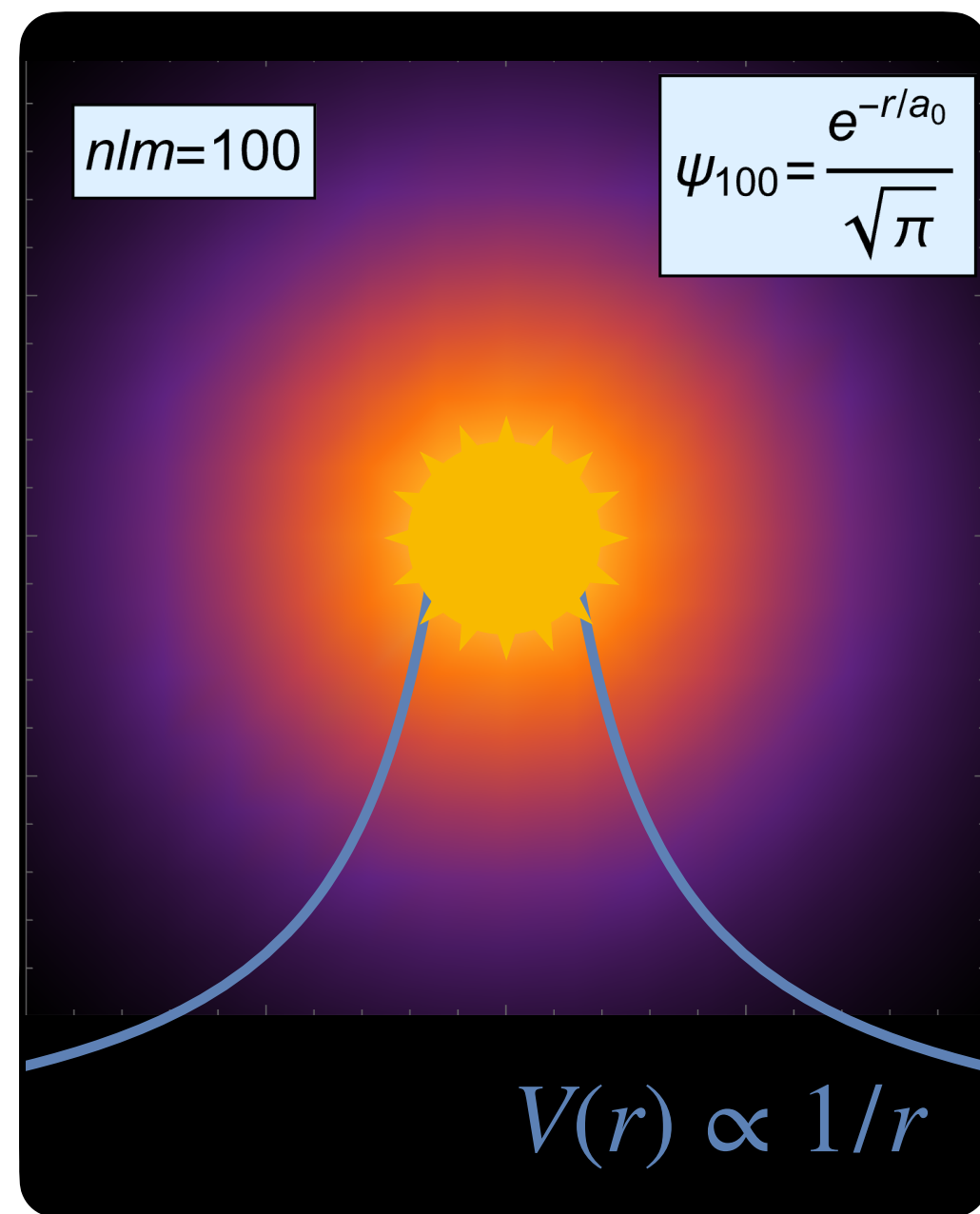
$$i \frac{\partial \psi_k}{\partial t} - \left[-\frac{\nabla^2}{2m_\phi} + \frac{\alpha_g}{r} \right] \psi_k = 0$$

$$\xi_{\text{foc}} \equiv \frac{\lambda_{\text{dB}}}{R_\star} = \frac{2\pi\alpha_g}{v_{\text{dm}}} \quad (\text{dictates effectiveness of 'gravitational focusing'})$$



A Solar Halo

$$i \frac{\partial \psi}{\partial t} - \left[-\frac{\nabla^2}{2m_\phi} + \frac{\alpha_g}{r} \right] \psi = 0$$



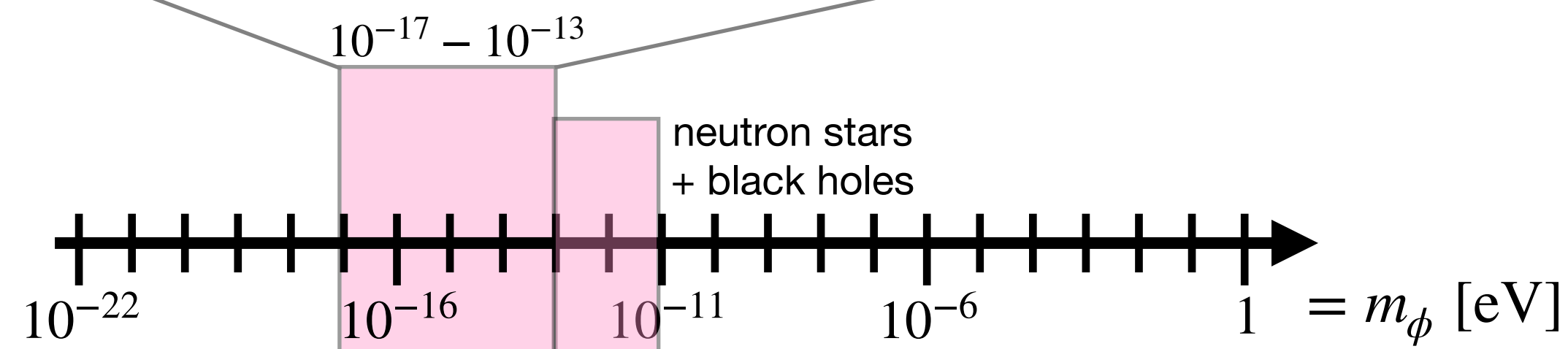
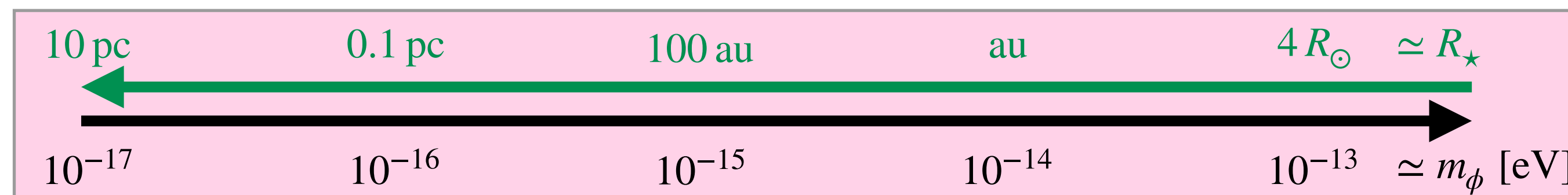
$$\alpha_g \equiv GMm_\phi \simeq 10^{-4} \left(\frac{m_\phi}{10^{-14} \text{ eV}} \right) \left(\frac{M}{M_\odot} \right)$$

$$\xi_{\text{foc}} \equiv \frac{2\pi\alpha_g}{v_{\text{dm}}} \sim \left(\frac{m_\phi}{10^{-14} \text{ eV}} \right) \left(\frac{M}{M_\odot} \right) \left(\frac{240 \text{ km/sec}}{v_{\text{dm}}} \right)$$

$$R_\star \equiv \frac{1}{m_\phi \alpha_g} \simeq 1 \text{ au} \left(\frac{10^{-14} \text{ eV}}{m_\phi} \right) \left(\frac{M_\odot}{M} \right)$$

distance between stars

radius of Sun



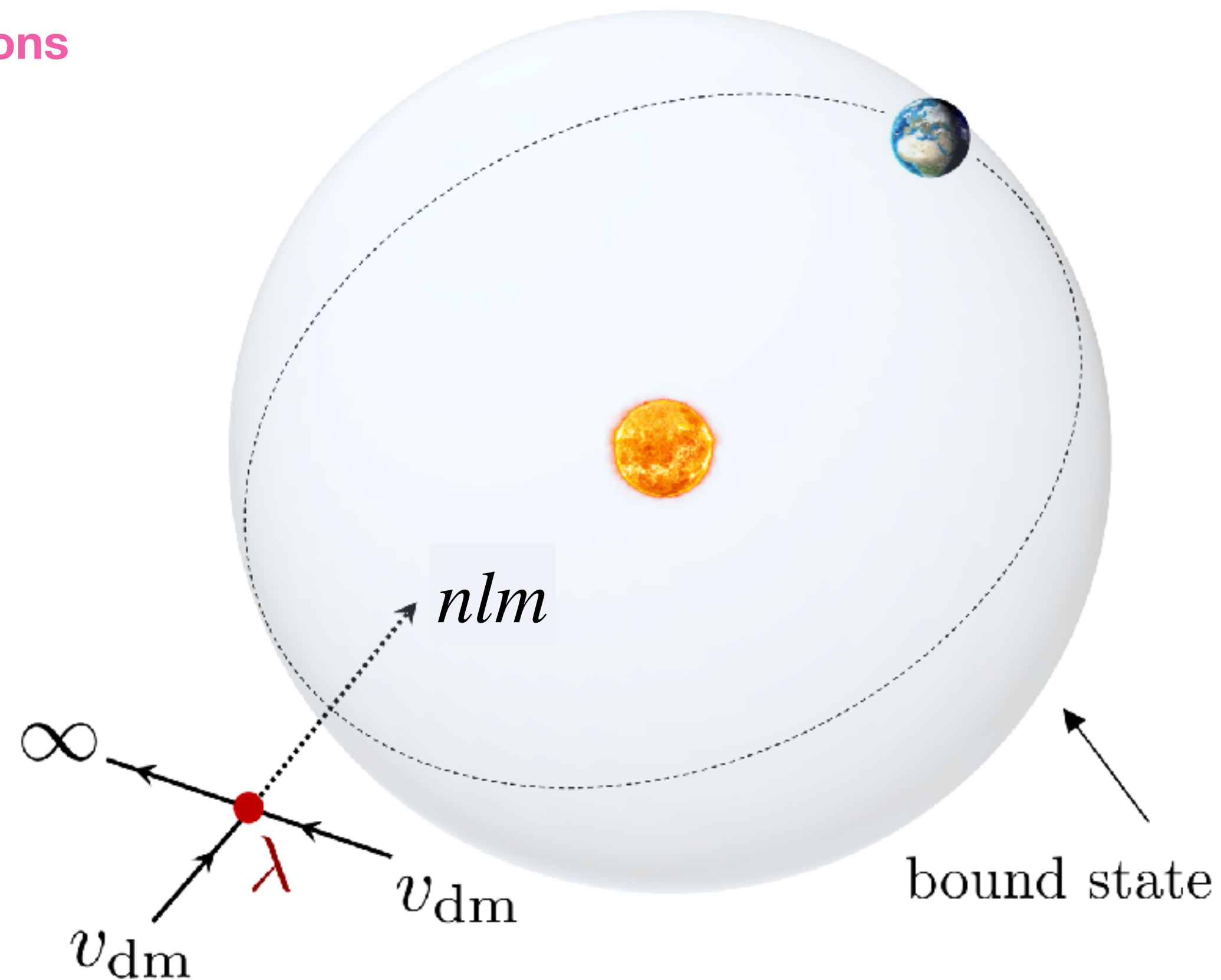
Budker, **JE**, Gorghetto,
Jiang, Perez (2306.12477)

Self-Interactions

$$i \frac{\partial \psi}{\partial t} - \left[-\frac{\nabla^2}{2m_\phi} + \frac{\alpha_g}{r} \right] \psi = \frac{\lambda}{8m_\phi^2} |\psi|^2 \psi$$

Self-interactions can move particles from scattering states to bound states (and vice versa)

Hydrogenic E.o.M. + self-interactions



Mass Growth

Initially: $N_{nlm}(t=0) = 0$

$$\frac{dM_\star}{dt} = m_\phi \sum_{nlm} \frac{dN_{nlm}}{dt} = \sum_{nlm} \frac{\pi\lambda^2}{16m_\phi^3} \int \frac{d^3k_1}{(2\pi)^3} \frac{d^3k_2}{(2\pi)^3} \frac{d^3k_3}{(2\pi)^3} \delta(\Delta\omega) |\mathcal{M}_{nlm}|^2 \mathcal{M}_{nlm} \equiv \int \psi_{k_1} \psi_{k_2} \psi_{k_3}^* \psi_{nlm}^* d^3x$$

phase space + energy conservation
 $\Delta\omega \equiv \omega_{k_1} + \omega_{k_2} - \omega_{k_3} - \omega_n$

matrix element

Boltzmann factors for the scattering rate

$$\times \left[n(k_1)n(k_2)(n(k_3)+1)(N_{nlm}+1) - n(k_3)N_{nlm}(n(k_1)+1)(n(k_2)+1) \right] + \mathcal{O}(N_{nlm}^2)$$

Excited states

$$\simeq \sum_{nlm} \frac{\pi\lambda^2}{16m_\phi^3} \int \frac{d^3k_1}{(2\pi)^3} \frac{d^3k_2}{(2\pi)^3} \frac{d^3k_3}{(2\pi)^3} \delta(\Delta\omega) |\mathcal{M}_{nlm}|^2 \left[n(k_1)n(k_2)n(k_3) + N_{nlm} (n(k_1)n(k_2) - 2n(k_2)n(k_3)) \right]$$

simplify:

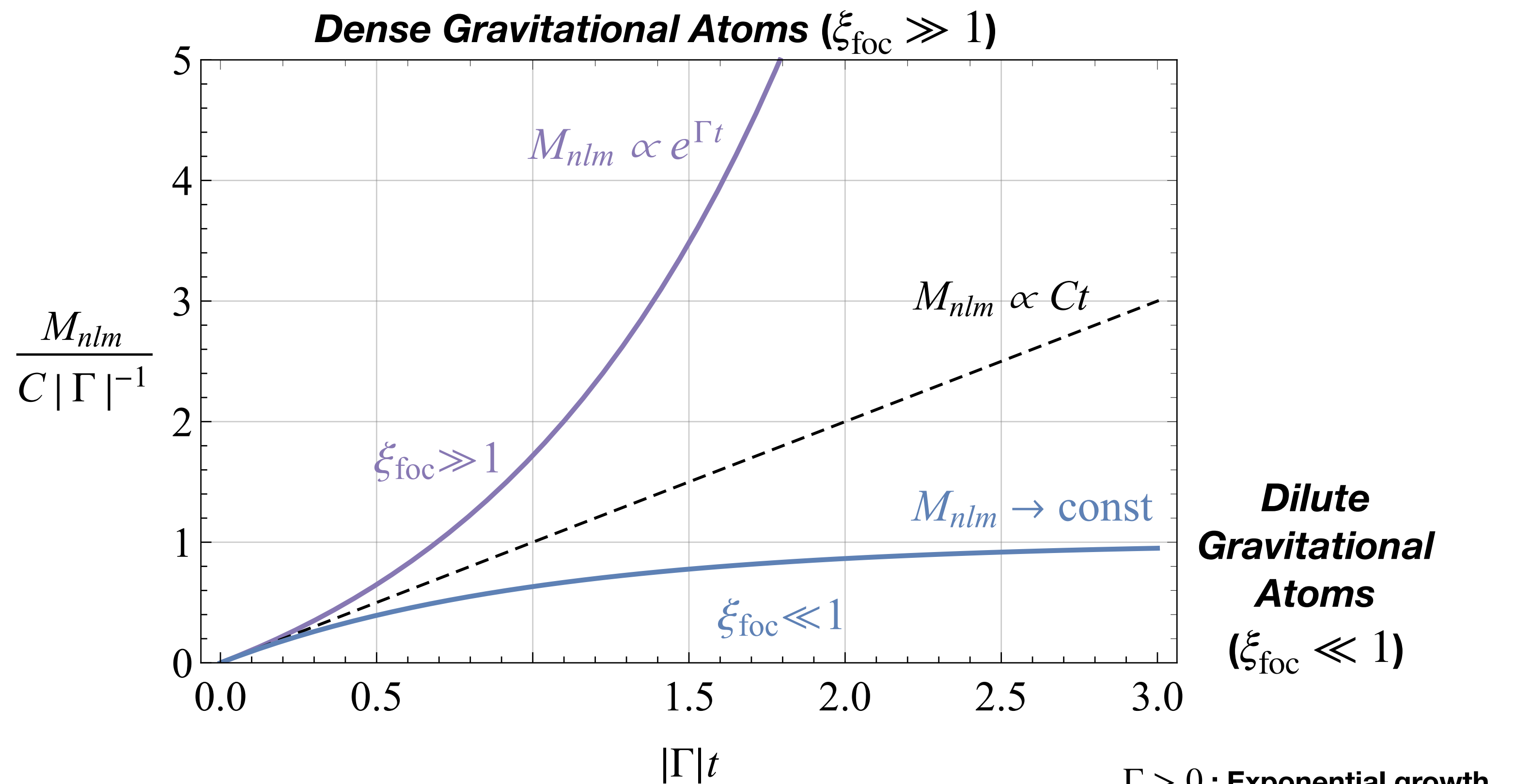
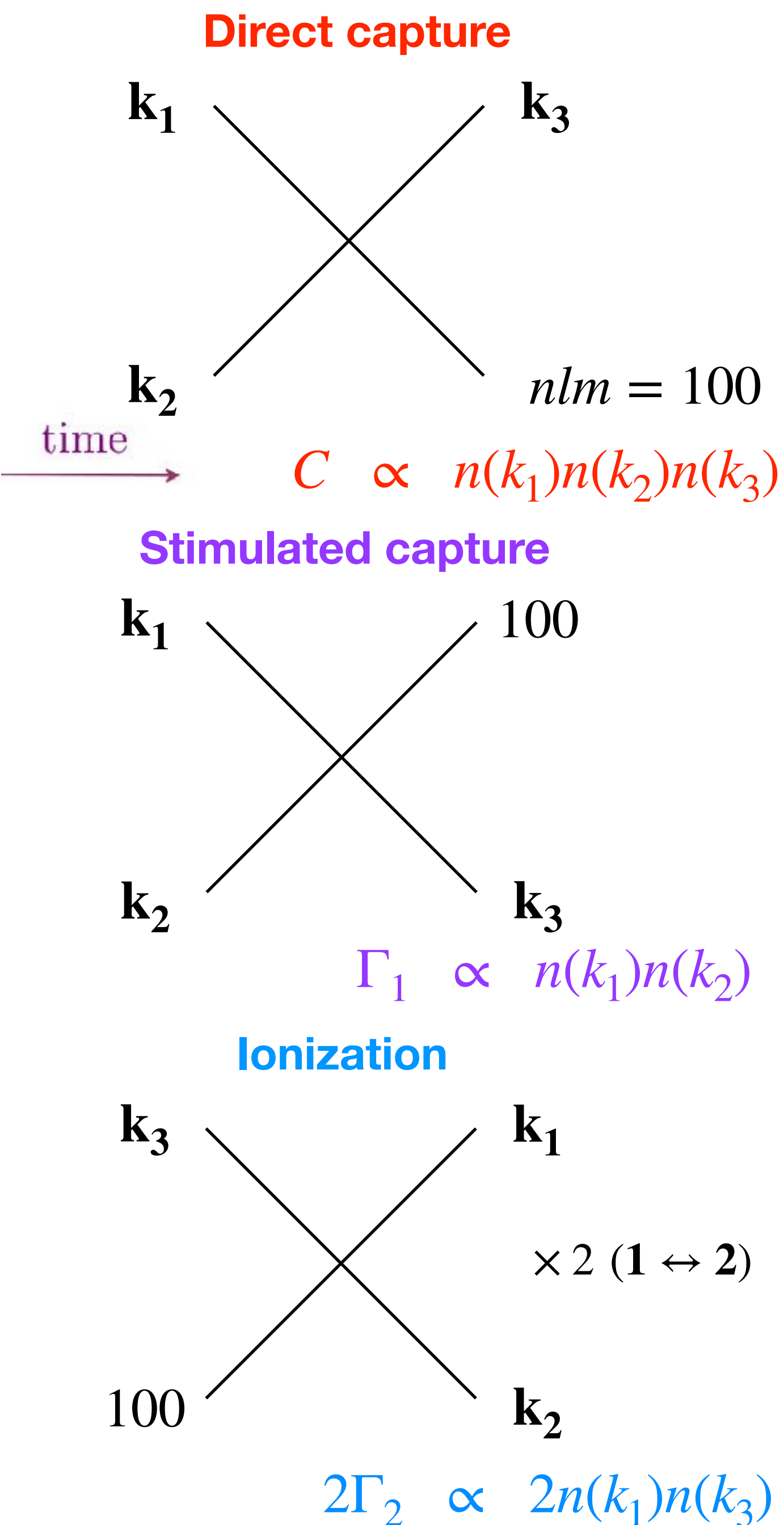
$$nlm \rightarrow 100$$

$$N_{nlm} \rightarrow N_{100} = M_\star / m_\phi$$

$$\frac{dM_\star}{dt} \simeq C + (\Gamma_1 - 2\Gamma_2) M_\star$$

Budker, JE, Gorghetto, Jiang, Perez (2306.12477)

Time Evolution



$\Gamma > 0$: Exponential growth
 $\Gamma < 0$: Saturation
 determines late-time behavior

$$\frac{dM_{\star}}{dt} \simeq C + (\Gamma_1 - 2\Gamma_2) M_{\star}$$

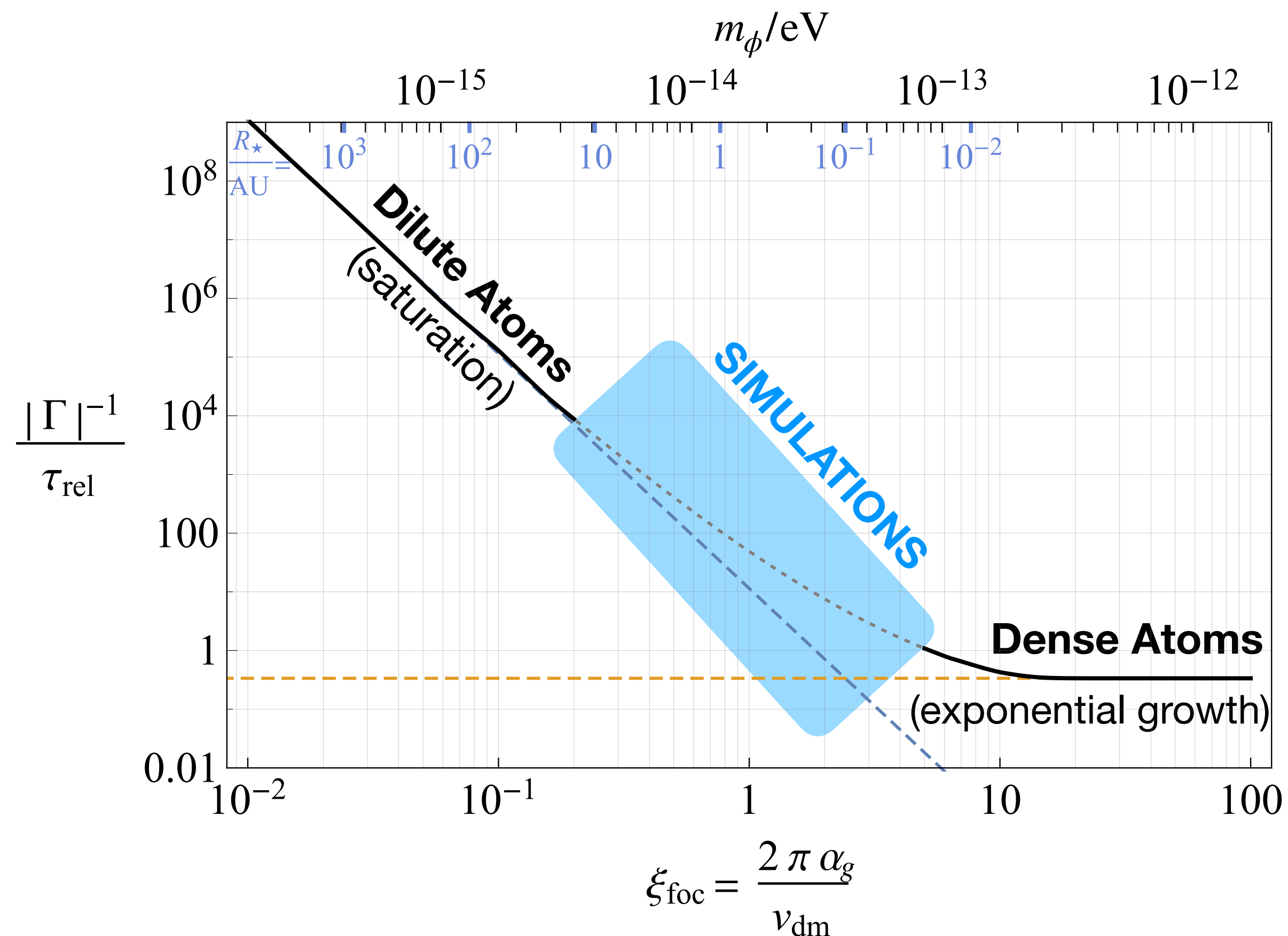
Budker, **JE**, Gorghetto,
Jiang, Perez (2306.12477)

Relaxation Timescale

$$\tau_{\text{rel}} \equiv \frac{64m_\phi^7 v_{\text{dm}}^2}{\lambda^2 \rho_{\text{dm}}^2} = \frac{64m_\phi^3 f_a^4 v_{\text{dm}}^2}{\rho_{\text{dm}}^2} \simeq 9 \text{ Gyr} \left(\frac{f_a}{10^8 \text{ GeV}} \right)^4 \left(\frac{m_\phi}{10^{-14} \text{ eV}} \right)^3 \left(\frac{0.4 \text{ GeV/cm}^3}{\rho_{\text{dm}}} \right)^2 \left(\frac{v_{\text{dm}}}{240 \text{ km/sec}} \right)^2$$

For axions,

$$\lambda = - \frac{m_\phi^2}{f_a^2}$$



Low decay constant f_a required
for astrophysical timescale

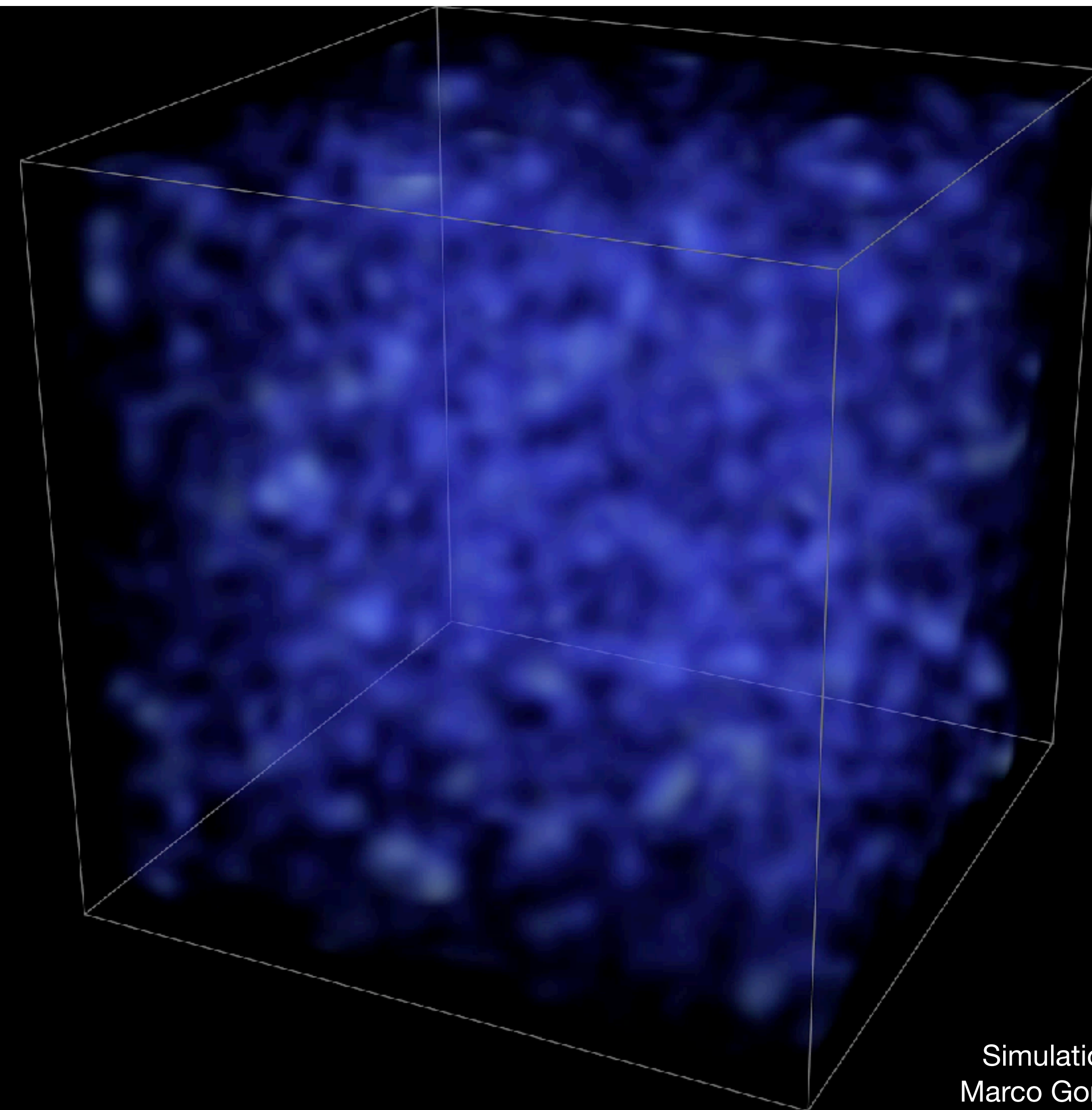
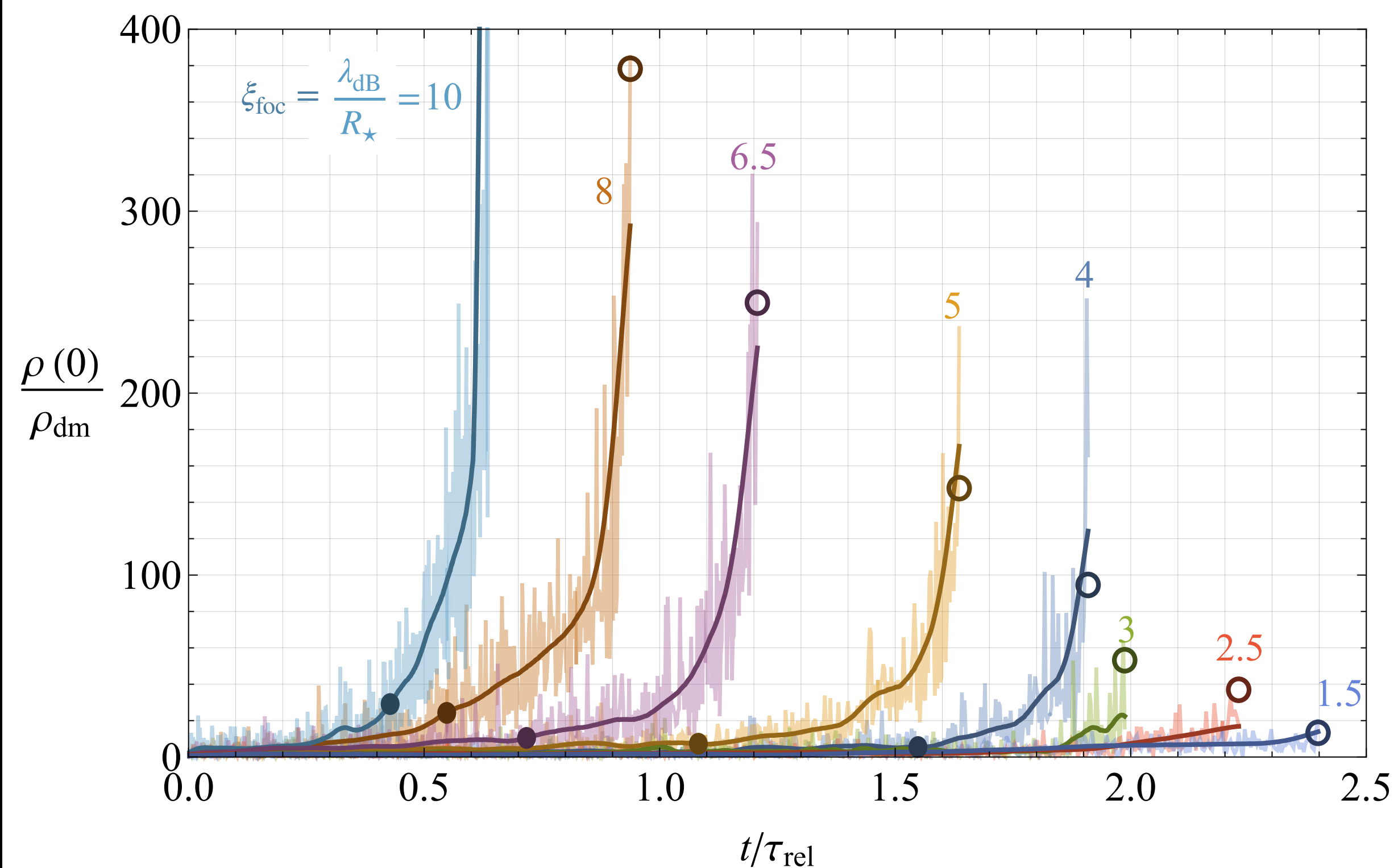
*Kinetic misalignment,
Non-periodic potentials,
etc.*

See G. Servant lectures (Monday)

Budker, **JE**, Gorghetto,
Jiang, Perez (2306.12477)

Numerical Simulations

$$i \frac{\partial \psi}{\partial t} = \left[-\frac{\nabla^2}{2m_\phi} + \frac{\alpha_g}{r} + \frac{\lambda}{8m_\phi^2} |\psi|^2 \right] \psi$$

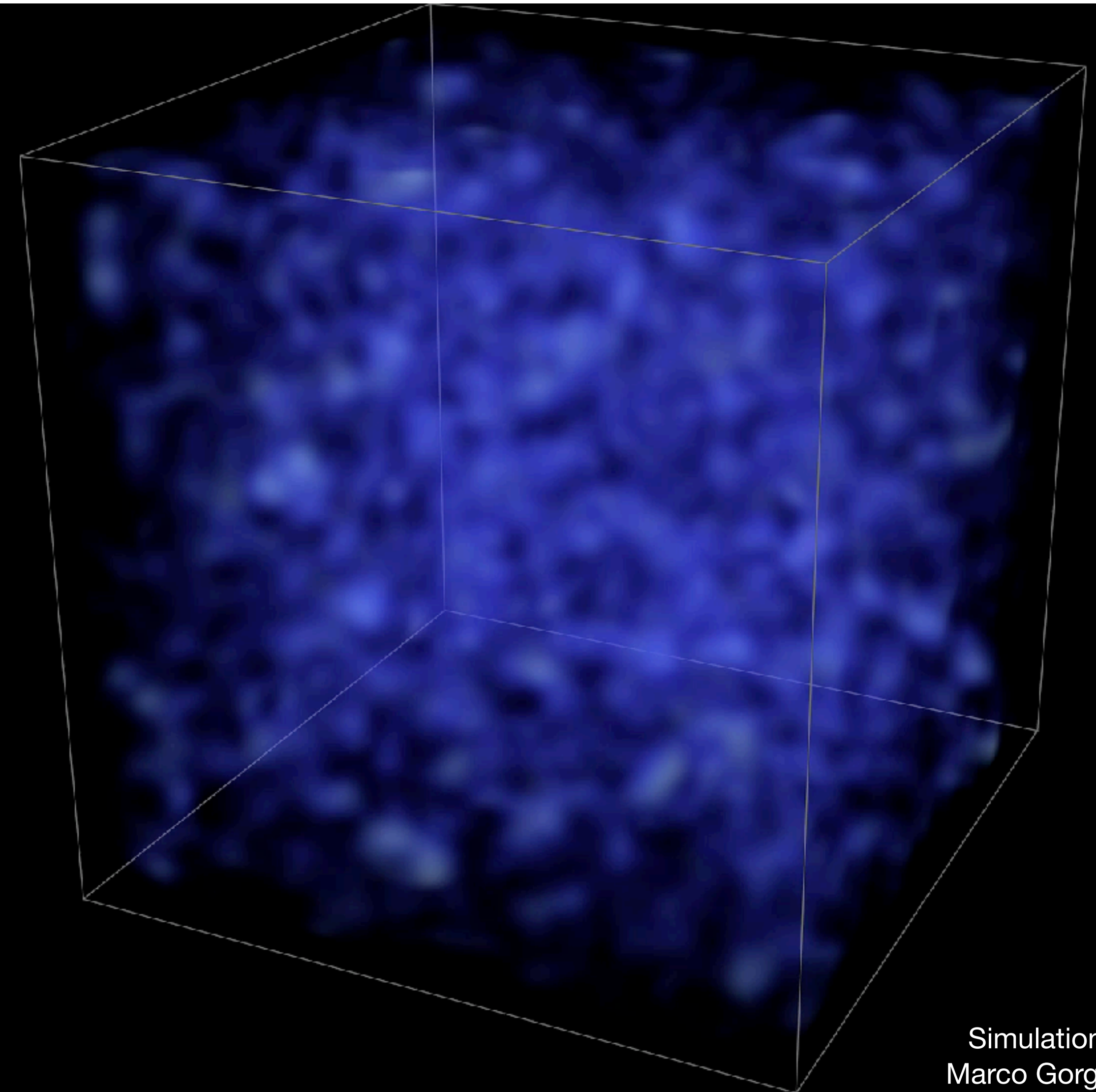
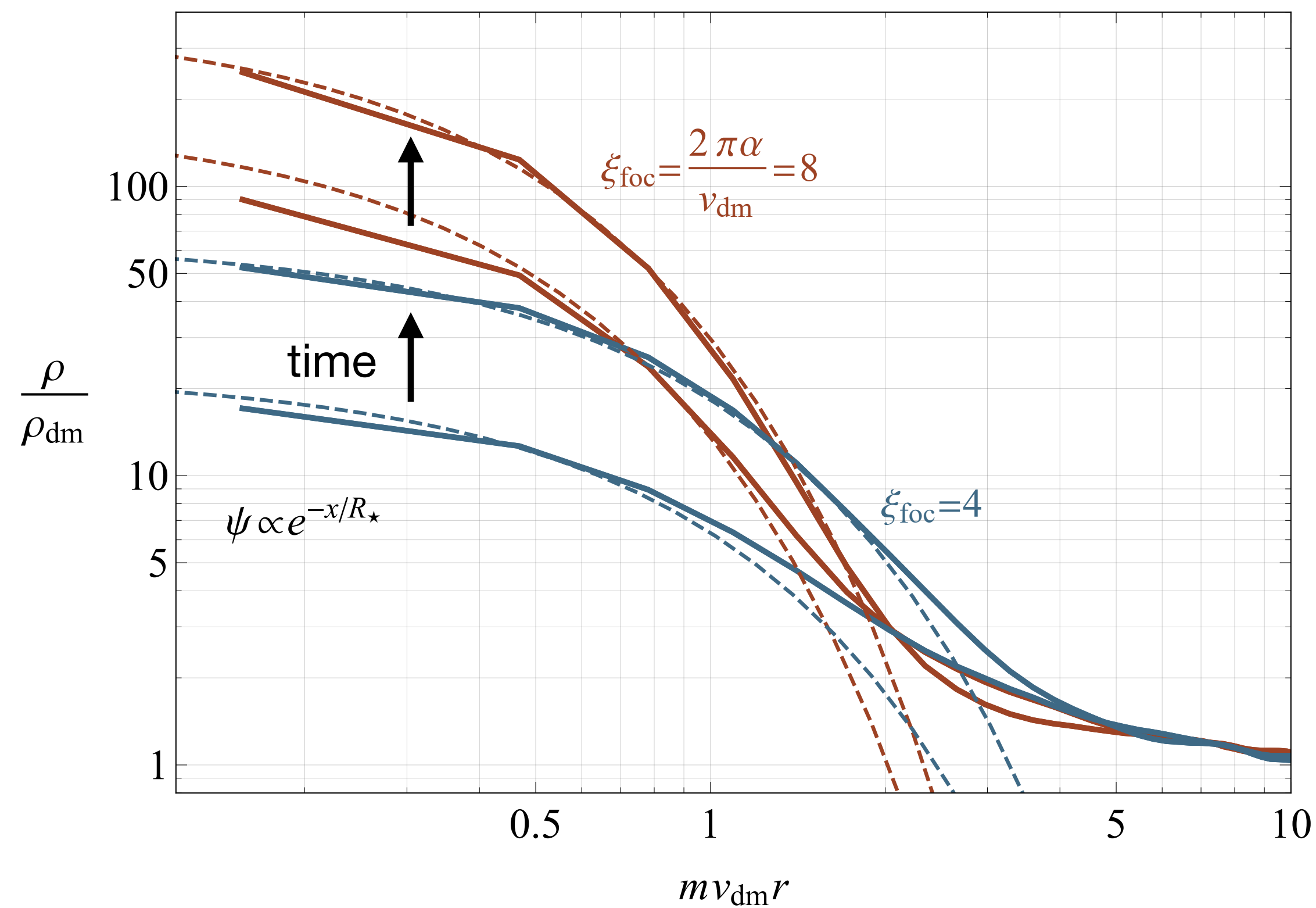


Simulation by
Marco Gorghetto

Budker, **JE**, Gorghetto,
Jiang, Perez (2306.12477)

Numerical Simulations

$$i \frac{\partial \psi}{\partial t} = \left[-\frac{\nabla^2}{2m_\phi} + \frac{\alpha_g}{r} + \frac{\lambda}{8m_\phi^2} |\psi|^2 \right] \psi$$

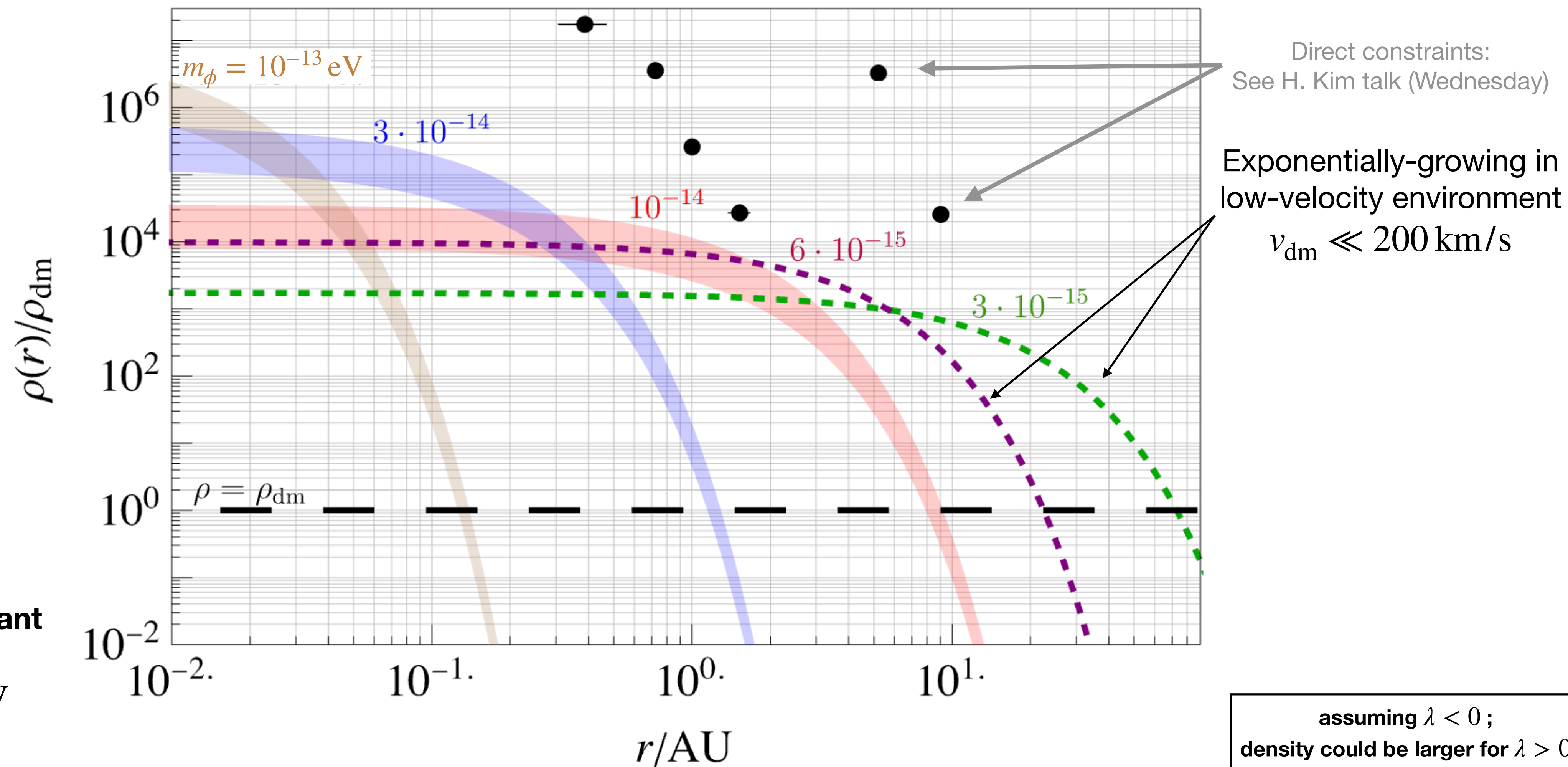


Simulation by
Marco Gorghetto

Budker, **JE**, Gorghetto, Jiang, Perez (2306.12477)

Solar Halo Density

Large modifications to the **very local density** from axions captured over **5Gyr (solar lifetime)**
 (i.e. in our solar system)



for axion decay constant

$$f_a \simeq 10^7 - 10^8 \text{ GeV}$$

Budker, JE, Gorghetto, Jiang, Perez (2306.12477)

Fate of Gravitational Atoms

$$i \frac{\partial \psi}{\partial t} = \left[-\frac{\nabla^2}{2m_\phi} + \frac{\alpha_g}{r} + \frac{\lambda}{8m_\phi^2} |\psi|^2 \right] \psi$$

Density grows \iff Self-interaction term grows

until $\frac{\alpha_g}{R_\star} \simeq \frac{|\lambda|}{8m_\phi^3} \rho$

Attractive self-interaction
 $\lambda < 0$

Collapse! when

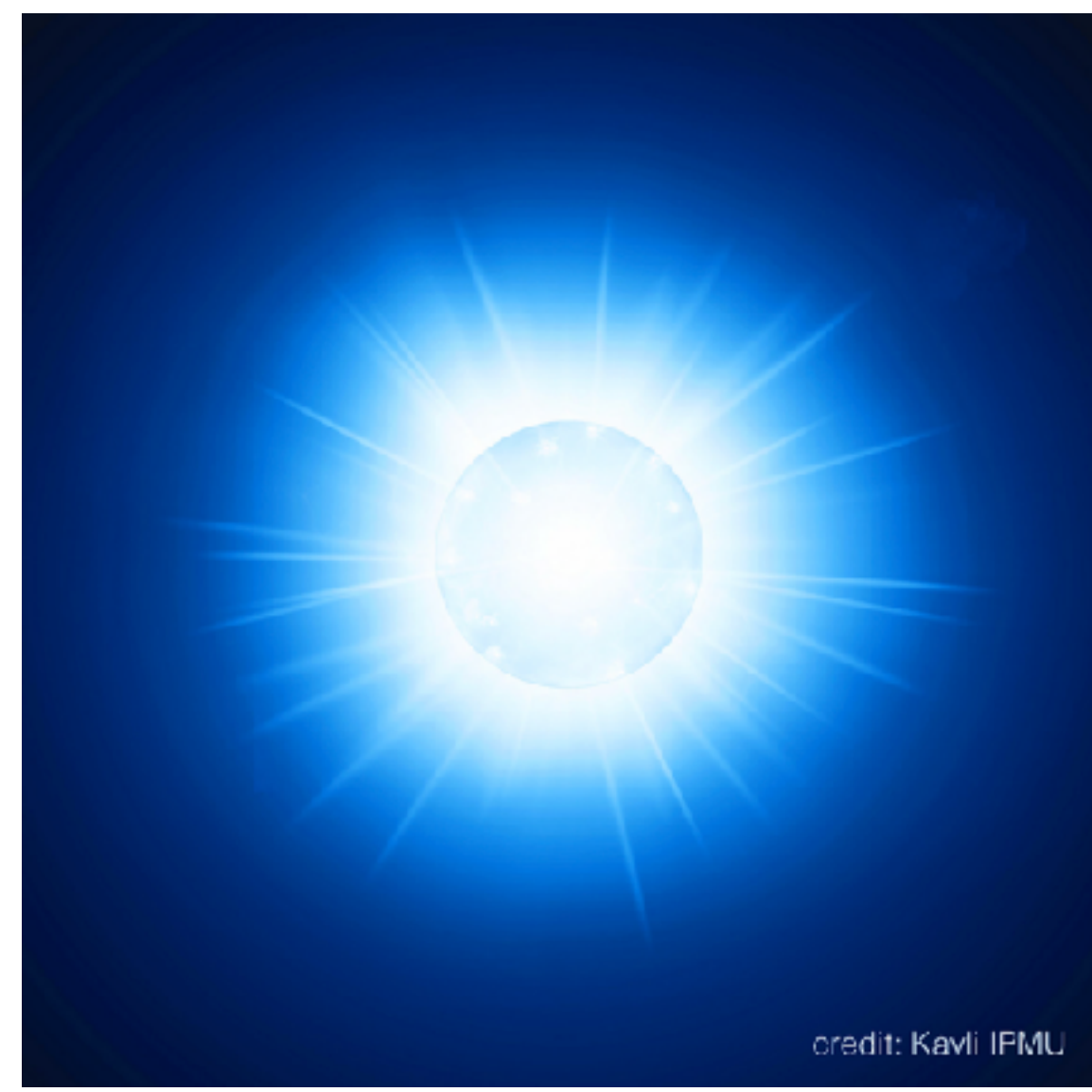
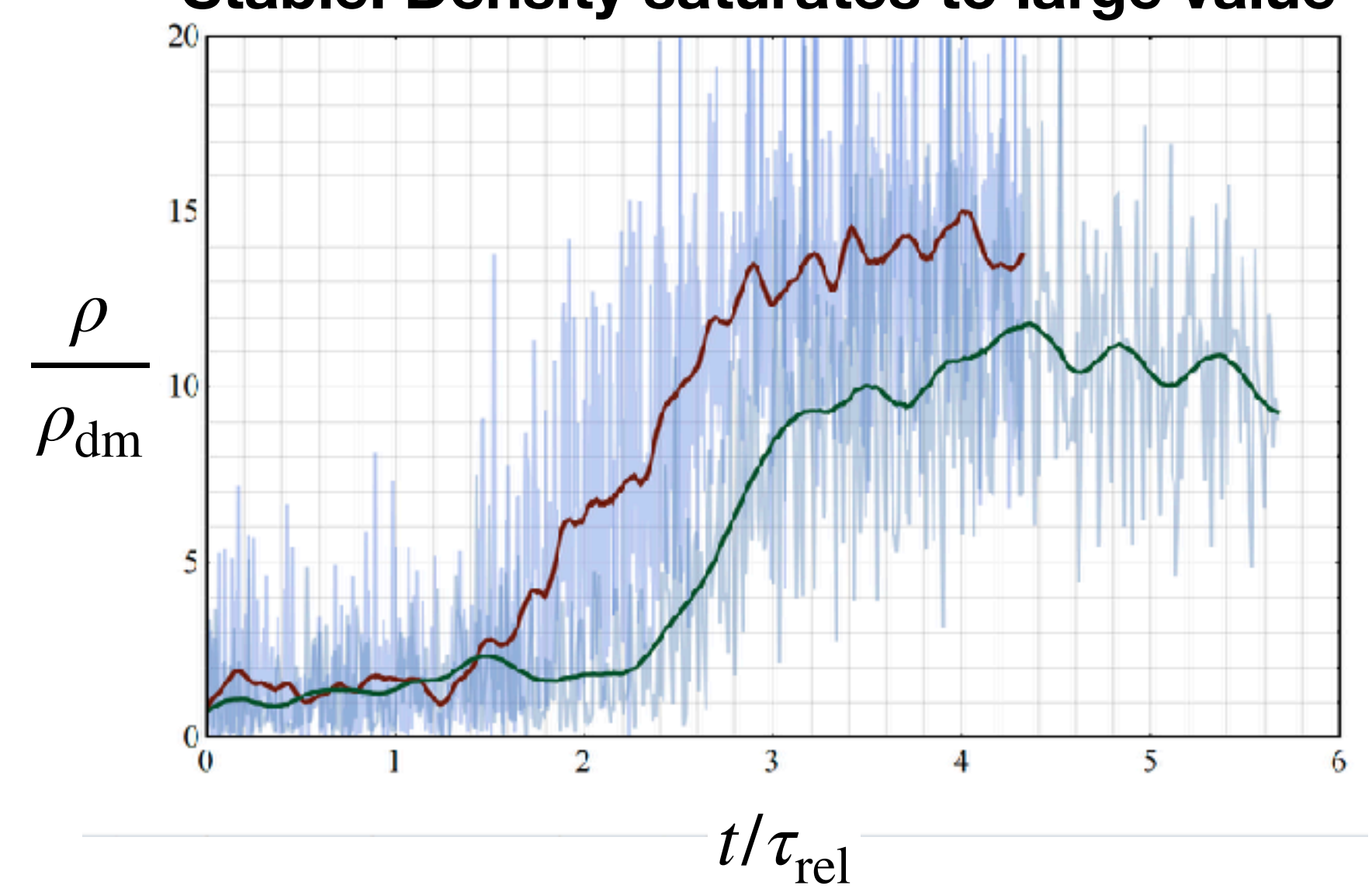
$$\rho = \rho_{\text{crit}} \sim \frac{\alpha_g^2 m_\phi^4}{|\lambda|}$$

\implies Bosenova emission of relativistic ULDM particles

Ongoing work!

Repulsive self-interaction
 $\lambda > 0$

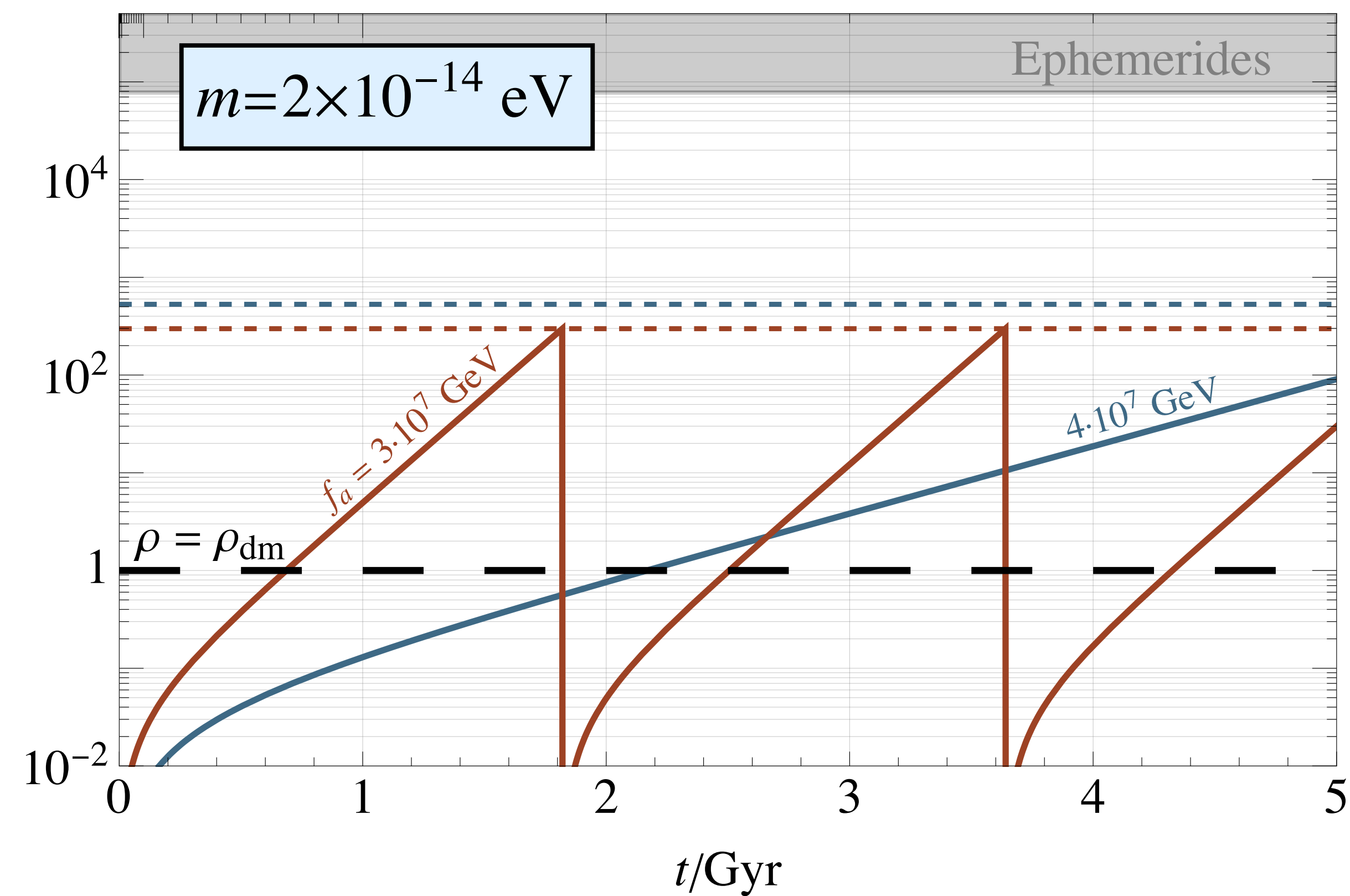
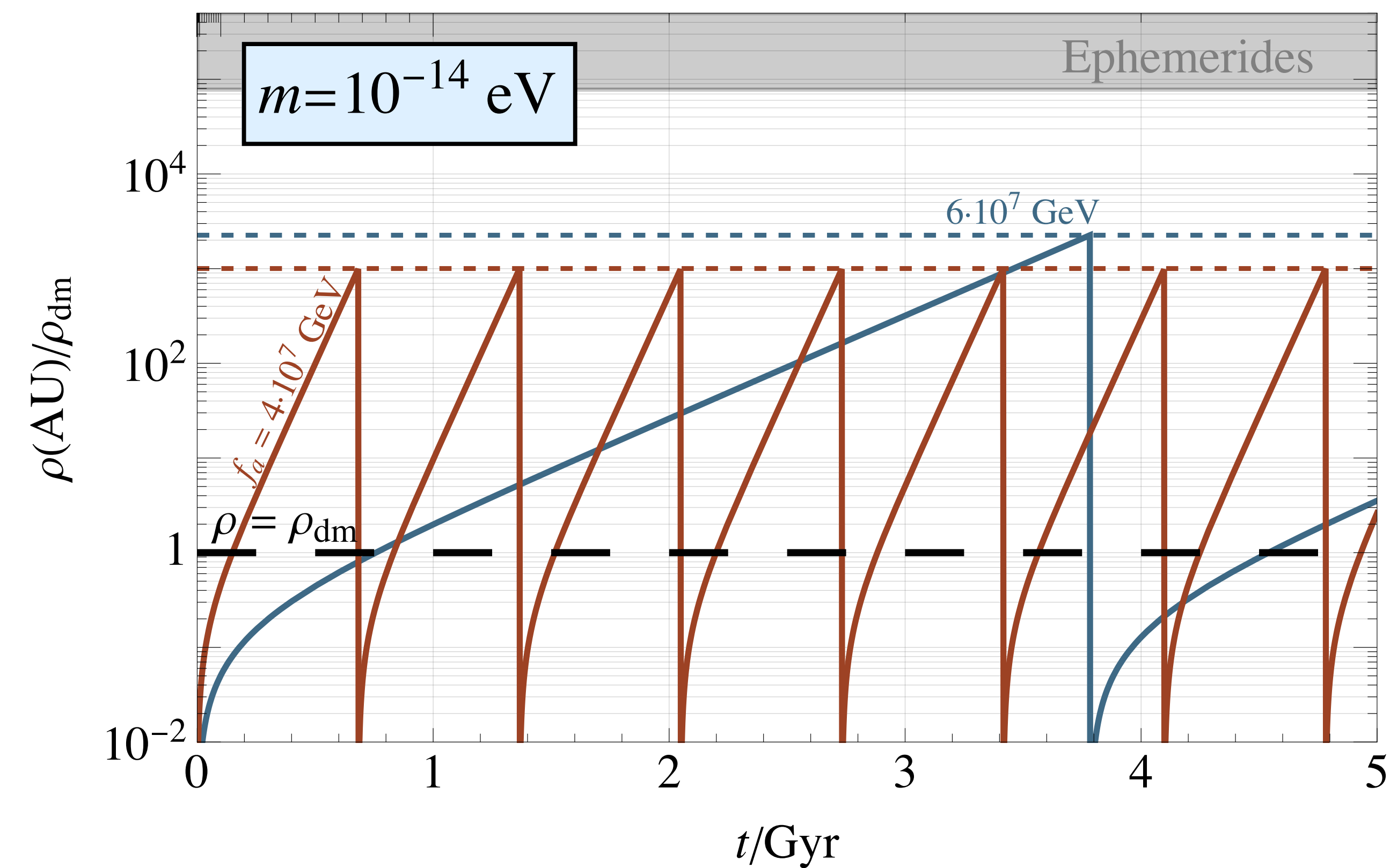
Stable! Density saturates to large value



credit: Kavli IPMU

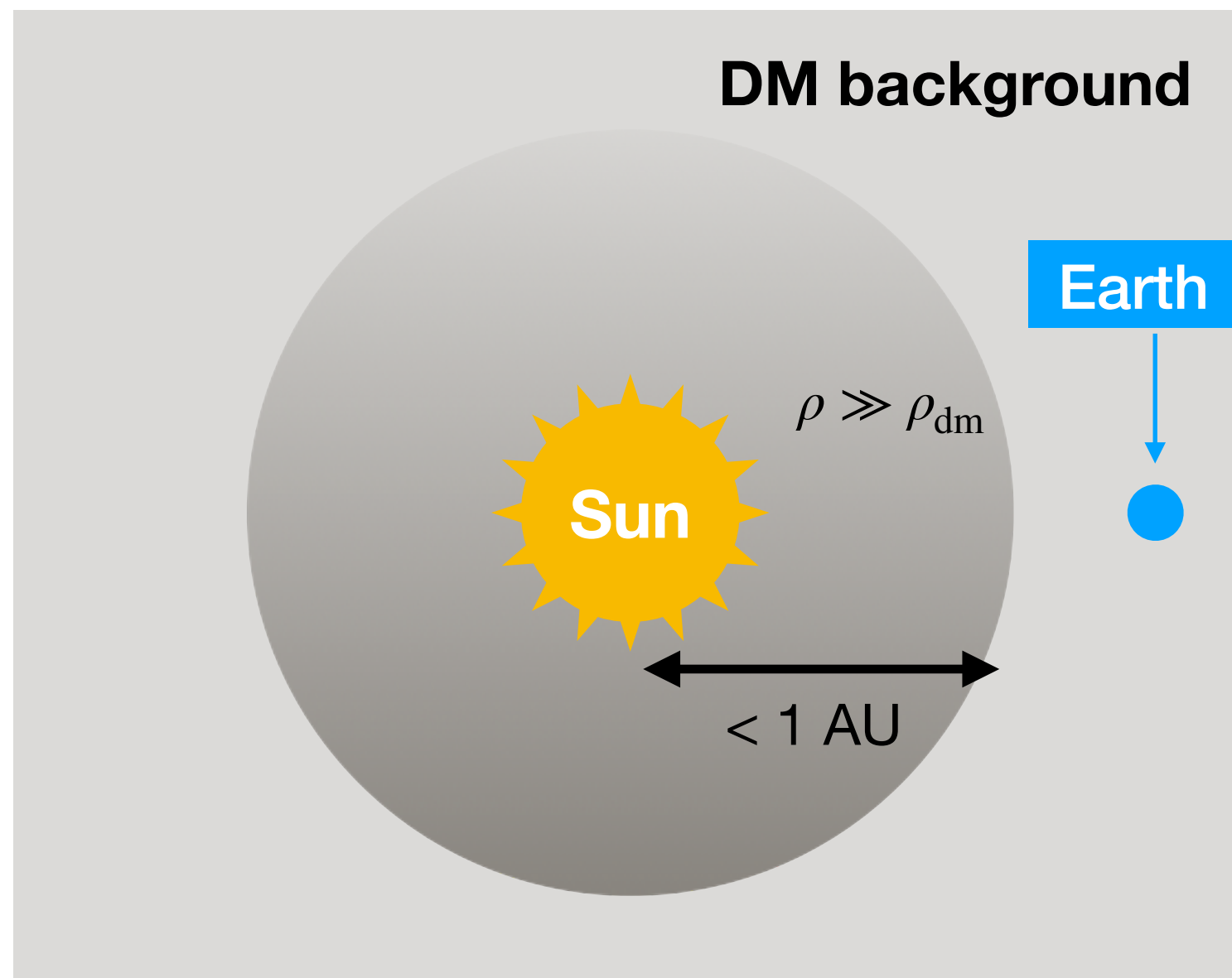
Budker, **JE**, Gorghetto,
Jiang, Perez (2306.12477)

Density over Time



assuming $\lambda < 0$;
density could be larger for $\lambda > 0$

Phenomenology of Solar / Stellar Halos



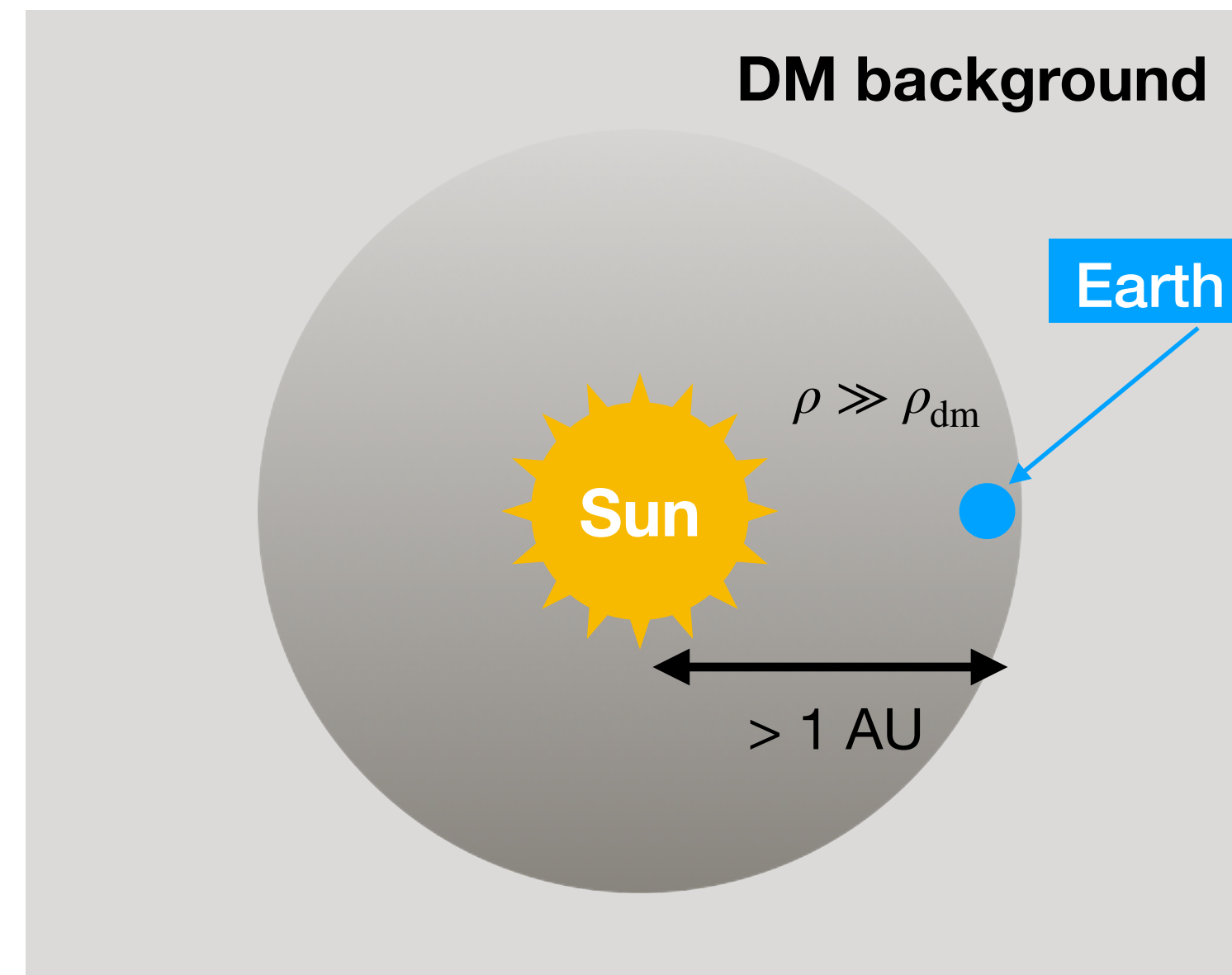
$$m_\phi \simeq 10^{-14} - 2 \cdot 10^{-13} \text{ eV}$$

Probes in Space

Tsai, **JE**, Safronova
(2112.07674)
+ Arakawa, Farnocchia
(2210.03749)

Signals from the Sun

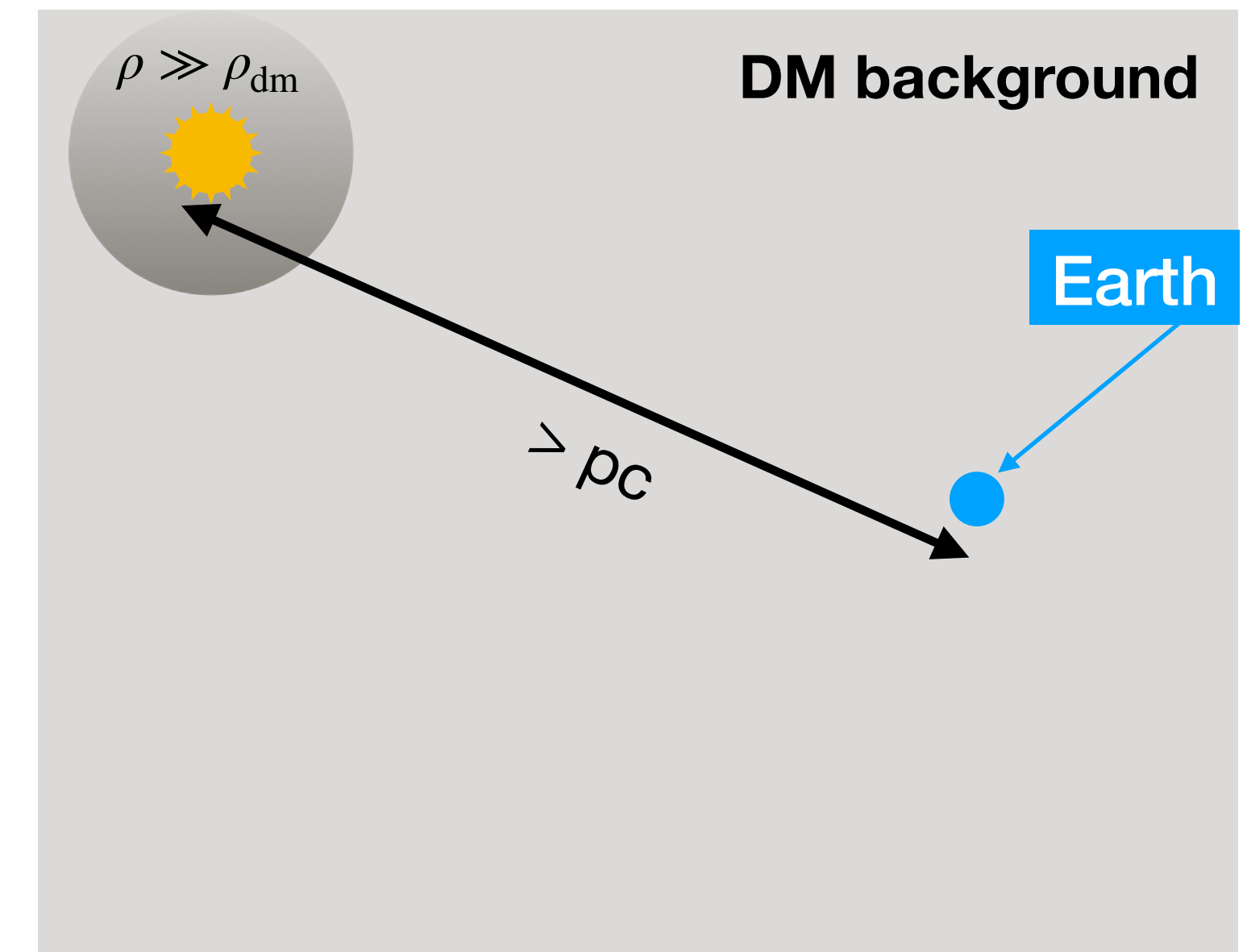
Ongoing work!



$$m_\phi \sim 10^{-15} - 3 \cdot 10^{-14} \text{ eV}$$

Probes on Earth

Banerjee, Budker, **JE**, Kim,
Perez (1902.08212)
+ Flambaum, Matsedonskyi
(1912.04295)



$$m_\phi \simeq 10^{-16} - 10^{-12} \text{ eV}$$

(supergiants to neutron stars)

Astrophysical signals

Ongoing work!

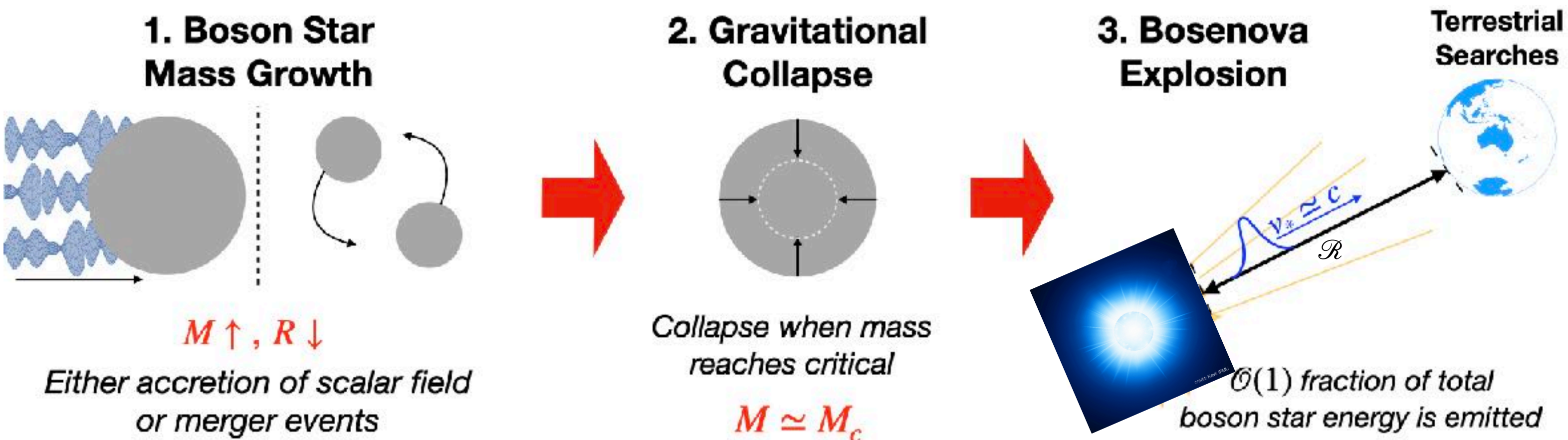
Bosenova signals

Ongoing work!

though see **JE**, Shirai, Stadnik, Takhistov (2106.14893)

Probes using Bosenovae

Widely-studied for bosenovae originating in axion star collapse



Transient Signals

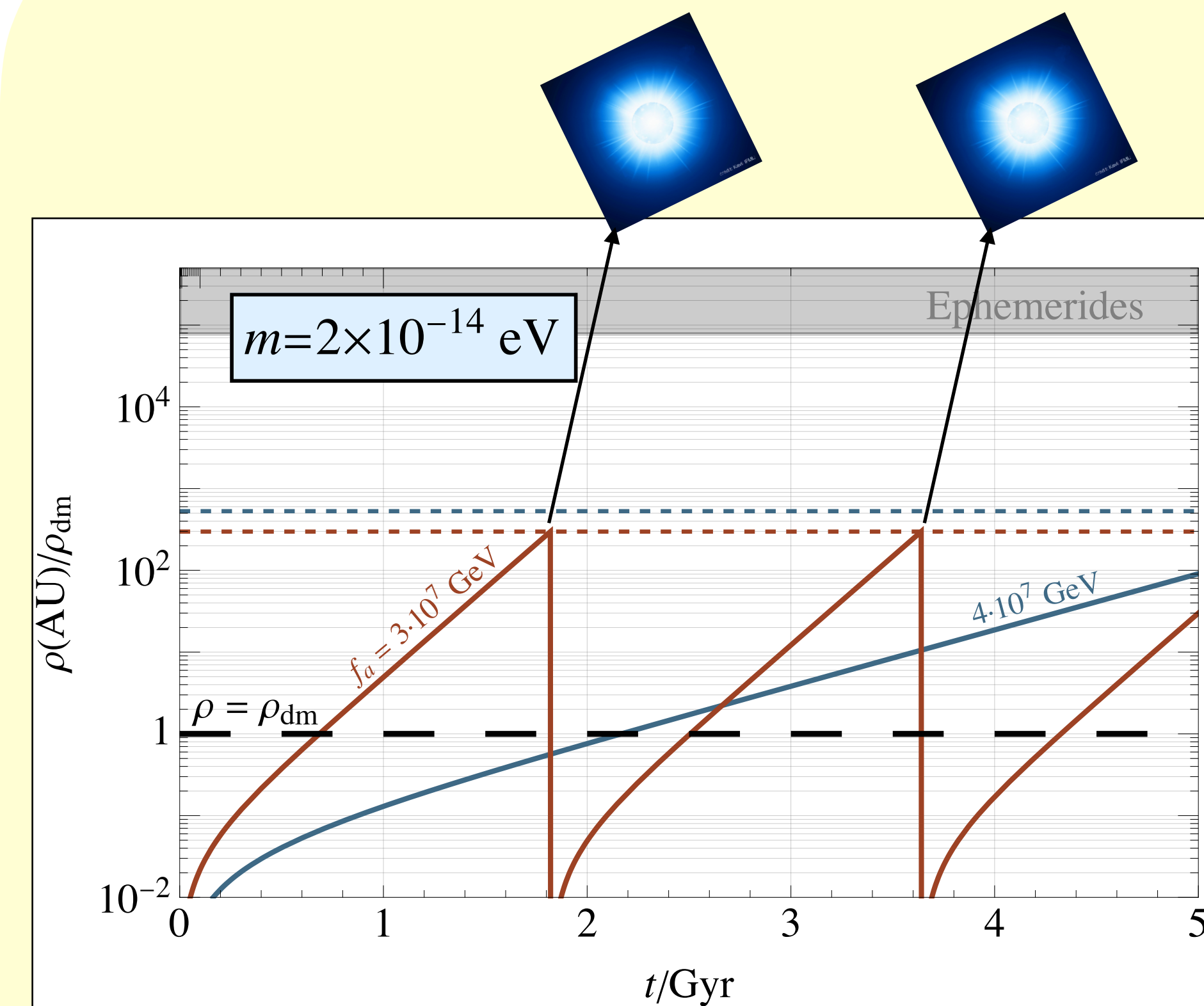
JE, Shirai, Stadnik, Takhistov (2106.14893)

Arakawa, **JE**, Safronova, Takhistov, Zaheer (2306.16468, 2402.06736)

Diffuse Axion Background

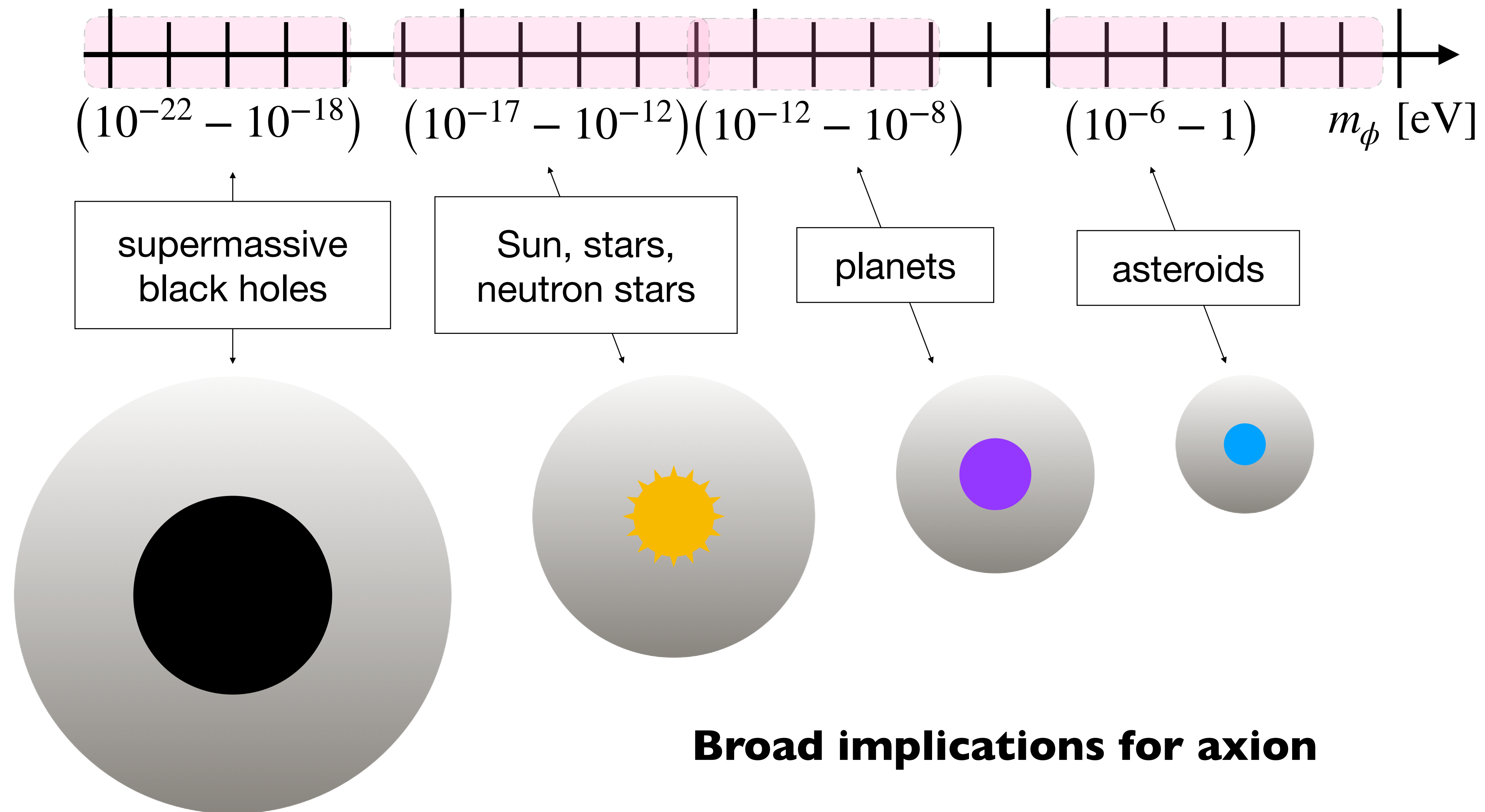
JE, Takhistov (2402.00100)

What about solar halo collapse?



Multiple, repeating bosenovae on stars all over the galaxy!

Gravitational Atoms Everywhere!



Broad implications for axion physics across model space!

Conclusions

Axions / ULDM **captured gravitationally** to astrophysical objects (e.g. the Sun), forming a gravitational atom

Complementary to superradiance: occurs in very different astrophysical environments

Formation occurs **through $2 \rightarrow 2$ scattering; generic mechanism** easily implies gravitational atoms elsewhere

When it forms, a **solar halo** gives rise to striking new targets for experiment:

- On Earth: enhanced density and coherence time for $m_\phi \sim \text{few} \cdot 10^{-15} - \text{few} \cdot 10^{-14} \text{ eV}$
- Future space missions can probe small, dense solar halos for $m_\phi \sim 10^{-13} - 10^{-14} \text{ eV}$
- **Astrophysical signals in other star systems: wider range of m_ϕ**

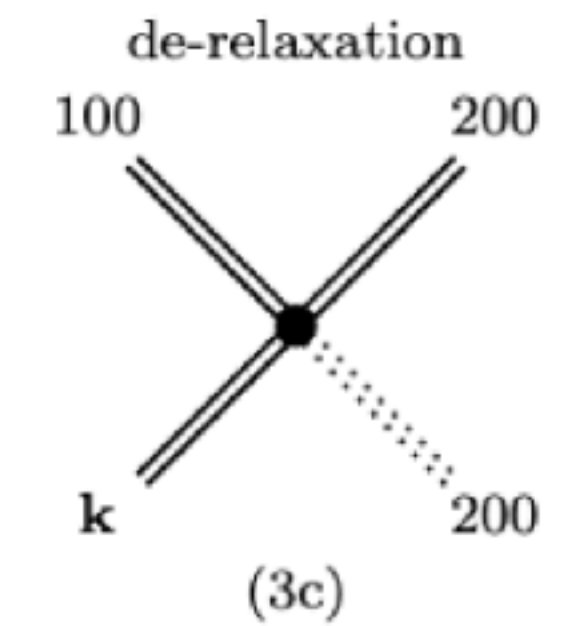
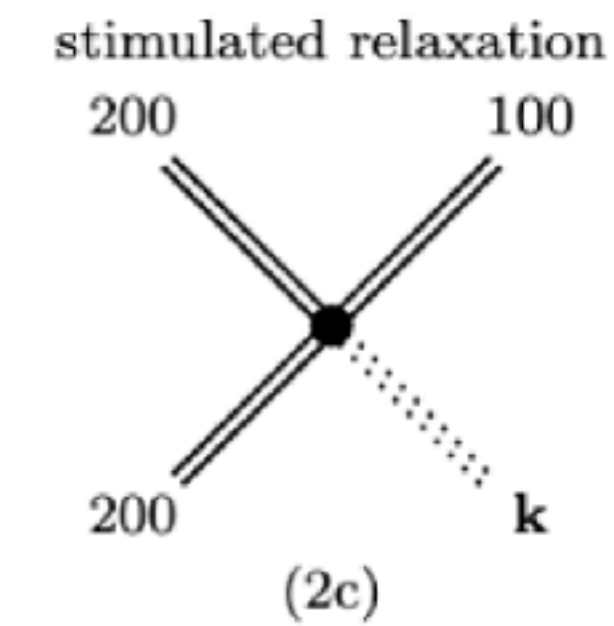
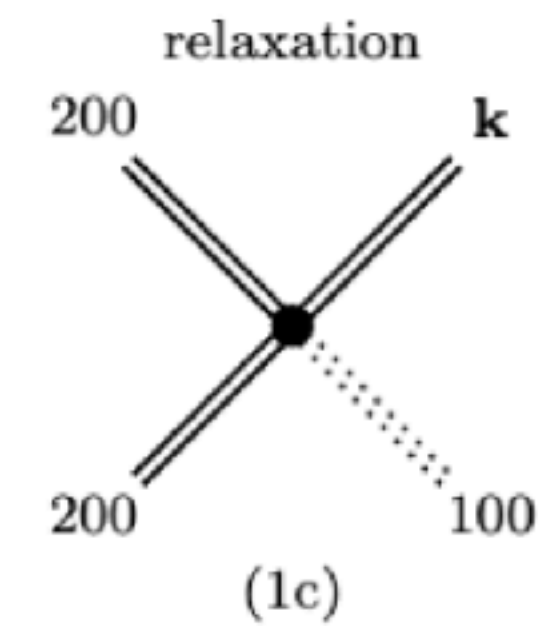
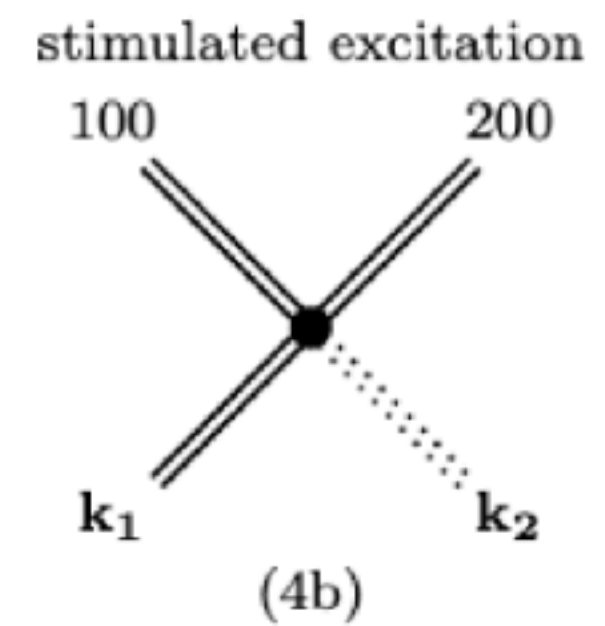
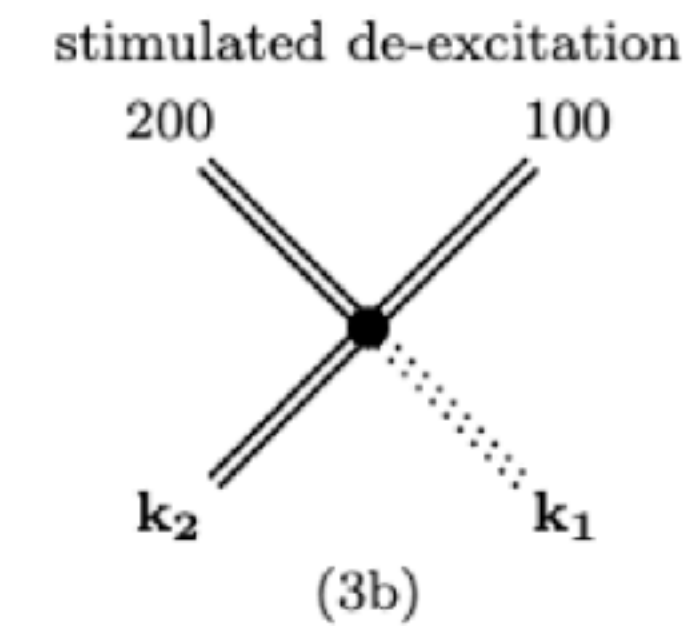
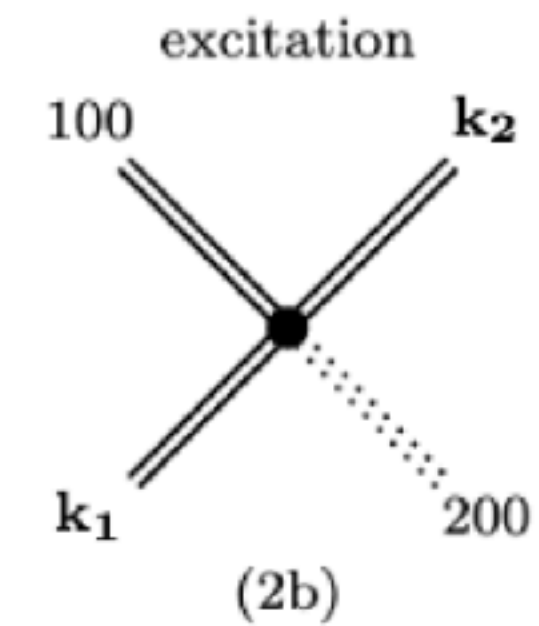
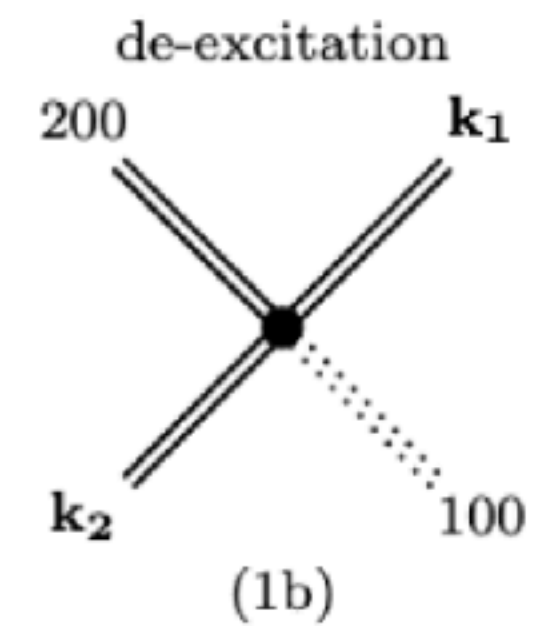
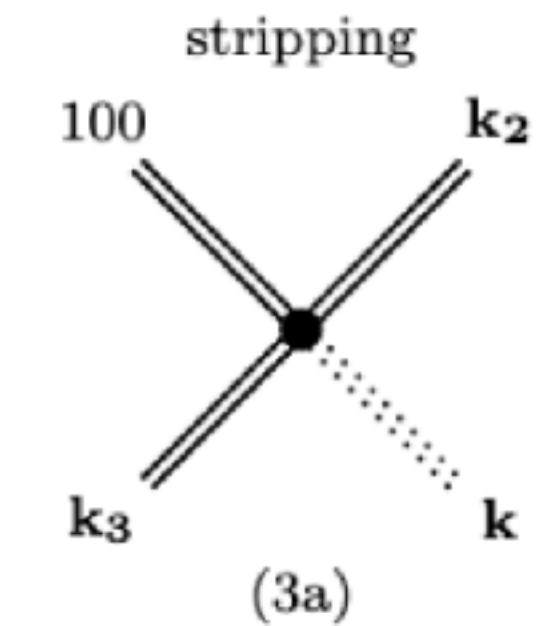
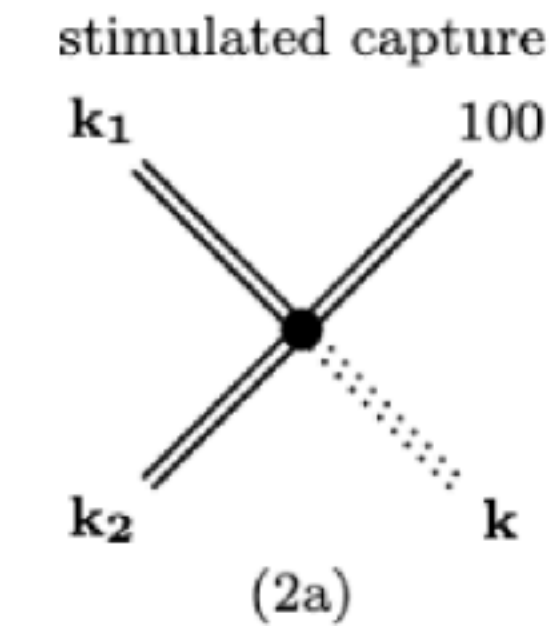
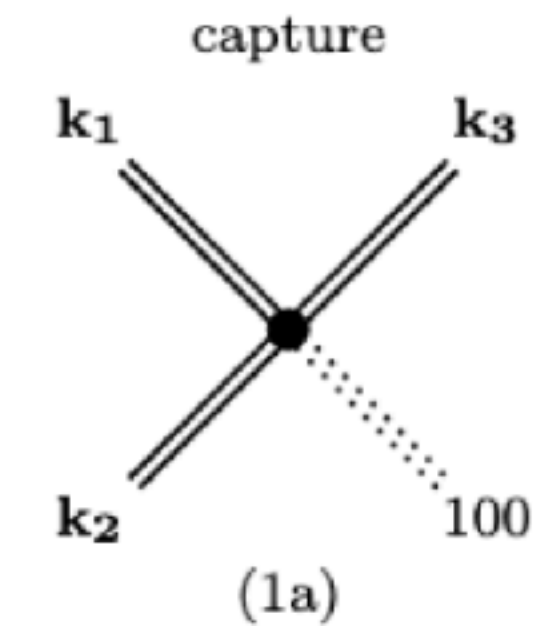
Thank you for your attention!

Bonus Round

“What about the Excited States?”

C Role of excited states and two-level system

In the main text we studied the capture of ULDM to the ground state, but neglected excited states in most of our discussions. In this Appendix, we show that these are likely to only enhance the accretion of DM onto the ground state, and not to change the overall conclusions of Section 4.

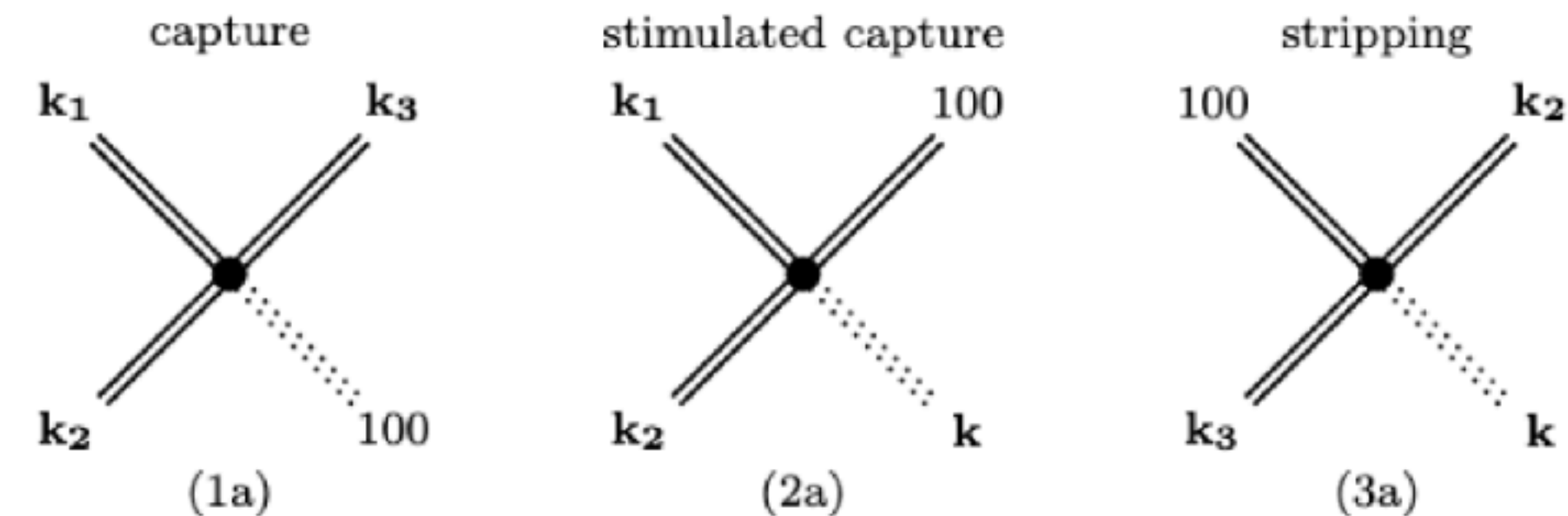


2-level system:
 $nlm = 100, 200$

“What about the Excited States?” (2)

C Role of excited states and two-level system

In the main text we studied the capture of ULDM to the ground state, but neglected excited states in most of our discussions. In this Appendix, we show that these are likely to only enhance the accretion of DM onto the ground state, and not to change the overall conclusions of Section 4.



de-excitation

excitation

stimulated de-excitation

stimulated excitation

$$\begin{aligned}
 \dot{N}_{100} = & 2g^2 \int [dk_1][dk_2][dk_3] |\mathcal{M}_{k_1+k_2 \rightarrow k_3+100}|^2 (2\pi) \delta(\omega_{k_1} + \omega_{k_2} - \omega_{100} - \omega_{k_3}) \\
 & \{f(\mathbf{k}_1)f(\mathbf{k}_2)f(\mathbf{k}_3) + N_{100}[f(\mathbf{k}_1)f(\mathbf{k}_2) - 2f(\mathbf{k}_2)f(\mathbf{k}_3)]\} \\
 & + 4g^2 \int [dk_1][dk_2] |\mathcal{M}_{k_1+100 \rightarrow k_2+200}|^2 (2\pi) \delta(\omega_{k_1} + \omega_{100} - \omega_{200} - \omega_{k_2}) \\
 & \{N_{100}N_{200}[f(\mathbf{k}_2) - f(\mathbf{k}_1)] - f(\mathbf{k}_1)f(\mathbf{k}_2)[N_{100} - N_{200}]\} \\
 & + 2g^2 \int [dk] |\mathcal{M}_{200+200 \rightarrow 100+k}|^2 (2\pi) \delta(2\omega_{200} - \omega_{100} - \omega_k) \{f(\mathbf{k})N_{200}^2 + N_{100}N_{200}^2 - 2N_{200}N_{100}f(\mathbf{k})\}.
 \end{aligned}
 \tag{103}$$

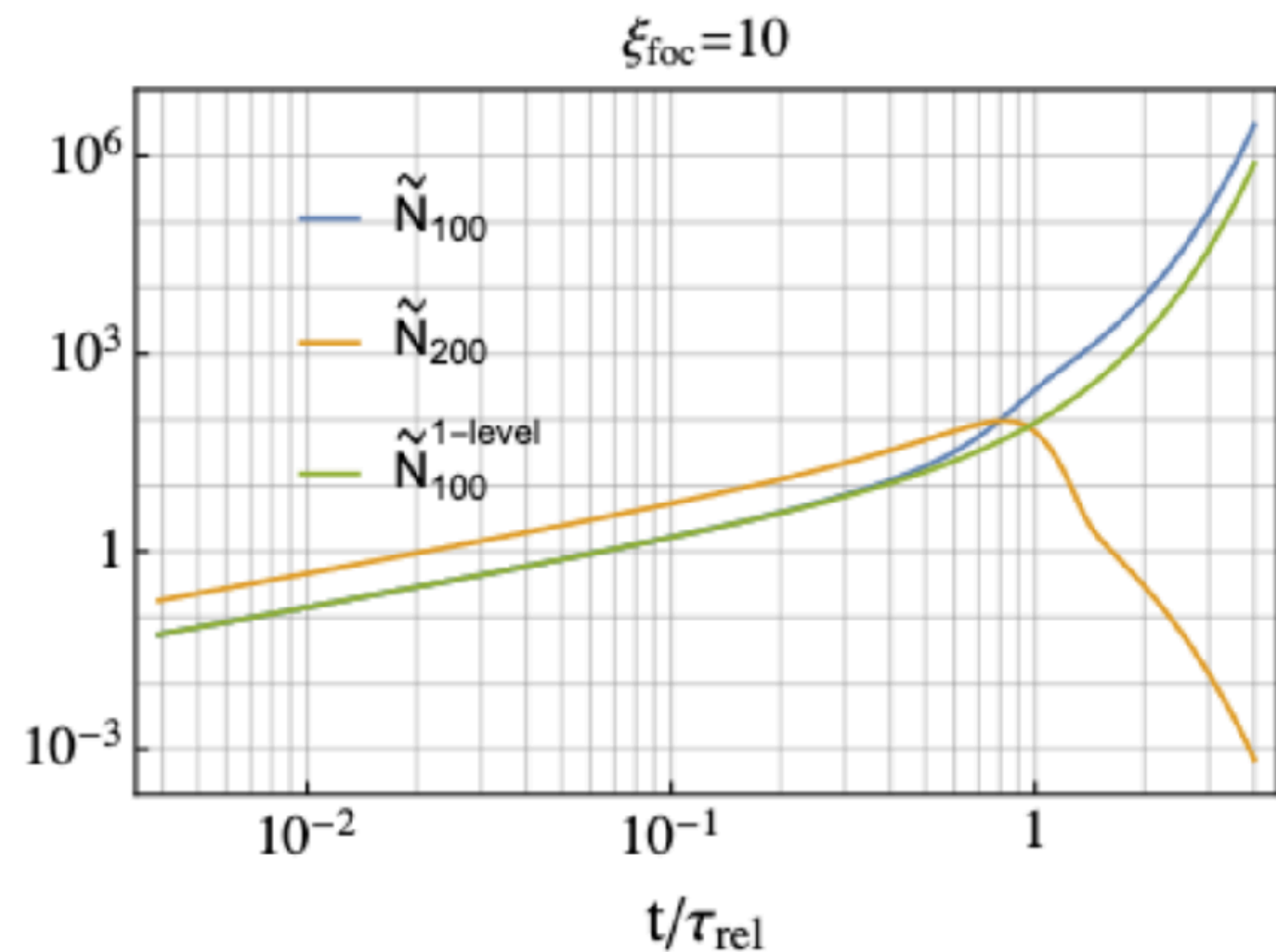
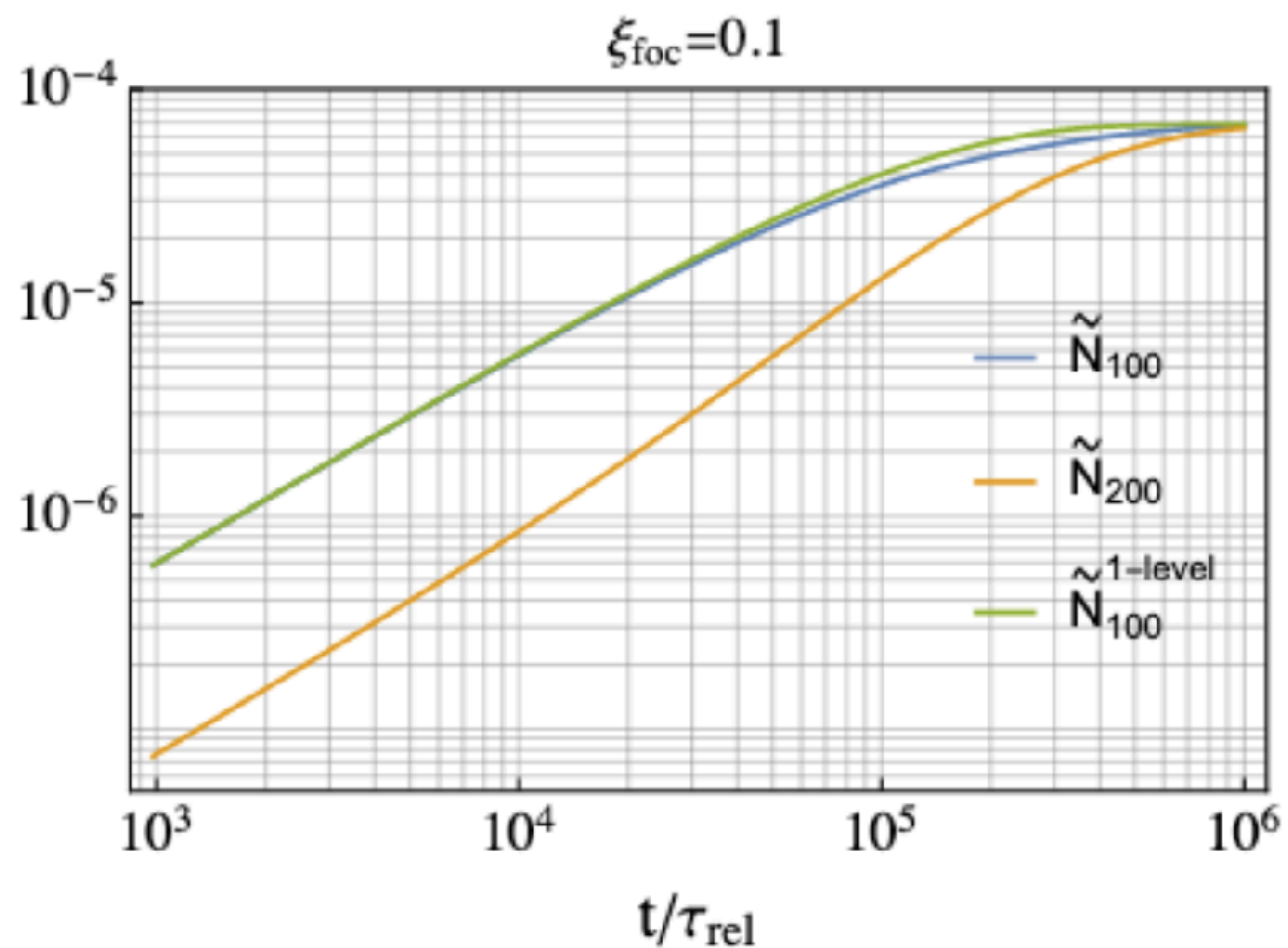
(1c)

(2c)

(3c)

2-level system:
 $nlm = 100, 200$

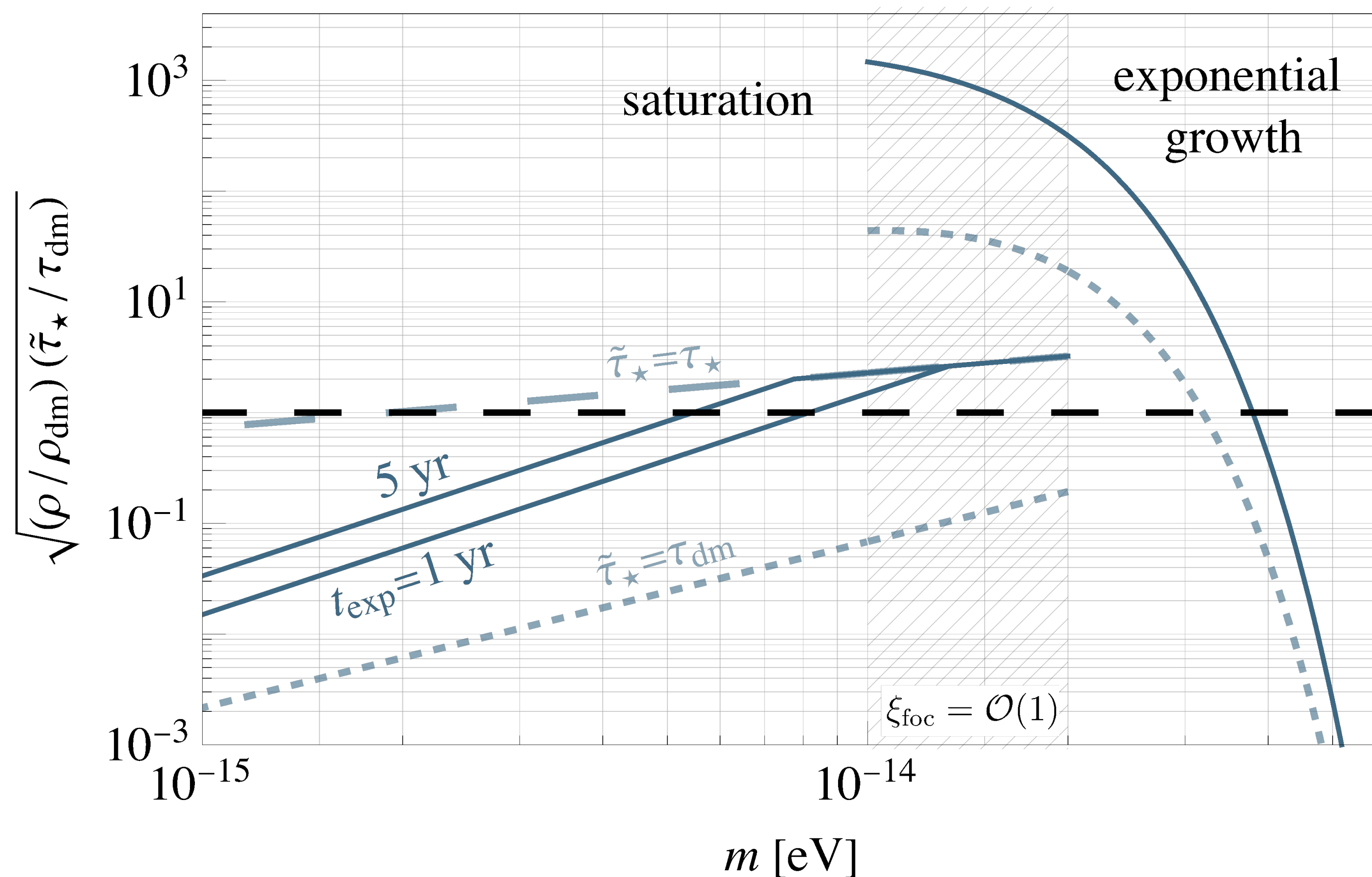
“What about the Excited States?” (3)



“What about the Dilute Atoms?”

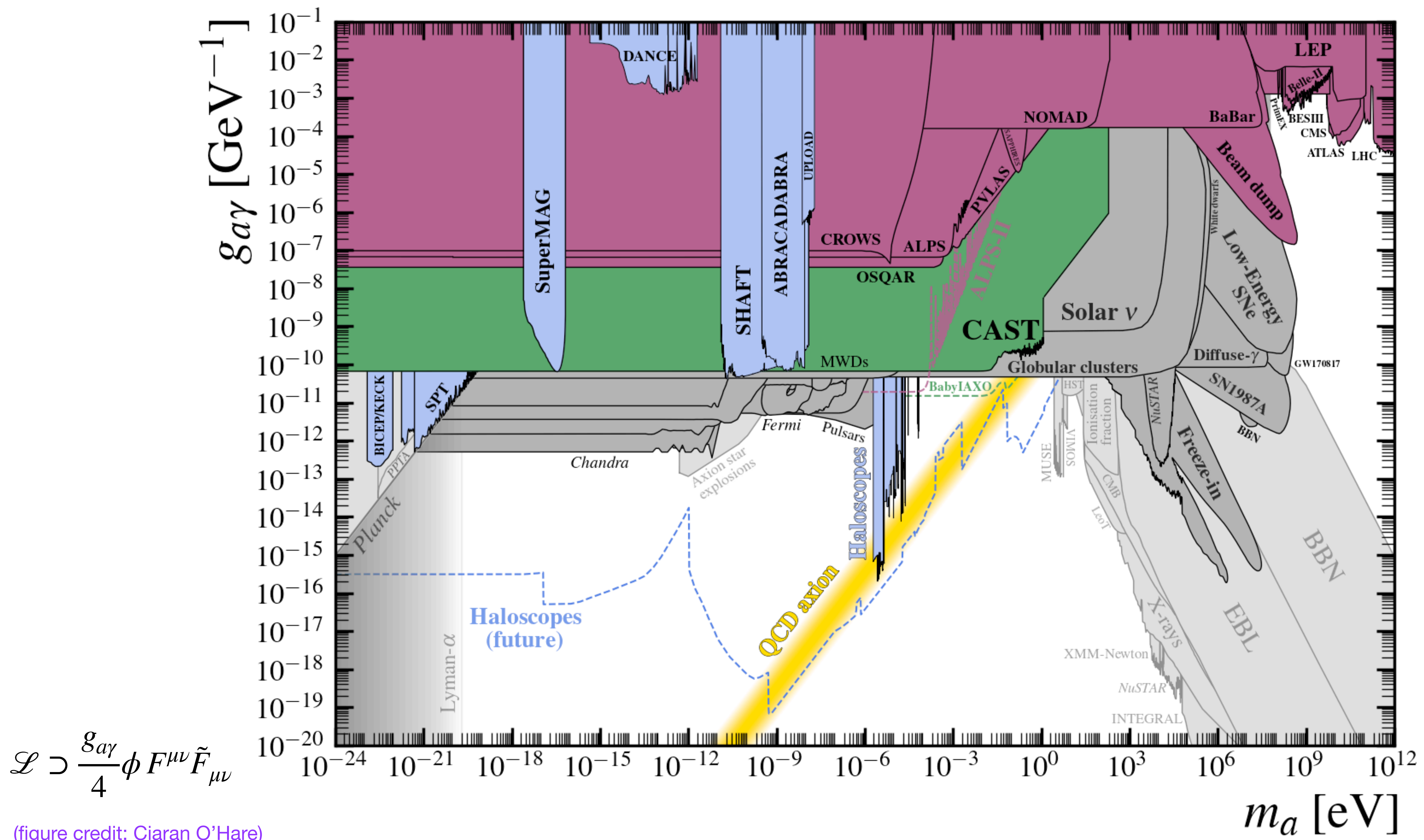
Signal in direct search typically grows as $\propto \sqrt{\rho t}$ as long as

$$t \lesssim \tau_\star \simeq \frac{2\pi}{m_\phi \alpha_g^2} = \left[\frac{2\pi}{\xi_{\text{foc}}} \right]^2 \tau_{\text{dm}} \simeq 2 \text{ year} \left[\frac{10^{-14} \text{ eV}}{m_\phi} \right]^3$$



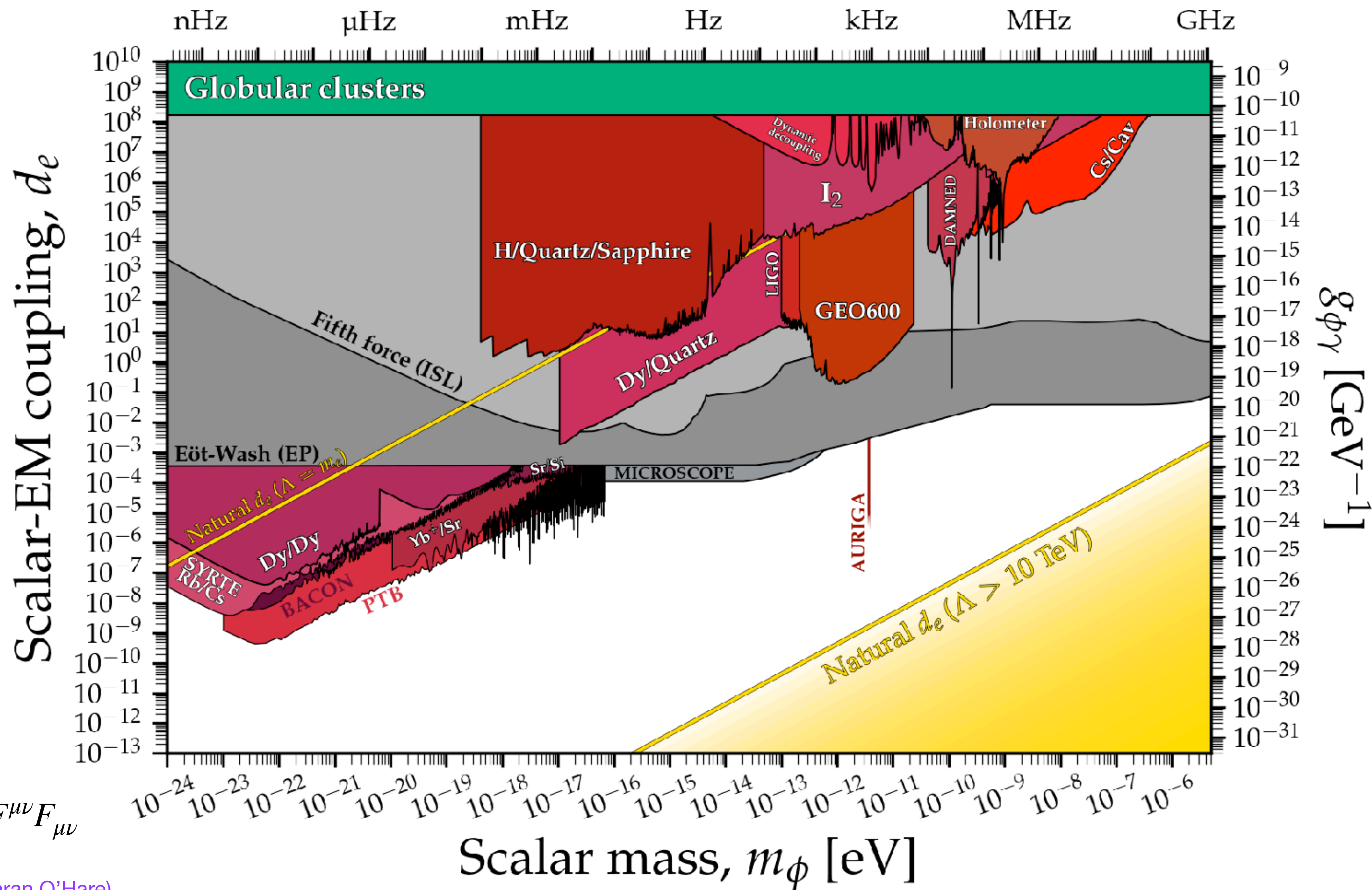
$$\tilde{\tau}_\star \equiv \min(\tau_\star, t_{\text{exp}})$$

assuming $\lambda < 0$;
density could be larger for $\lambda > 0$



$$\mathcal{L} \supset \frac{g_{a\gamma}}{4} \phi F^{\mu\nu} \tilde{F}_{\mu\nu}$$

(figure credit: Ciaran O'Hare)



$$\mathcal{L} \supset \frac{g_{\phi\gamma}}{4} \phi F^{\mu\nu} F_{\mu\nu}$$

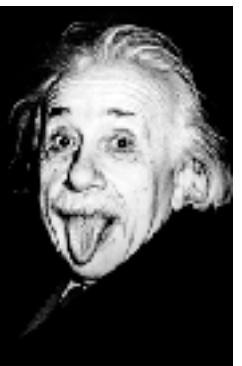
(figure credit: Ciaran O'Hare)

Probing the Very Local DM Density

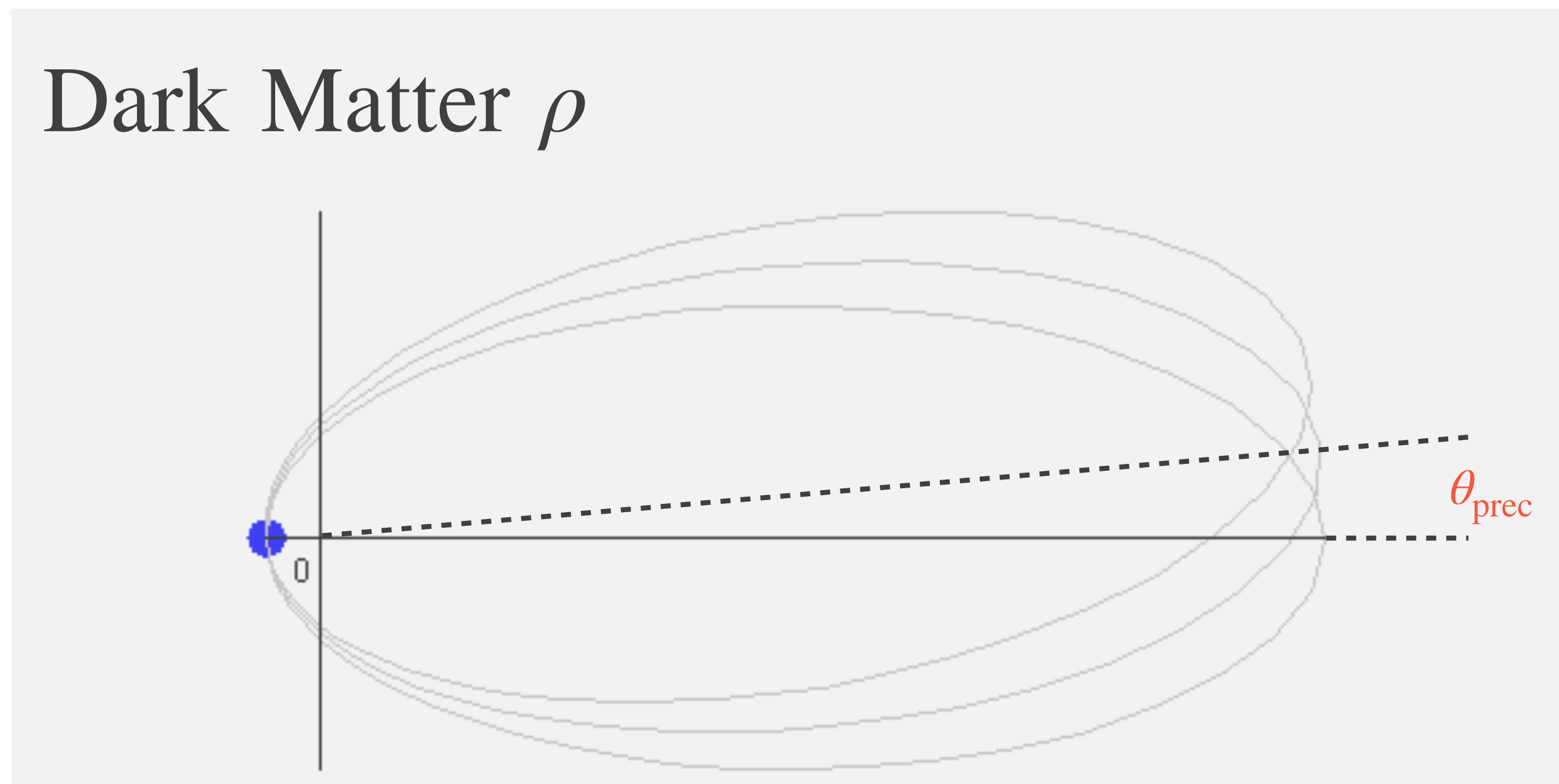
(i.e. in our solar system)

Sources of perihelion precession:

- ☉ 3-body interactions
 - Jupiter, Venus, Earth 91 %
 - Other planets 1 %
 - Asteroids 10^{-6}
- ☉ Solar oblateness 0.1 %
- ☉ General Relativistic effects 8 %



- ☉ **“Extra” (dark) mass** $< 10^{-7}$



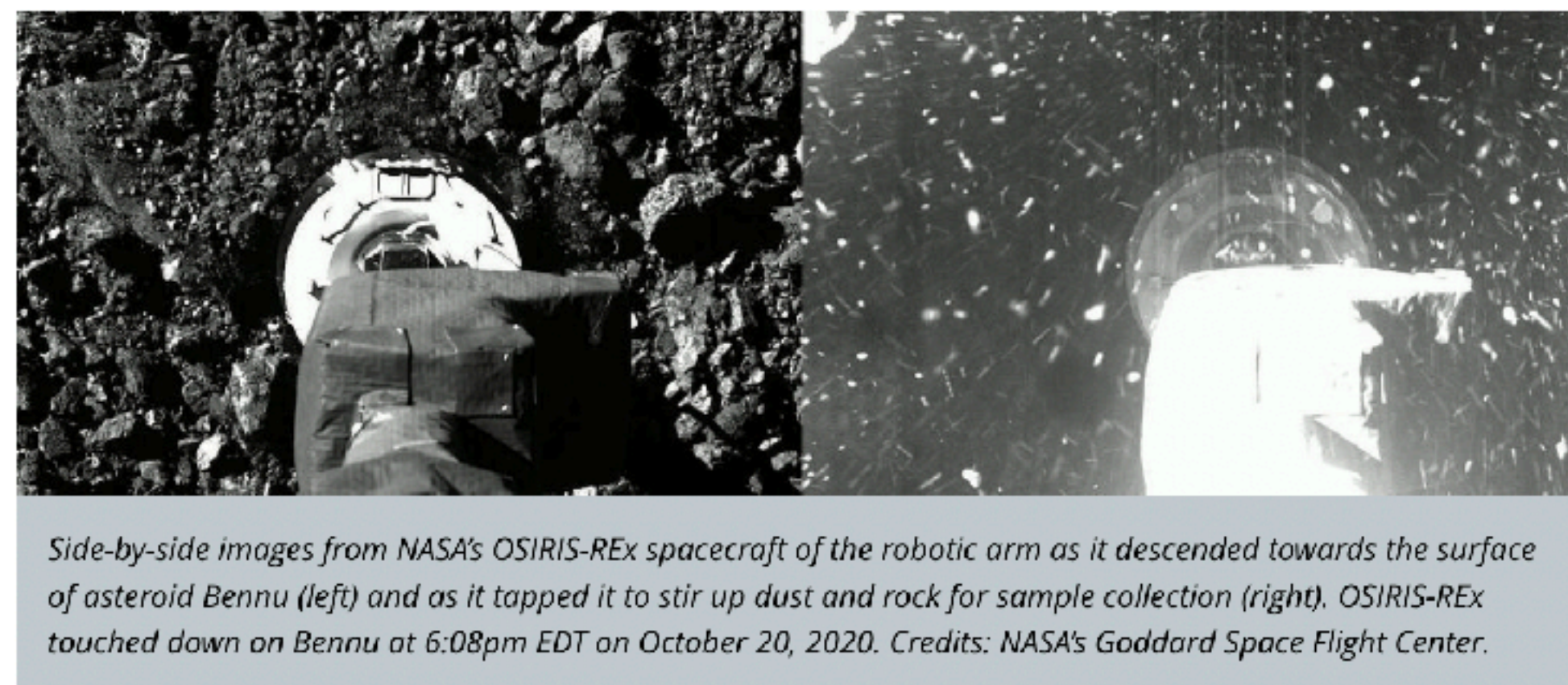
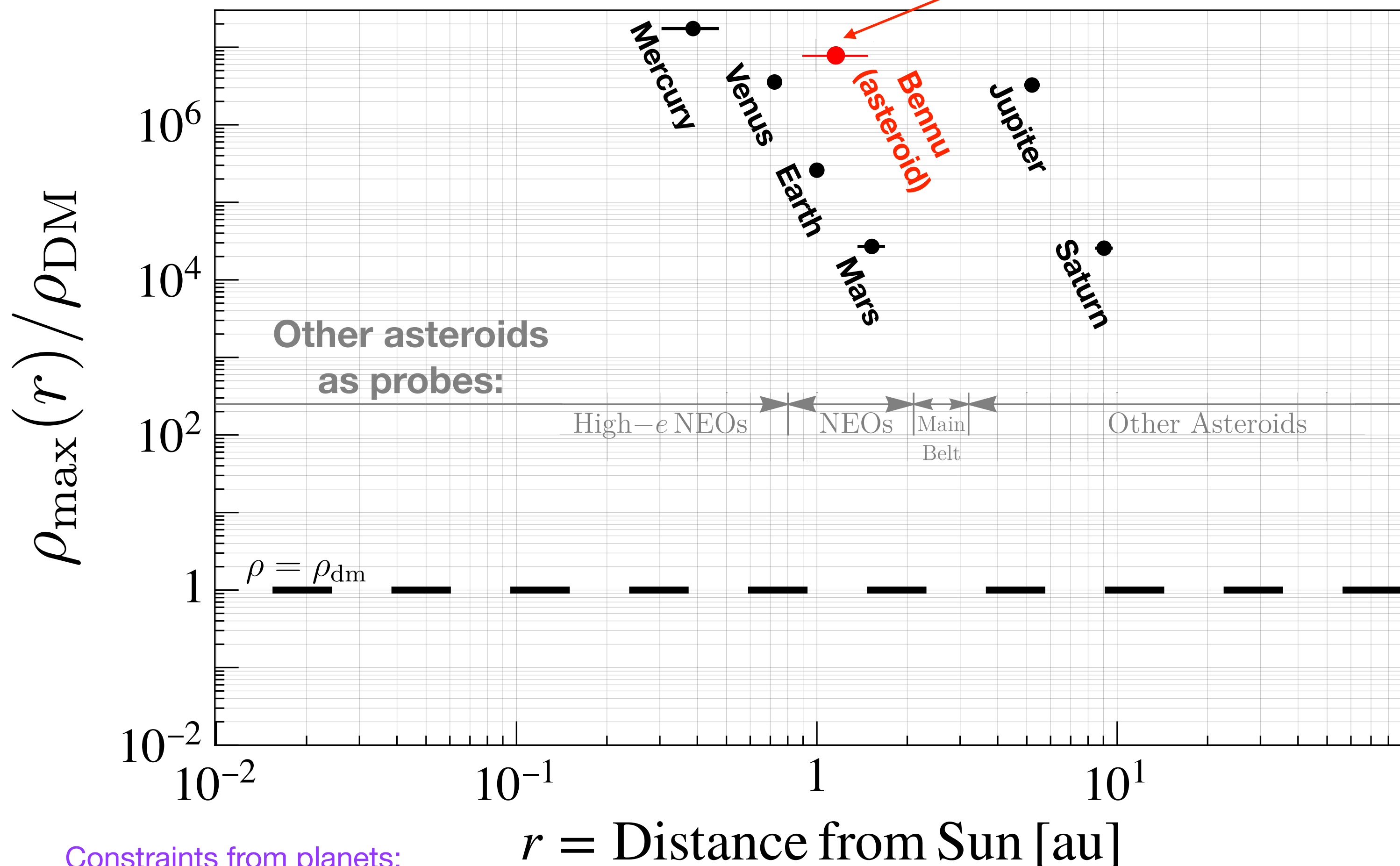
Consider **Mercury**: $\dot{\theta}_{\text{prec}} = 575.3100 \pm 0.0015 \frac{\text{arcsec}}{\text{century}}$

See Park et al. (*MESSENGER* collaboration),
The Astronomical Journal, 153:121 (2017)

Constraints on DM in our Solar System

First constraint using asteroids:
Tsai, **JE**, Arakawa, Farnocchia, Safronova (2210.03749)

Bennu: exceptional tracking data including dedicated OSIRIS-REx mission from NASA



Side-by-side images from NASA's OSIRIS-REx spacecraft of the robotic arm as it descended towards the surface of asteroid Bennu (left) and as it tapped it to stir up dust and rock for sample collection (right). OSIRIS-REx touched down on Bennu at 6:08pm EDT on October 20, 2020. Credits: NASA's Goddard Space Flight Center.

Of general interest:

Is there DM in the Solar System? How much?

How do we strengthen such constraints?

New ideas?

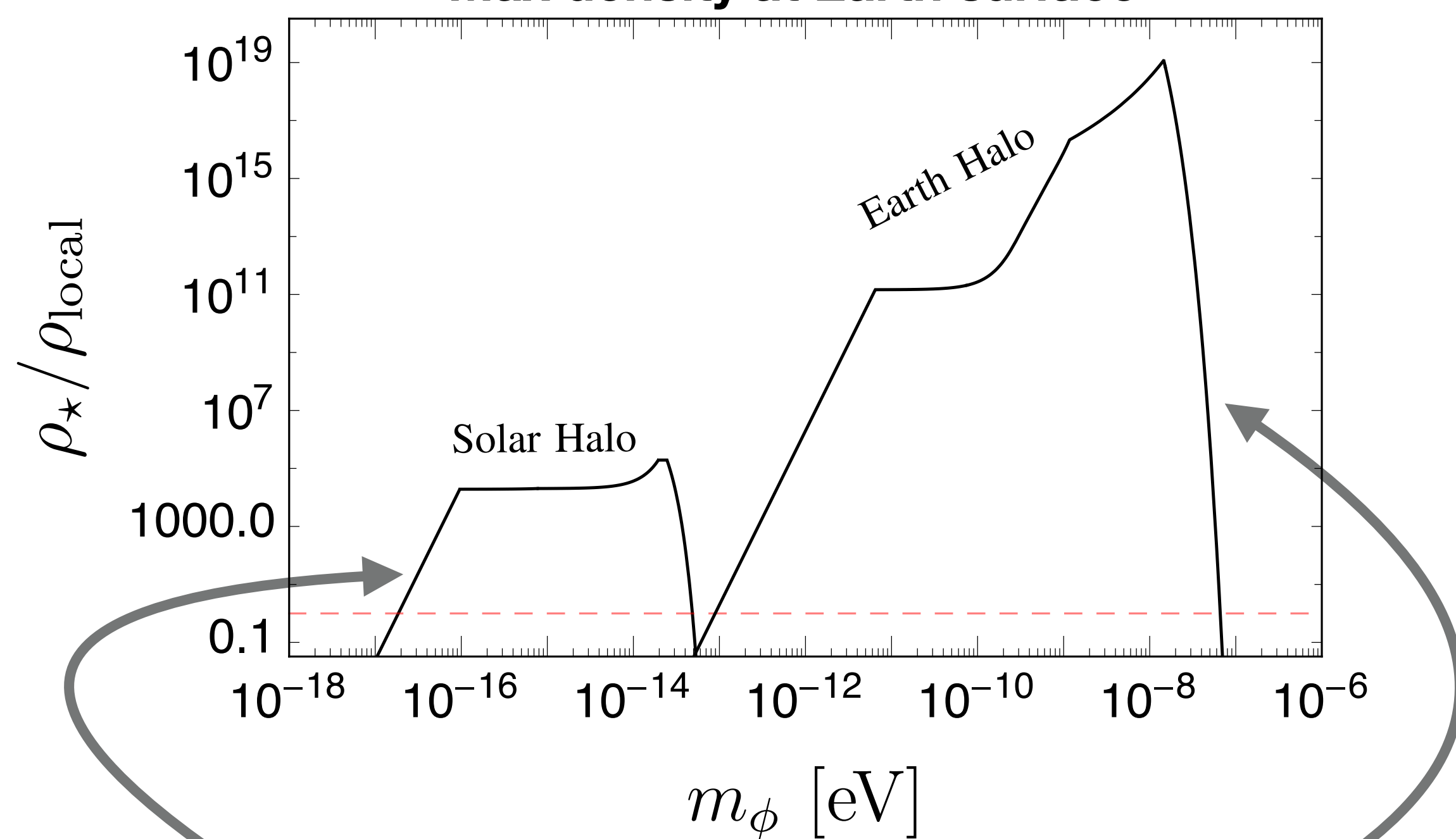
Constraints from planets:
see Pitjeva and Pitjeva (1306.5534)

Possible Overdensity

Banerjee, Budker, JE, Kim, Perez (1902.08212)

Density can be very much enhanced relative to 'naive' expectation ρ_{local}

Max density at Earth surface



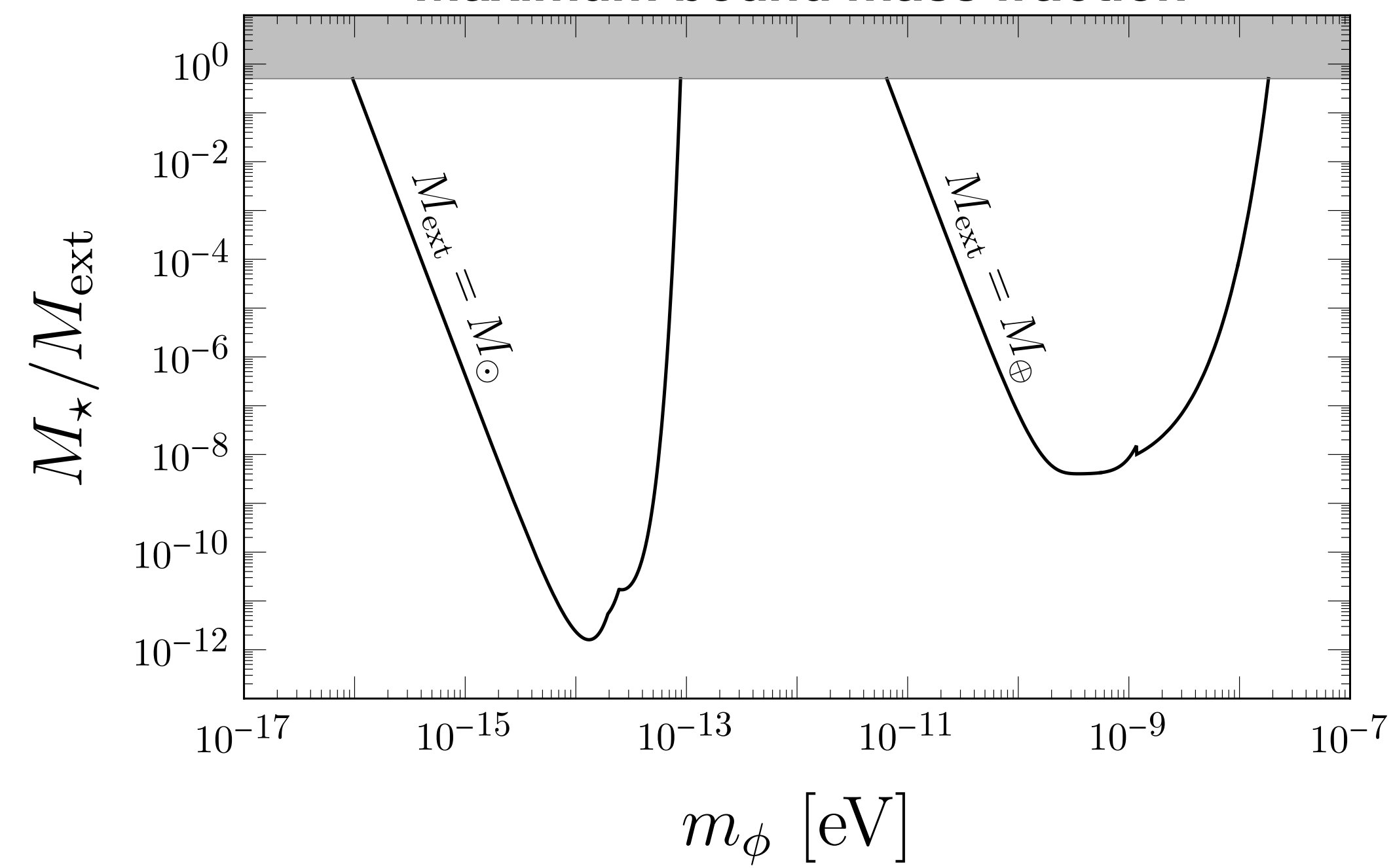
*Solar System Ephemerides
(Mercury, Mars, Saturn)*

Pitjev and Pitjeva (1306.5534)

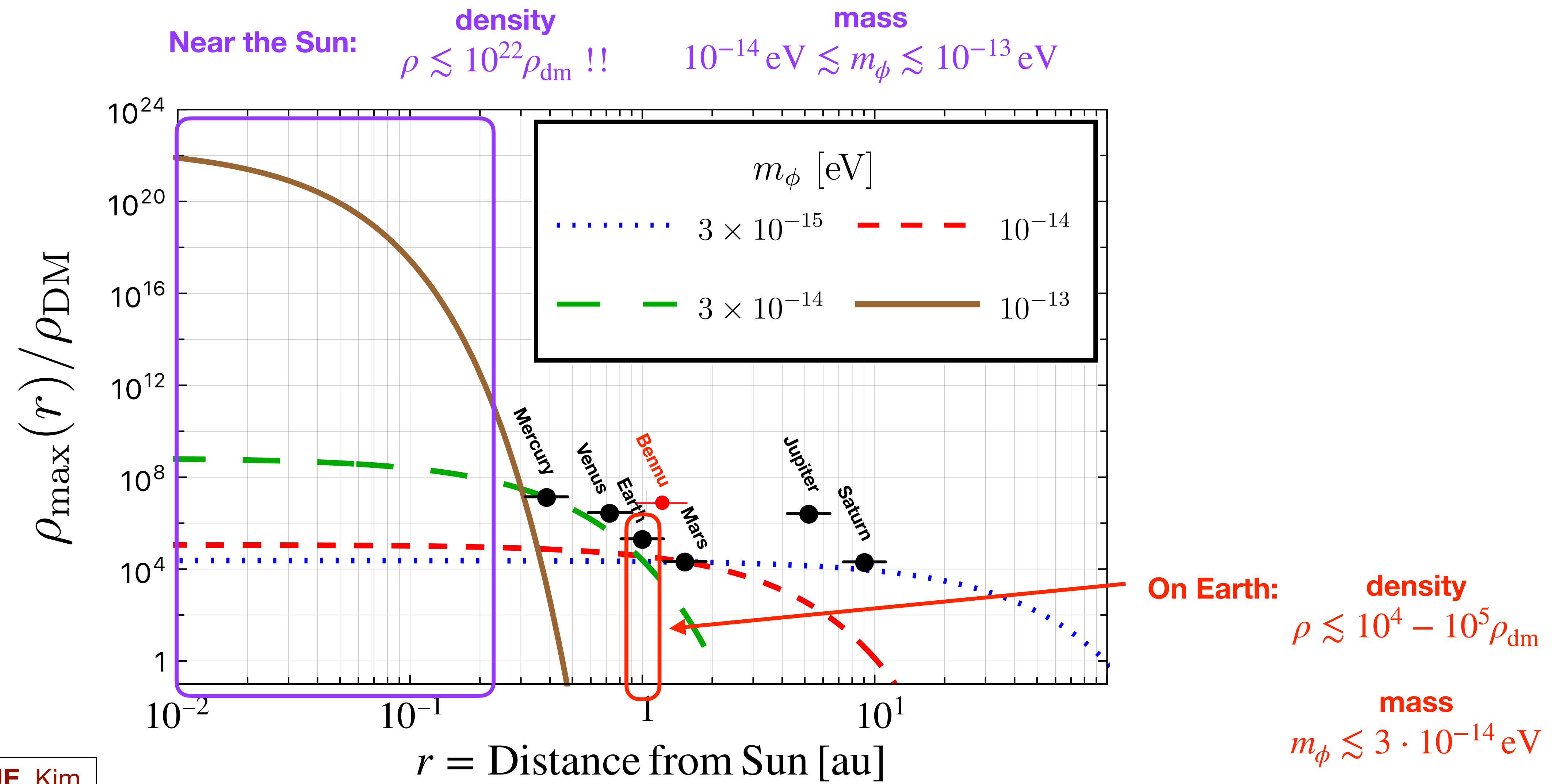
*Lunar Laser Ranging
+ LAGEOS Satellite*

Adler (0808.0899)

Maximum bound mass fraction



Maximum Density of Solar Halo



Banerjee, Budker, **JE**, Kim, Perez (1902.08212)

“But can bound states really form?”

Formed by “gravitational cooling”

Seidel and Suen (gr-qc/9309015)

Guzman and Urena-Lopez (astro-ph/0603613)

Can be understood analytically as gravitational relaxation of quasiparticles

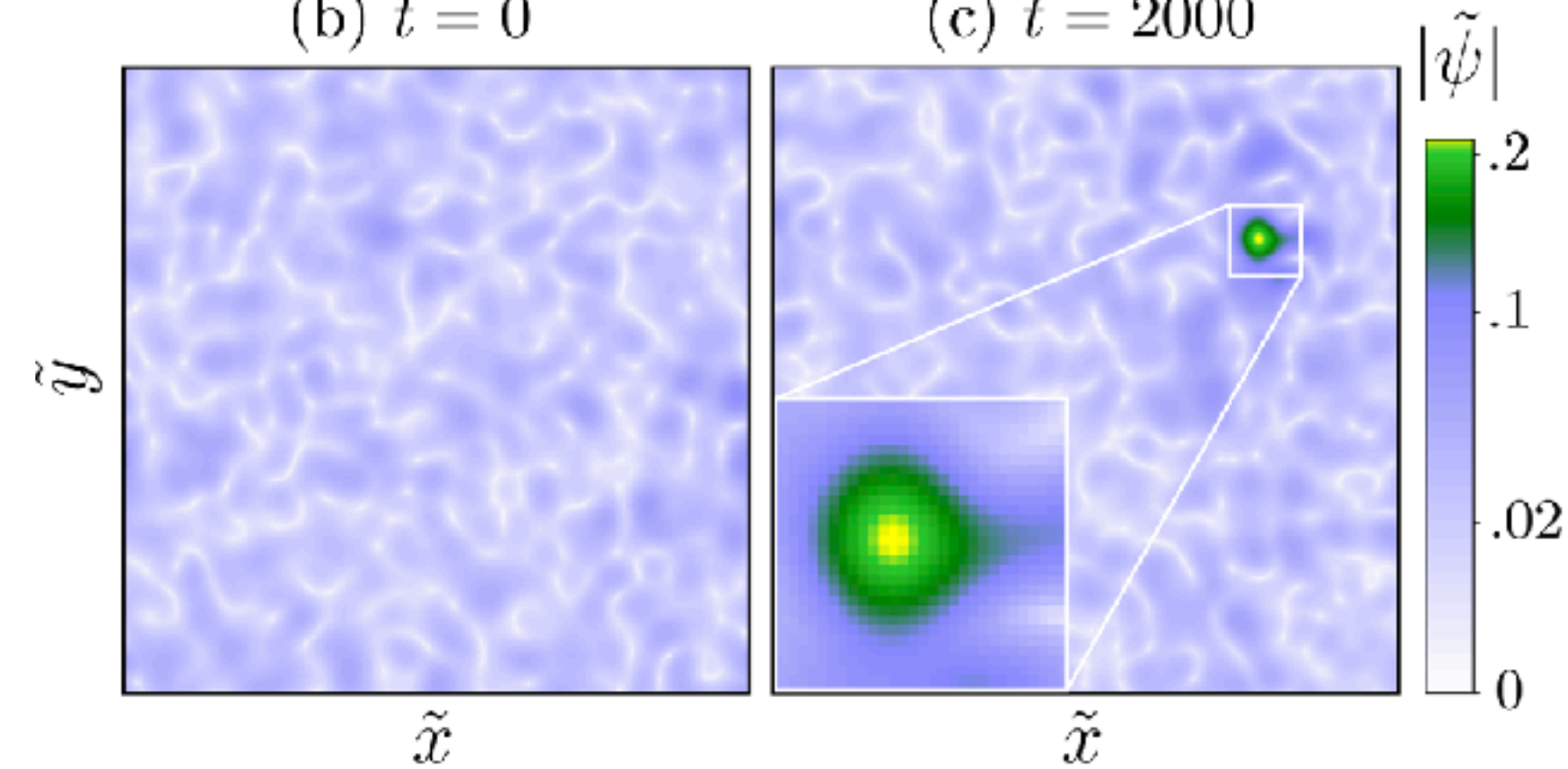
Hui, Ostriker, Tremaine, Witten (1610.08297)

Bar-Or, Fouvry, Tremaine (1809.07673)

Levkov, Panin, Tkachev (1804.05857)

(b) $\tilde{t} = 0$

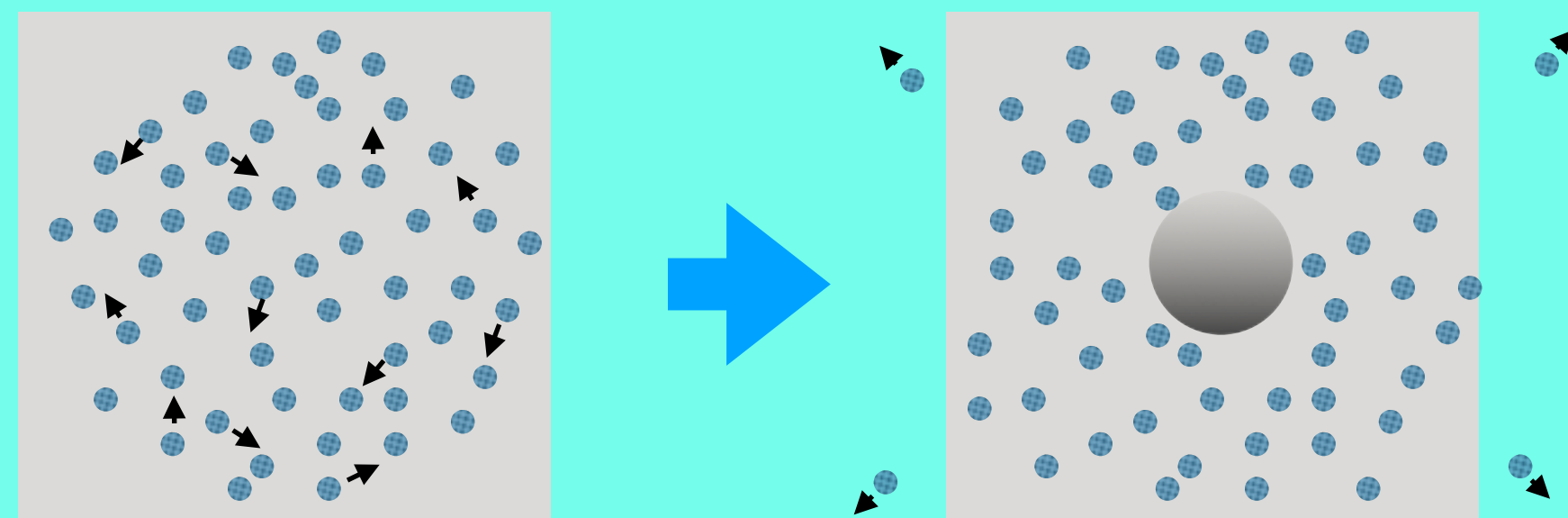
(c) $\tilde{t} = 2000$



For axion stars:
dedicated simulations confirm this picture

For gravitational atoms: so far, ours is the only dedicated analysis

N objects scatter gravitationally, exchange energy

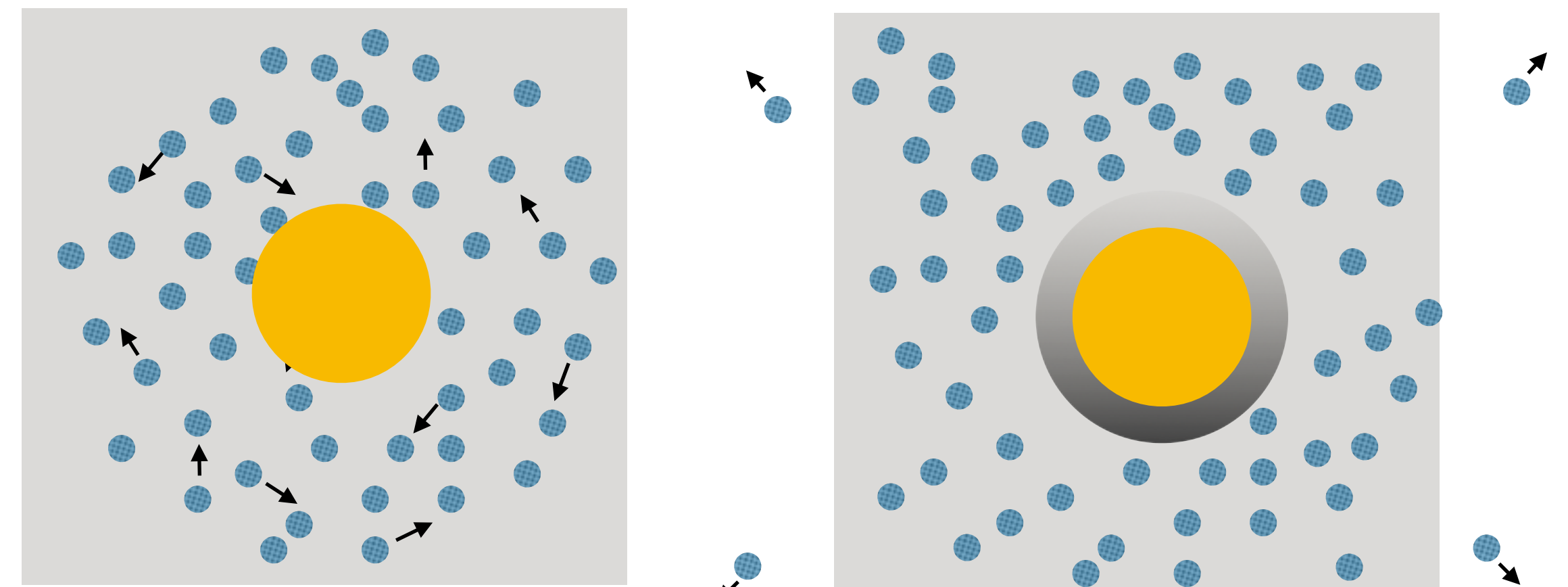


$$\Delta v^2 \simeq 8 N \left(\frac{GM}{R_{\text{gal}} v} \right) \ln N \quad \longrightarrow \quad \frac{\Delta v^2}{v^2} \simeq \frac{8 \ln N}{N}$$

Velocity change per crossing

$$\implies \quad t_{\text{relax}} \simeq \frac{0.1 N}{\ln N} t_{\text{cross}}$$

Binney and Tremaine, “Galactic Dynamics”



Quasiparticle scattering

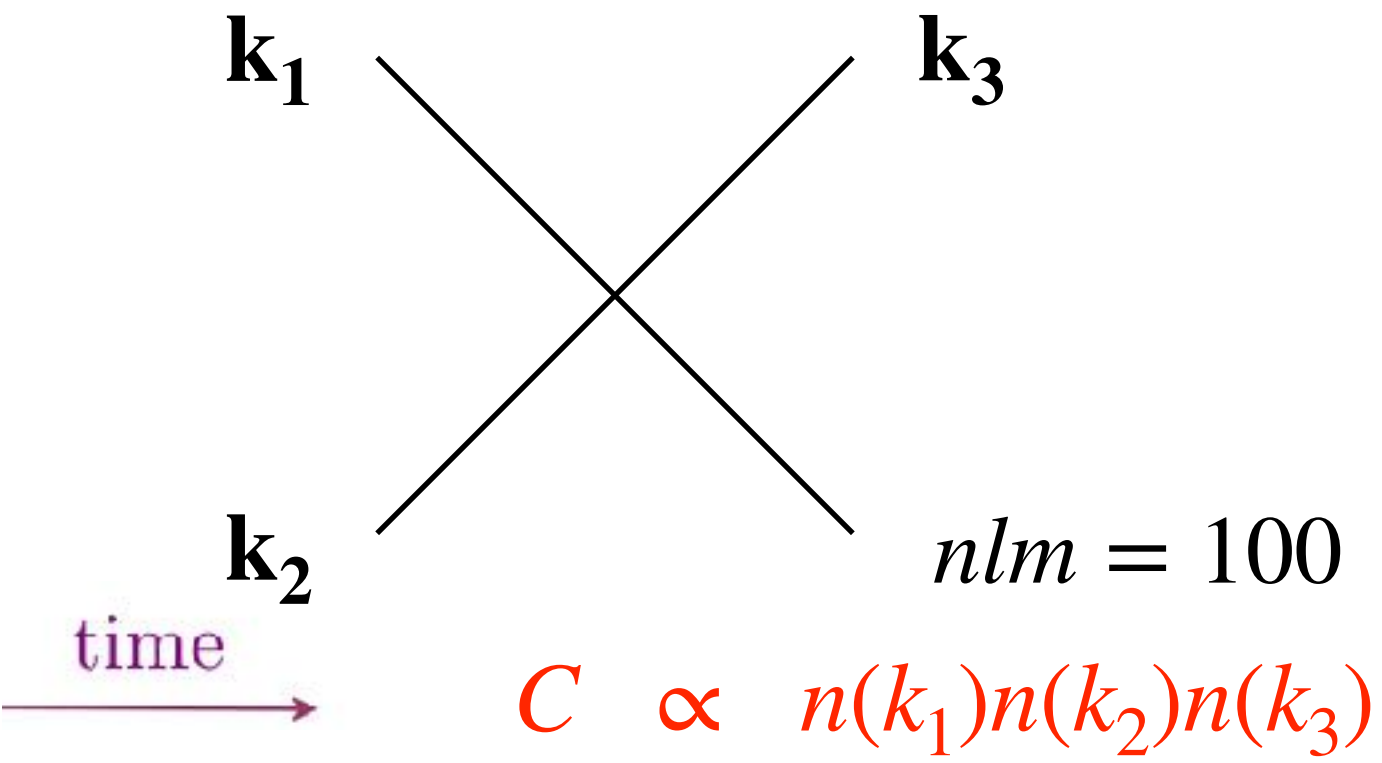
Bound halo formation??

Relaxation to ground state?

Budker, JE, Gorghetto, Jiang, Perez (2306.12477)

Dilute Gravitational Atoms

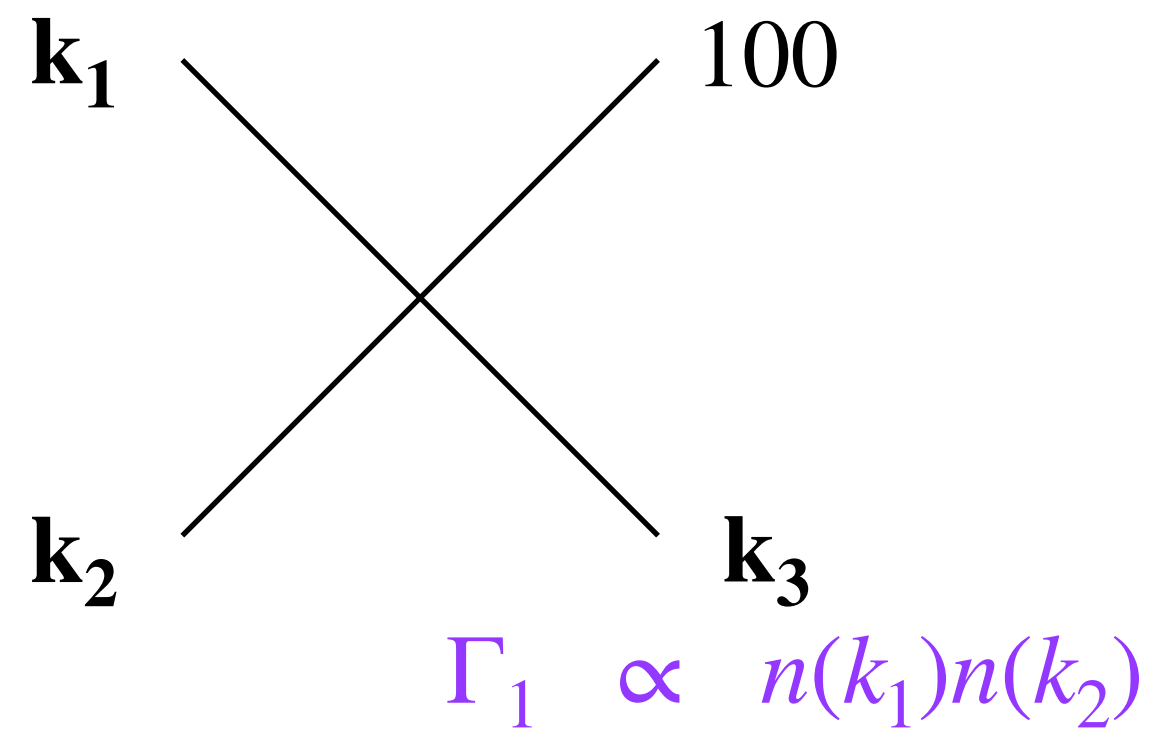
Direct capture



$$\frac{dM_\star}{dt} \simeq C + (\Gamma_1 - 2\Gamma_2)M_\star$$

$\Gamma > 0$: Exponential growth
 $\Gamma < 0$: Saturation
 determines late-time behavior

Stimulated capture

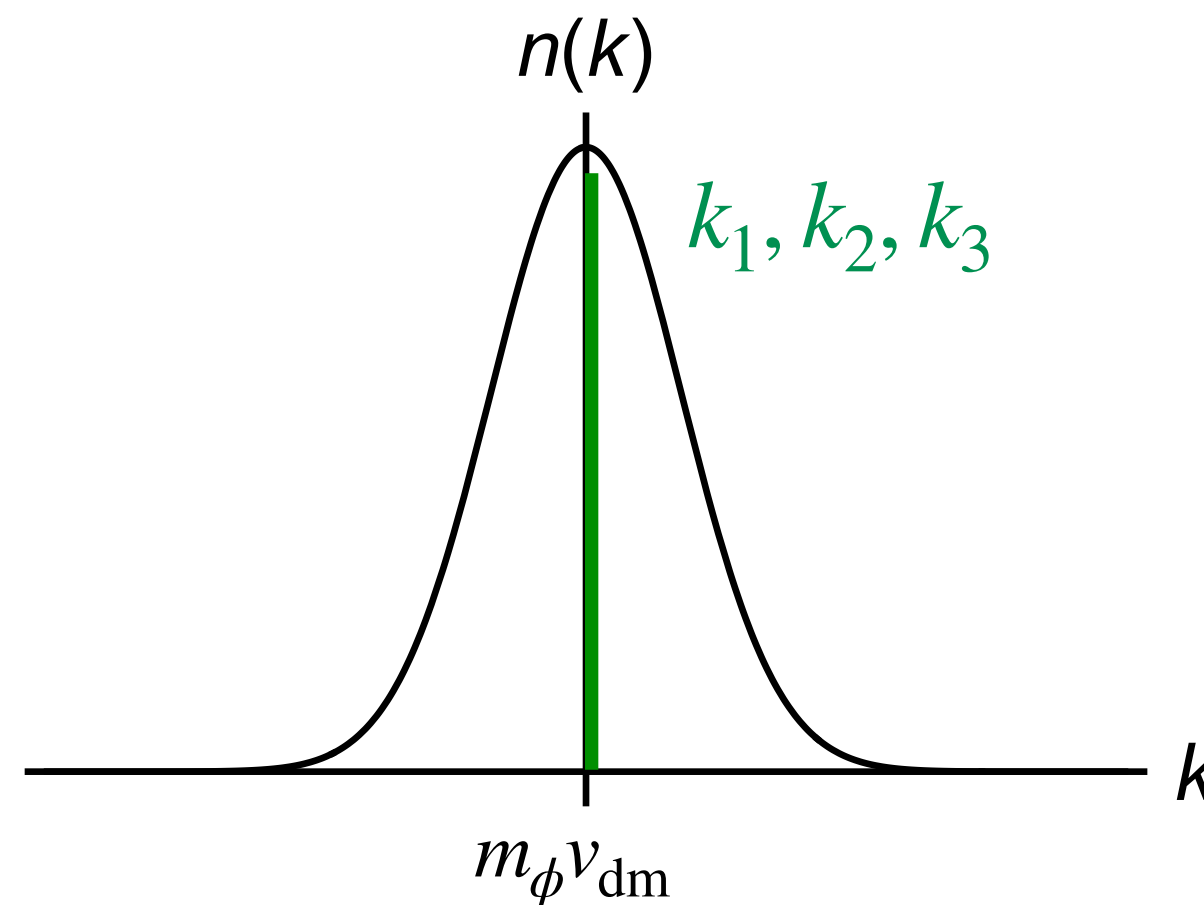
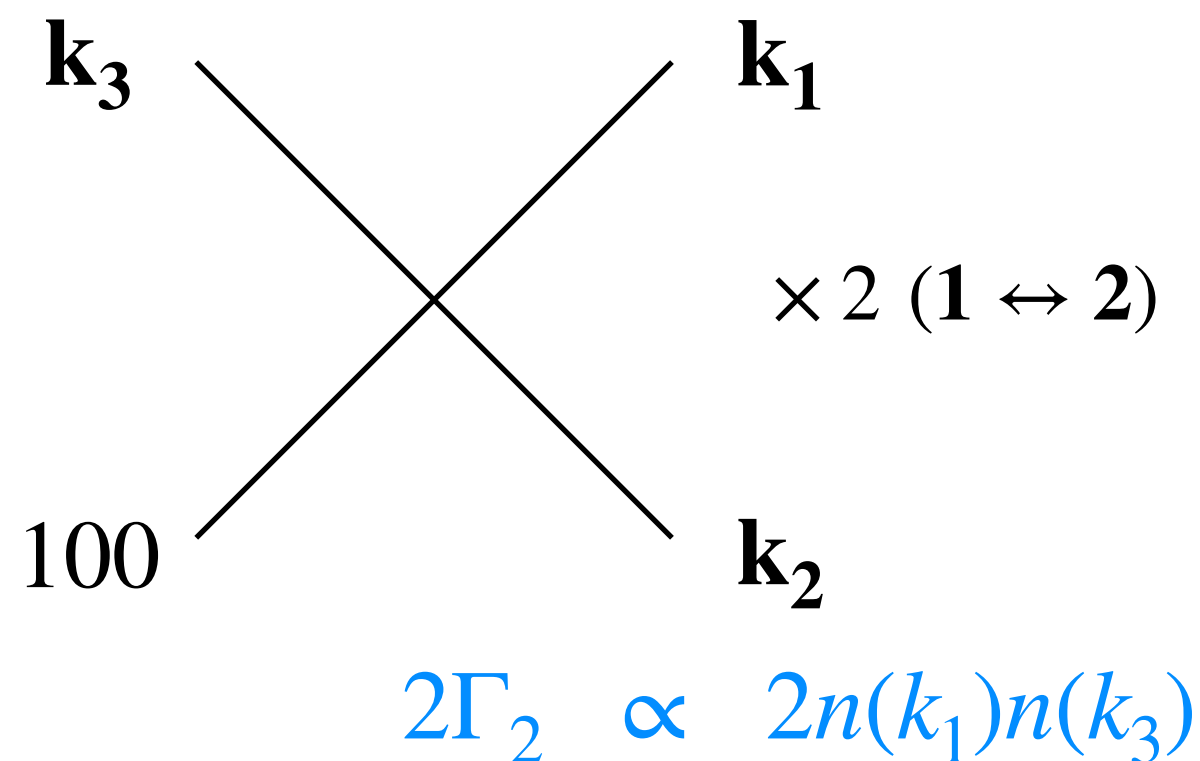


Recall $\xi_{\text{foc}} \equiv \frac{\lambda_{\text{dB}}}{R_\star} = \frac{2\pi\alpha_g}{v_{\text{dm}}} \ll 1$

implies $m_\phi v_{\text{dm}}^2 \gg m_\phi \alpha_g^2$

$k_1 \sim k_2 \sim k_3 \sim m_\phi v_{\text{dm}} \Rightarrow \Gamma_1 \simeq \Gamma_2$

Ionization



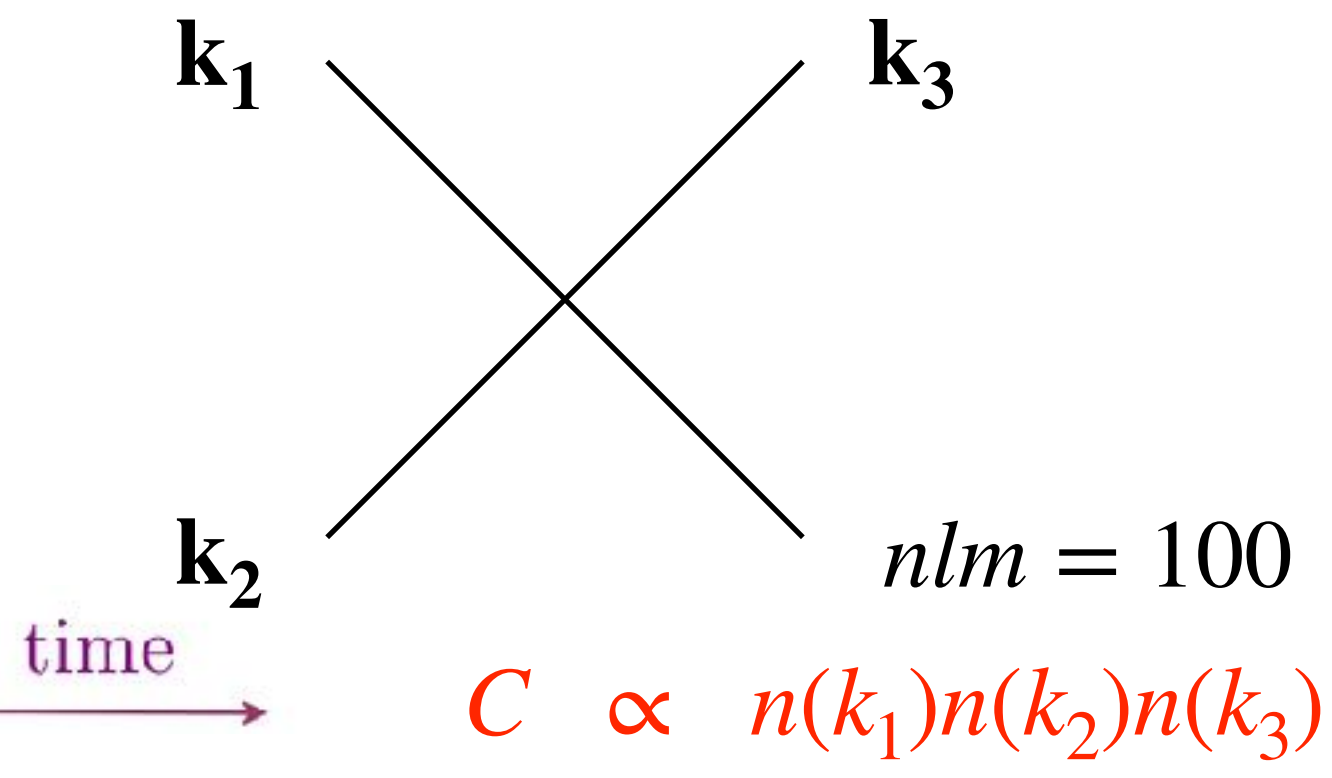
$\Rightarrow M_\star^{\text{eq}} = \frac{C}{\Gamma_2} \simeq 10\rho_{\text{dm}}\lambda_{\text{dB}}^3$

$\Rightarrow \rho \ll \rho_{\text{dm}}$
 "dilute"

Budker, JE, Gorghetto, Jiang, Perez (2306.12477)

Dense Gravitational Atoms

Direct capture



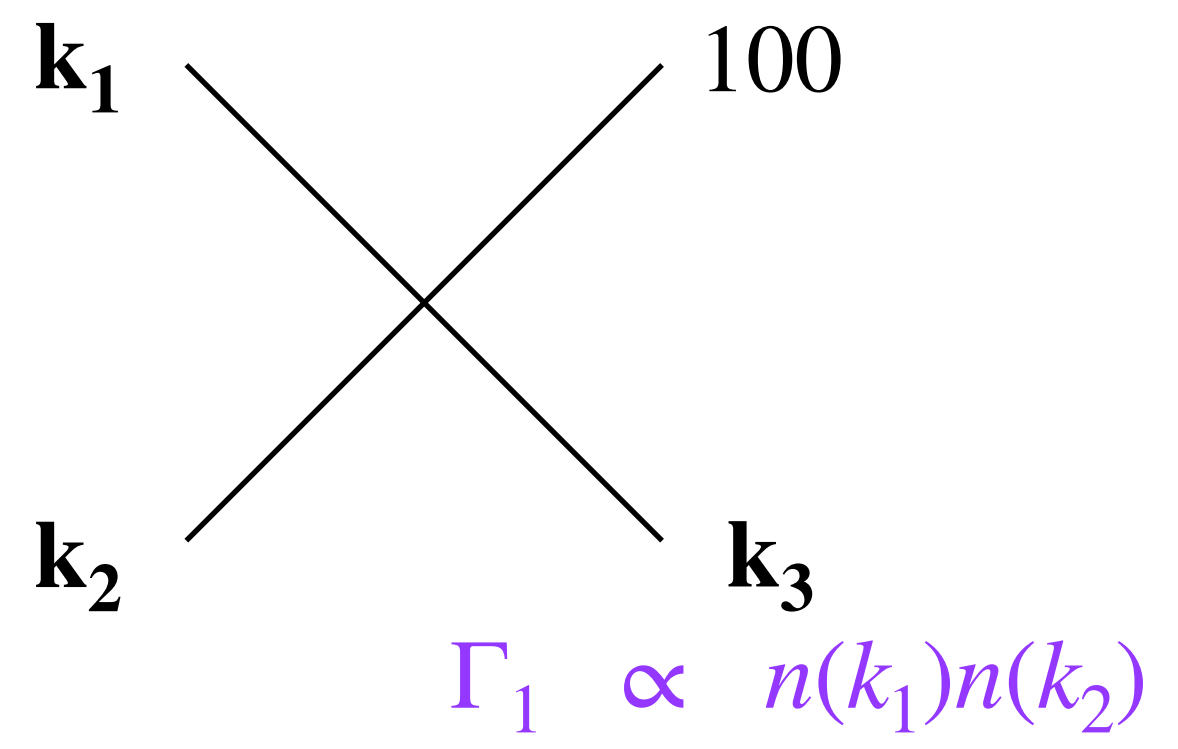
$$\frac{dM_\star}{dt} \simeq C + (\Gamma_1 - 2\Gamma_2)M_\star$$

$\Gamma > 0$: Exponential growth

$\Gamma < 0$: Saturation

determines late-time behavior

Stimulated capture

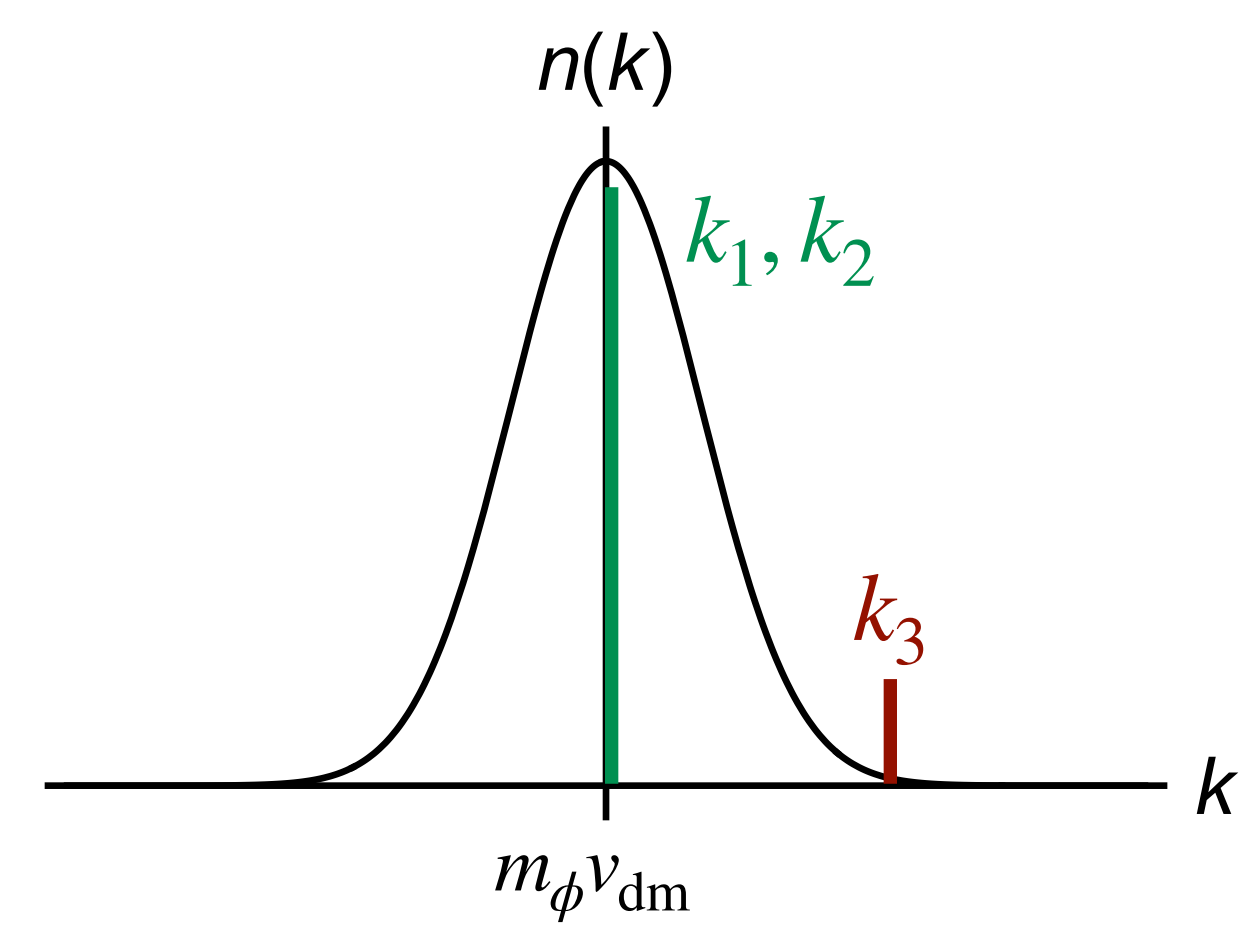
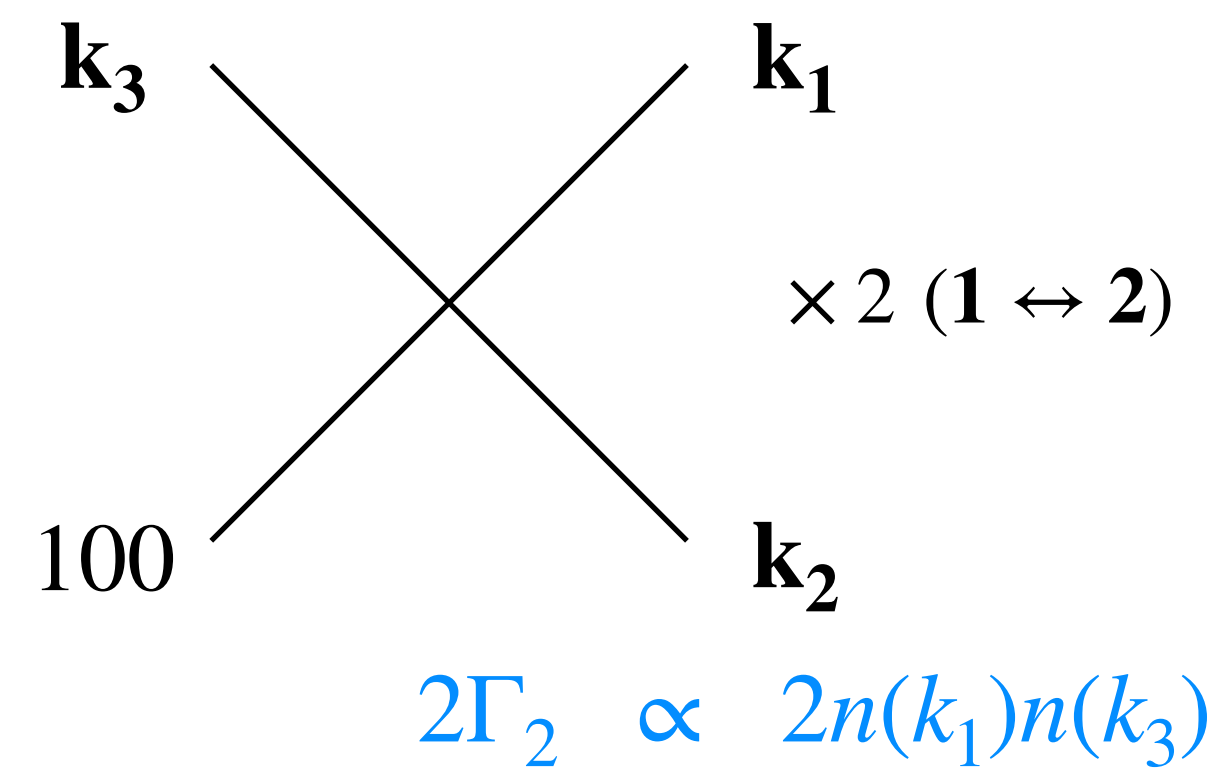


However, $\xi_{\text{foc}} \equiv \frac{\lambda_{\text{dB}}}{R_\star} = \frac{2\pi\alpha_g}{v_{\text{dm}}} \gg 1$

implies instead $m_\phi v_{\text{dm}}^2 \ll m_\phi \alpha_g^2$

$$k_1 \sim k_2 \sim m_\phi v_{\text{dm}} \ll k_3 \Rightarrow \Gamma_1 \gg \Gamma_2$$

Ionization

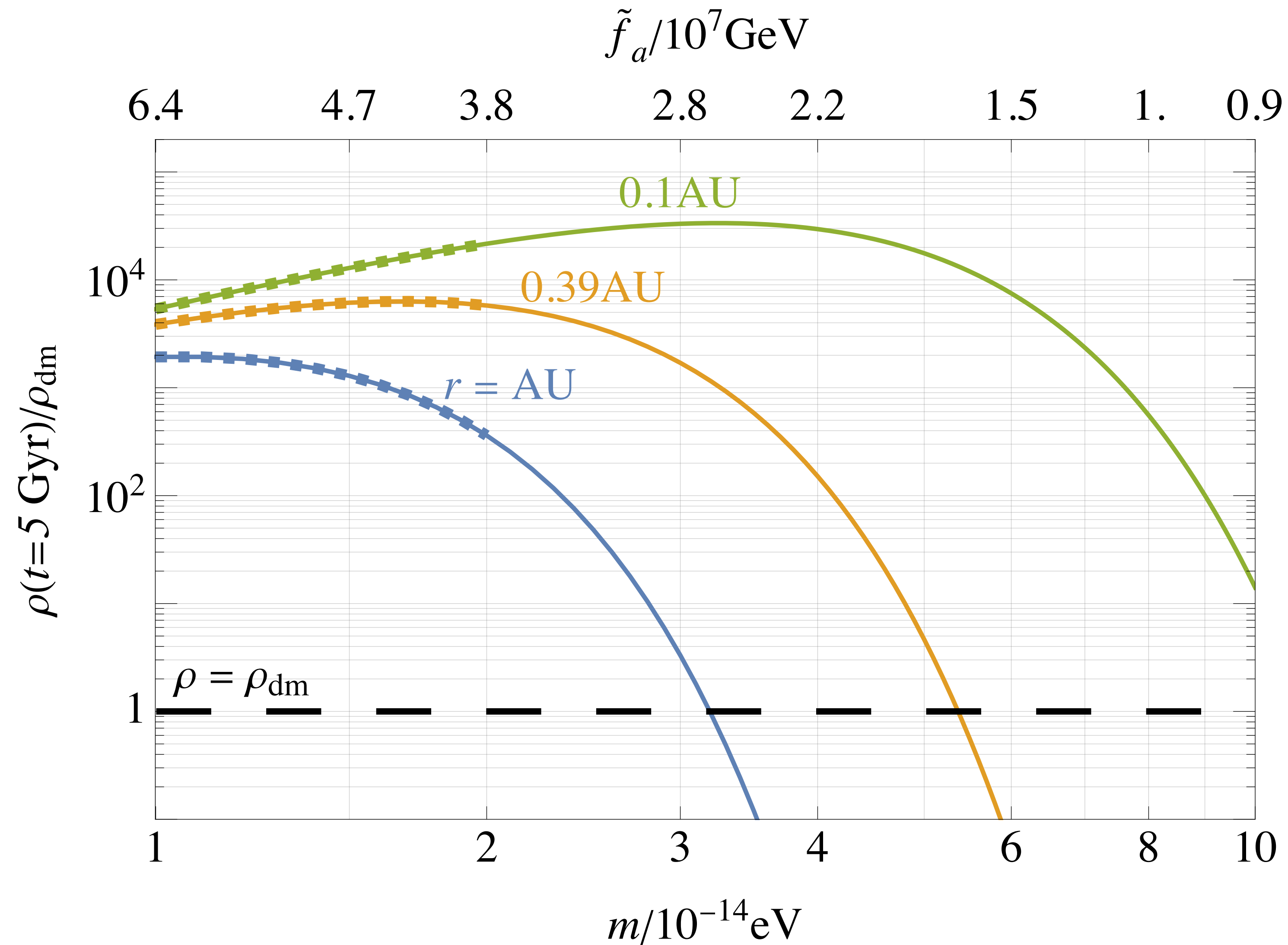


$$\Rightarrow M_\star(t) \simeq 10\rho_{\text{dm}}\lambda_{\text{dB}}^3 (e^{\Gamma_1 t} - 1)$$

$\rho \gg \rho_{\text{dm}}$
"dense"

Very Local Density from Capture

(i.e. in our solar system)



assuming $\lambda < 0$;
 density could be larger for $\lambda > 0$

Budker, JE, Gorghetto, Jiang, Perez (2306.12477)

Constraints from Large-Scale Structure

Evolution of cosmological density perturbations is modified by wavelike structure of ULDM fields

$$\ddot{\delta}_k + 2H\dot{\delta}_k - 4\pi G\bar{\rho} \left[1 - \left(\frac{k}{k_J}\right) \pm \left(\frac{k}{k_\lambda}\right)^2 \right]$$

suppression of structure on 'small' scales due to ultralight mass

$$k_J \simeq 10 \text{ Mpc}^{-1} \left(\frac{a}{a_{\text{eq}}}\right)^{1/4} \left(\frac{m}{10^{-22} \text{ eV}}\right)^{1/2}$$

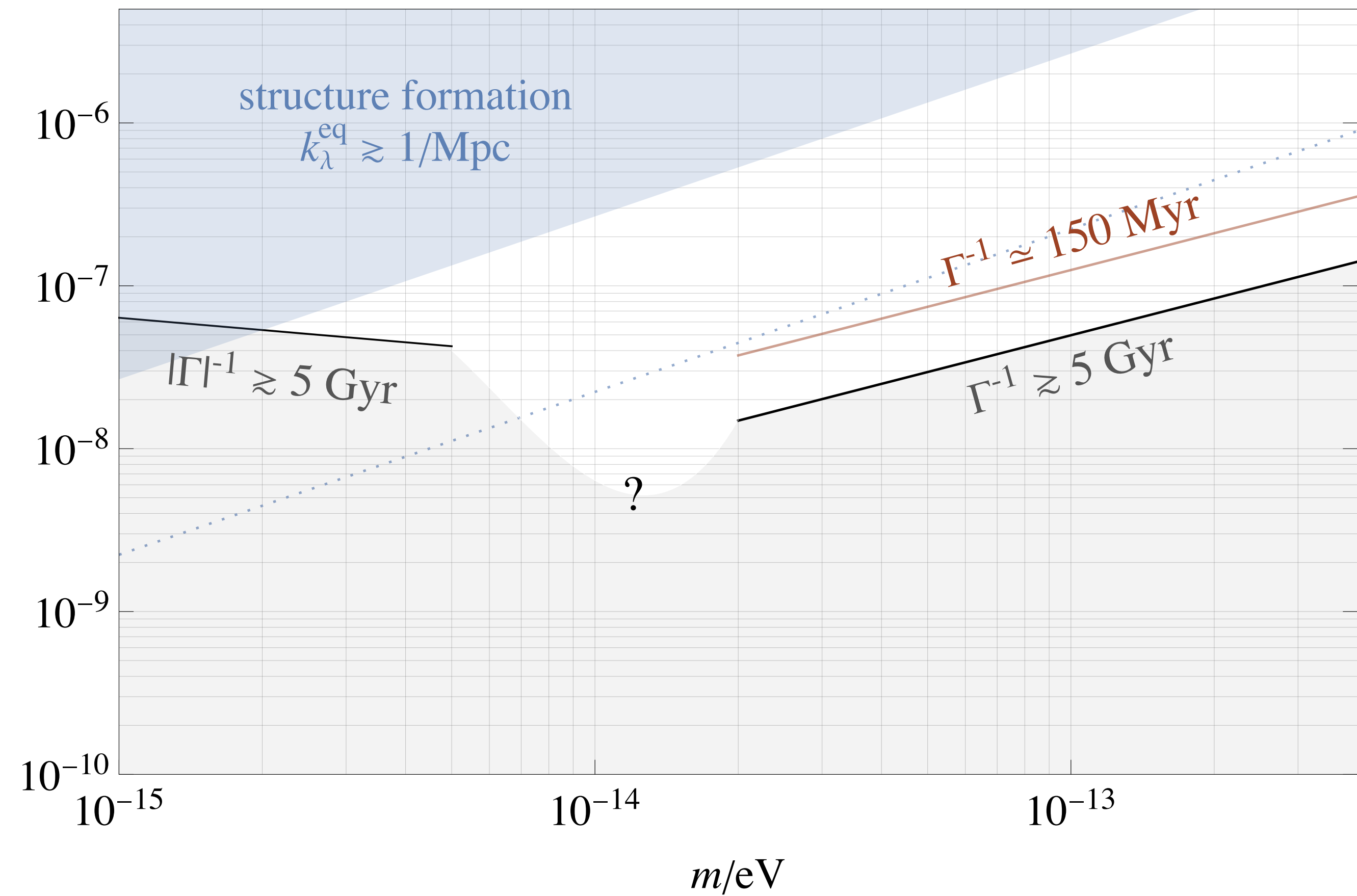
e.g. Lyman- α constraints, see Iršič++ (1703.04683), Rogers, Peiris (2007.12705)

suppression (or enhancement) of structure due to strong self-interactions in the early universe

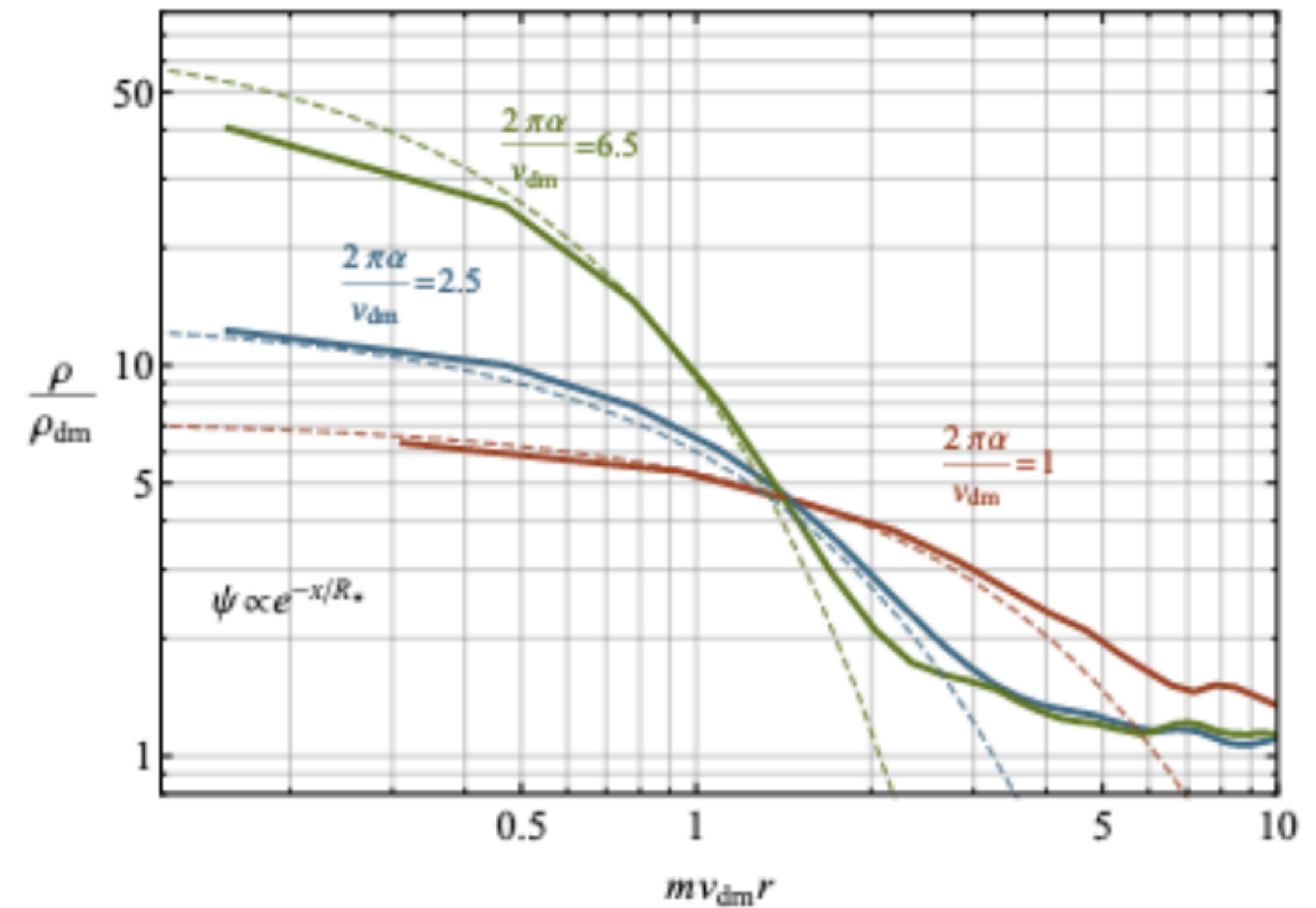
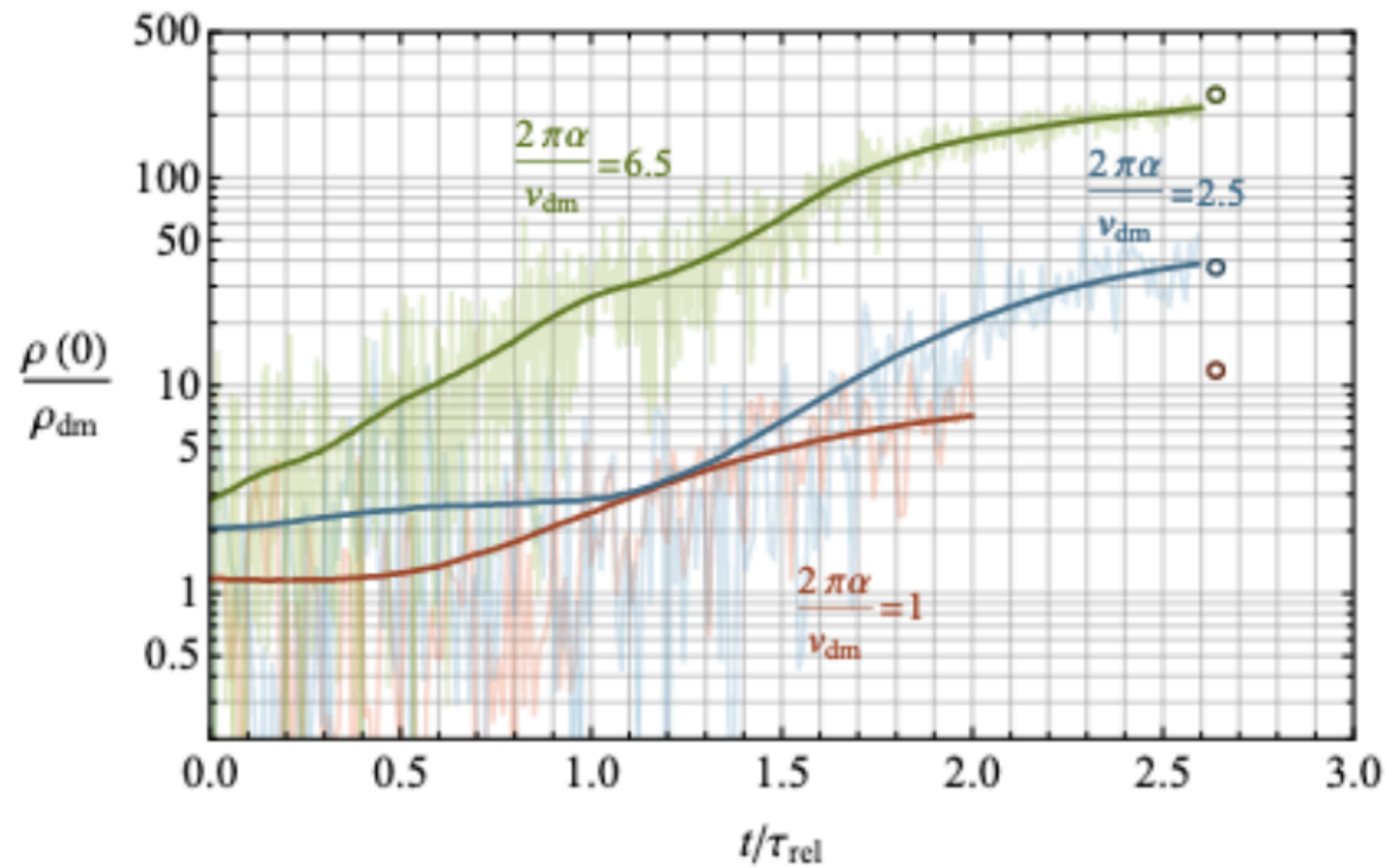
$$k_\lambda \simeq 2.7 \text{ Mpc}^{-1} \left(\frac{a}{a_{\text{eq}}}\right) \left(\frac{m}{10^{-14} \text{ eV}}\right) \left(\frac{f_a}{10^7 \text{ GeV}}\right)$$

see e.g. Arvanitaki, Huang, Van Tilburg (1405.2925) Fan (1603.06580) Cembranos++ (1805.08112)

$\frac{\text{GeV}}{f_a}$



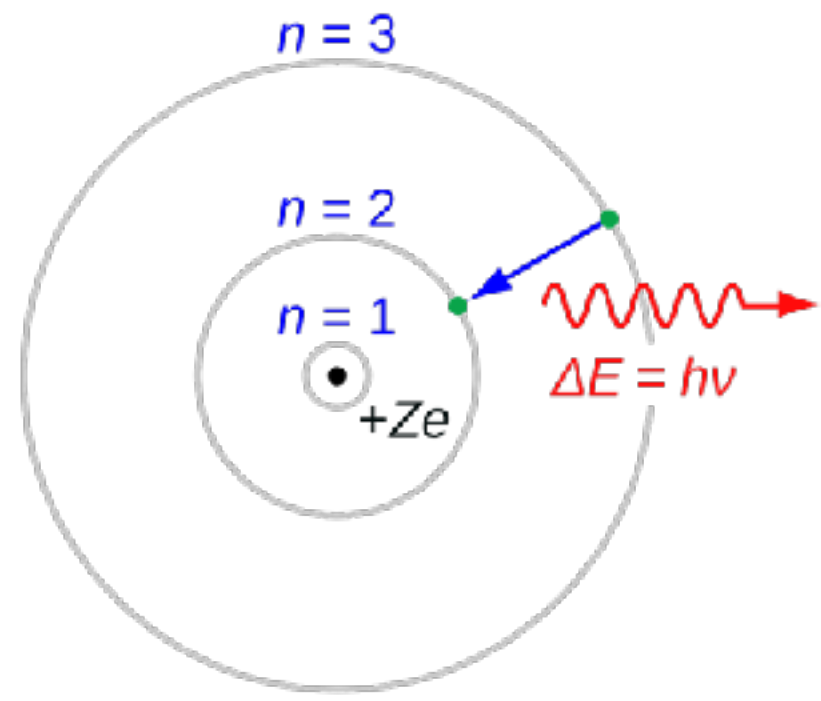
Simulations



Banerjee, Budker, JE, Kim, Perez (1902.08212)

Probes on Earth

ULDM-photon interaction (example):



$$\mathcal{L} \supset \frac{d_\alpha}{4\tilde{M}_P} \phi F^{\mu\nu} F_{\mu\nu}$$

Oscillation of $\phi = \phi(t)$ induces oscillation of fundamental constants of nature

$$\alpha(t) = \alpha_0 \left(1 - d_\alpha \frac{\phi(t)}{\tilde{M}_P} \right)$$

$$\Rightarrow \frac{\delta\alpha}{\alpha_0} \simeq \frac{d_\alpha \sqrt{2\rho}}{m_\phi \tilde{M}_P} \simeq 10^{-15} d_\alpha \left(\frac{10^{-15} \text{ eV}}{m_\phi} \right) \sqrt{\frac{\rho}{\rho_{\text{local}}}}$$

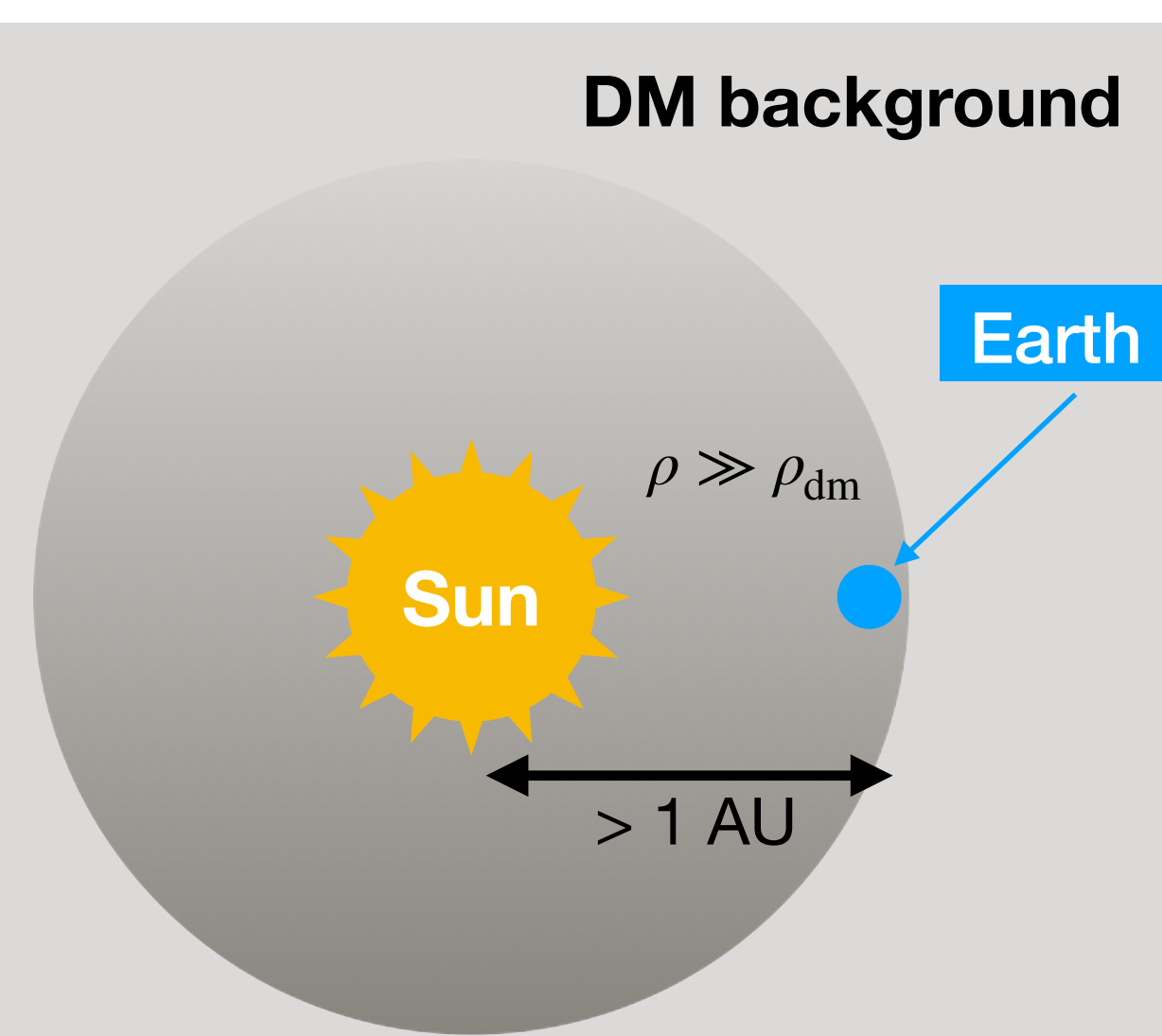
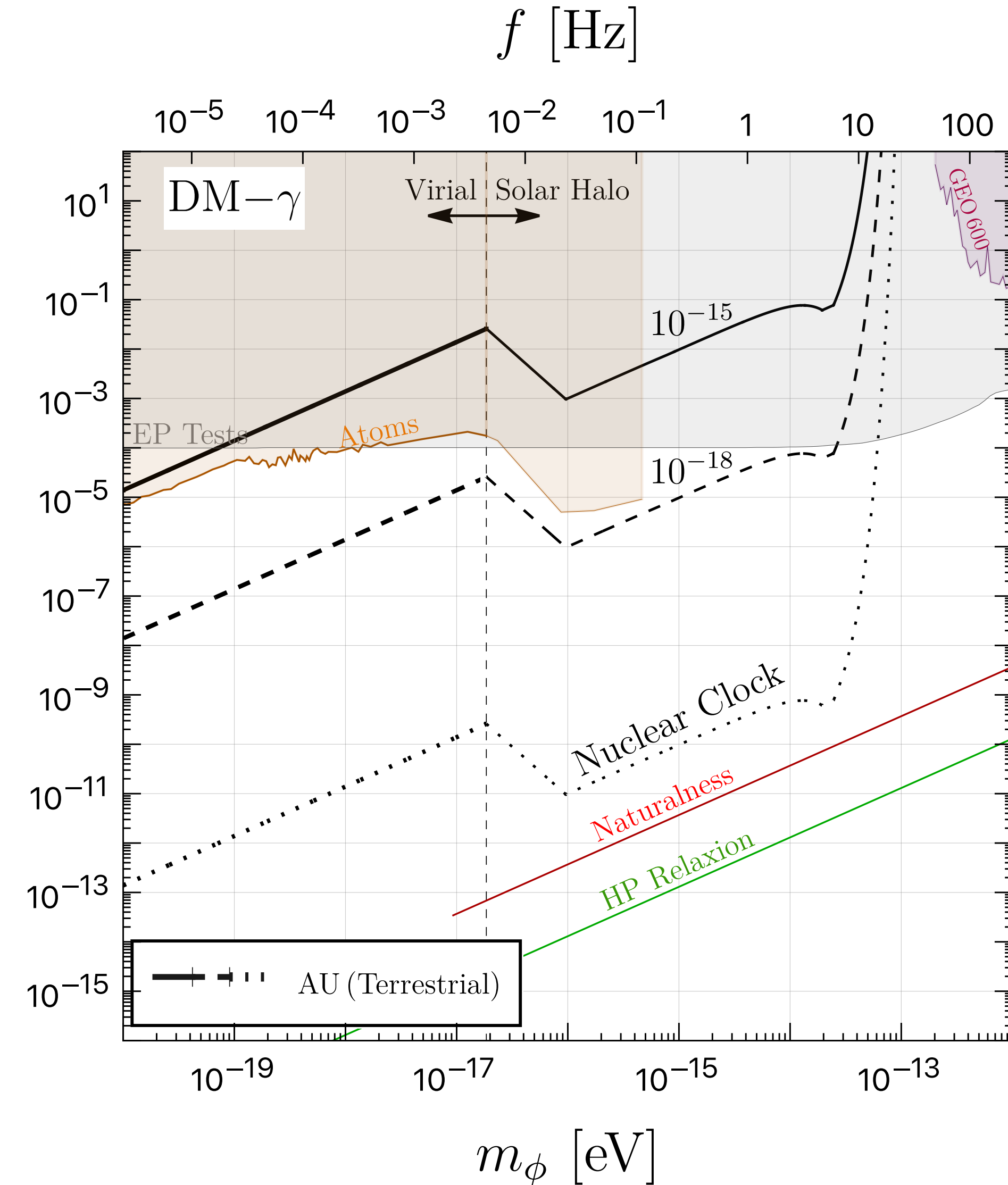
w/ oscillation frequency $\omega_\phi \simeq m_\phi \simeq \text{few Hz} \left(\frac{m_\phi}{10^{-15} \text{ eV}} \right)$

Modern optical clocks have achieved

$$\frac{\delta\alpha}{\alpha_0} \lesssim 10^{-18}$$

Future nuclear clock projected to reach

$$\frac{\delta\alpha}{\alpha_0} \lesssim 10^{-23}$$



Going to Space

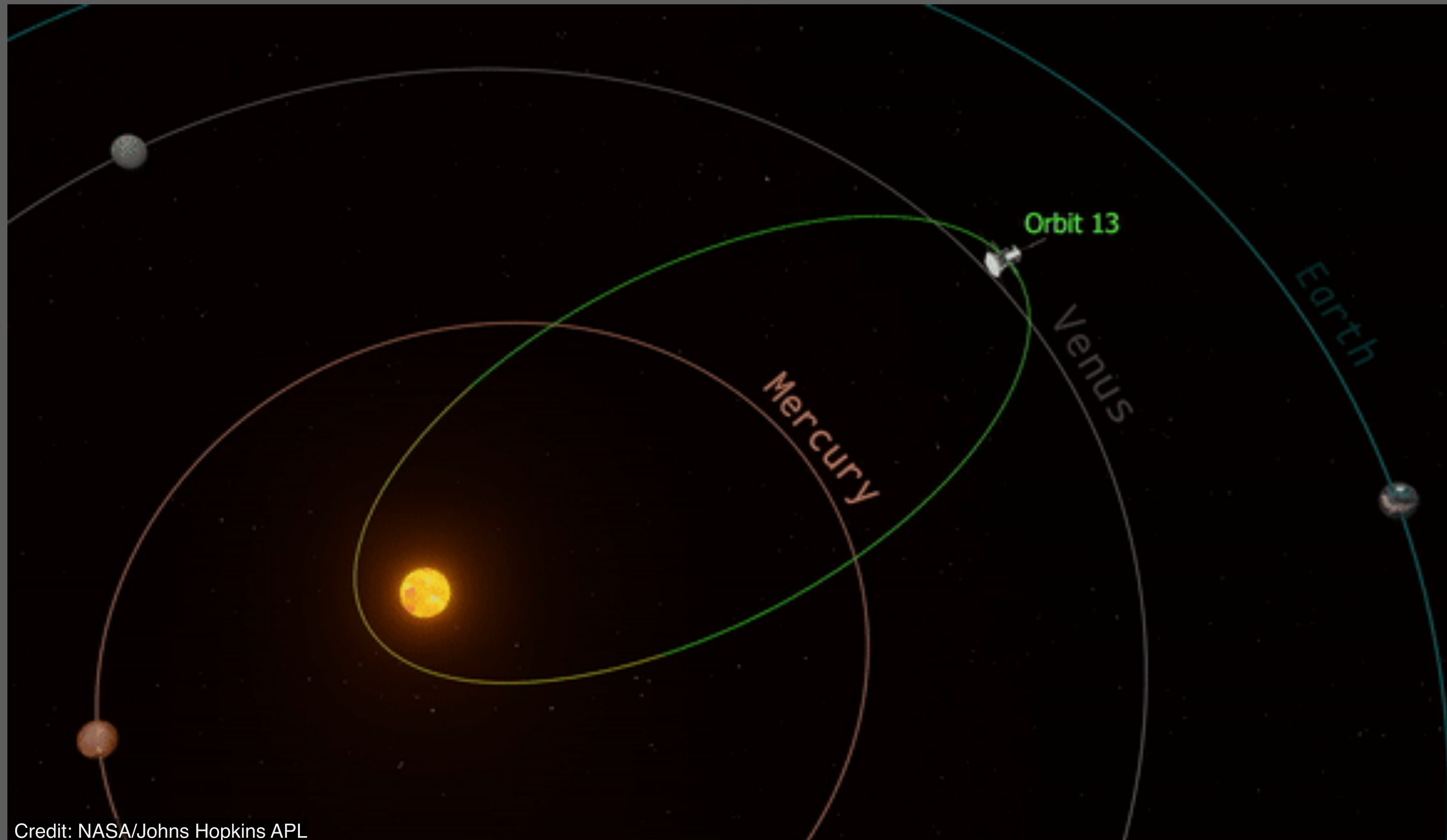
Parker Solar Probe (NASA)

<https://blogs.nasa.gov/parkersolarprobe/>

Launch: **Aug 12, 2018**

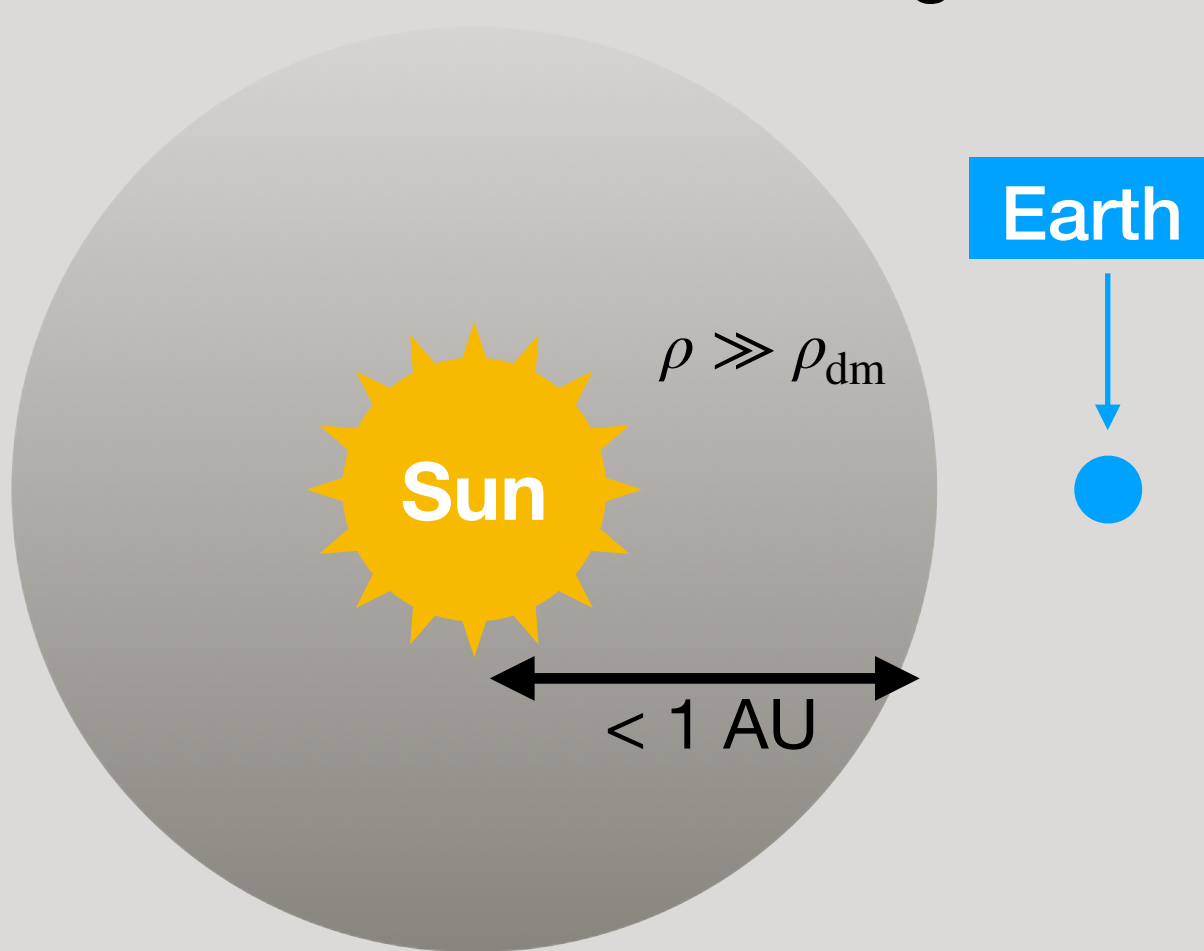
Most recent perihelion:

- **March 17, 2023 (15 of 24 planned)**
- **Approach: 0.057 au from Sun**
- **Temperature:**
2500°F (1370°C) heat shield,
85°F (30°C) interior



Credit: NASA/Johns Hopkins APL

DM background



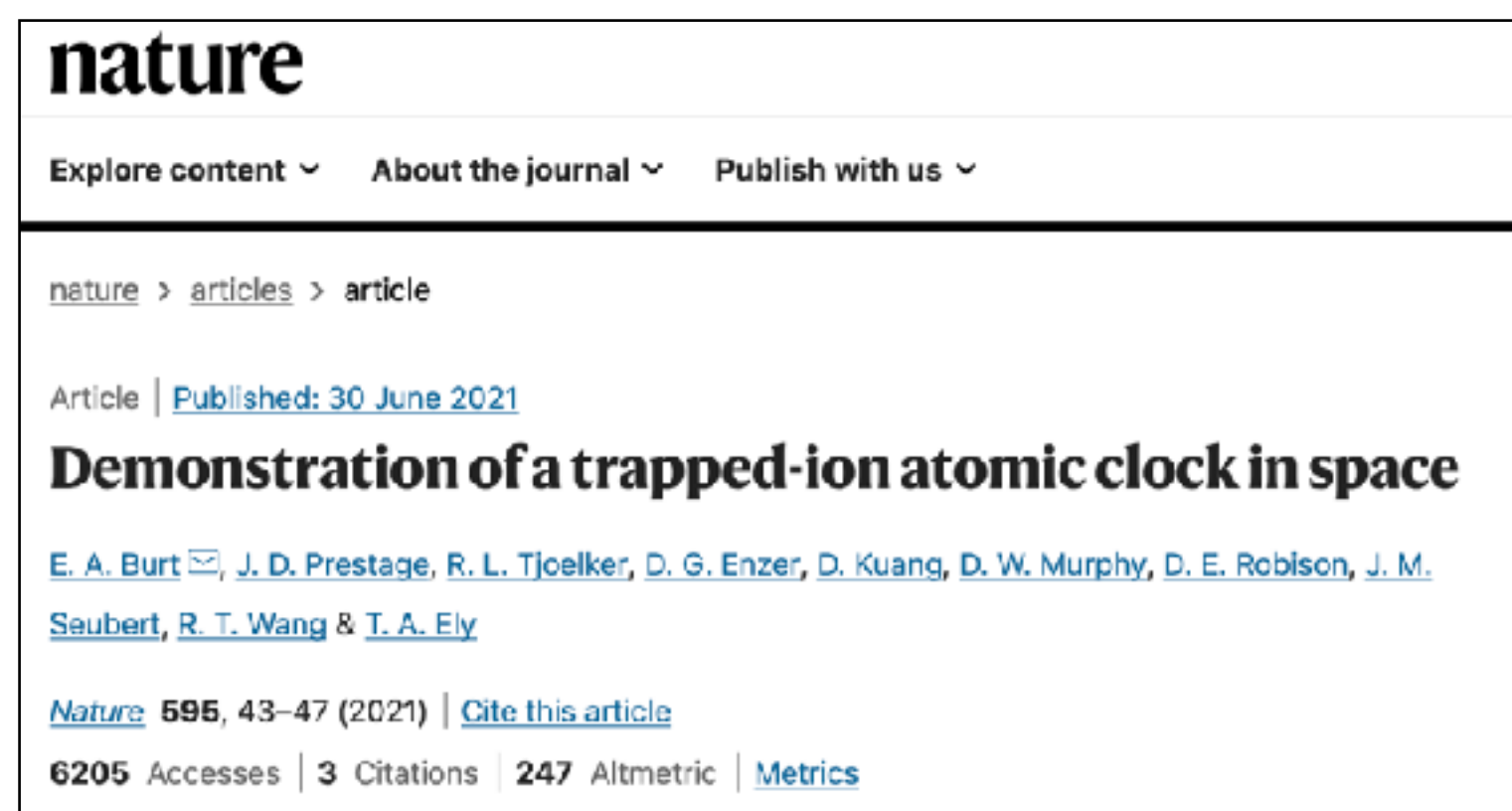
Can we send an atomic clock to the Sun?

(Should we?)

Quantum Clocks in Space

International community rapidly developing technologies to put atomic clocks in space

NASA Deep Space Atomic Clock (DSAC) 🇺🇸



Demonstrated stability in space $\sim 10^{-14}$

DSAC-2 may visit Venus, launch ~2028!

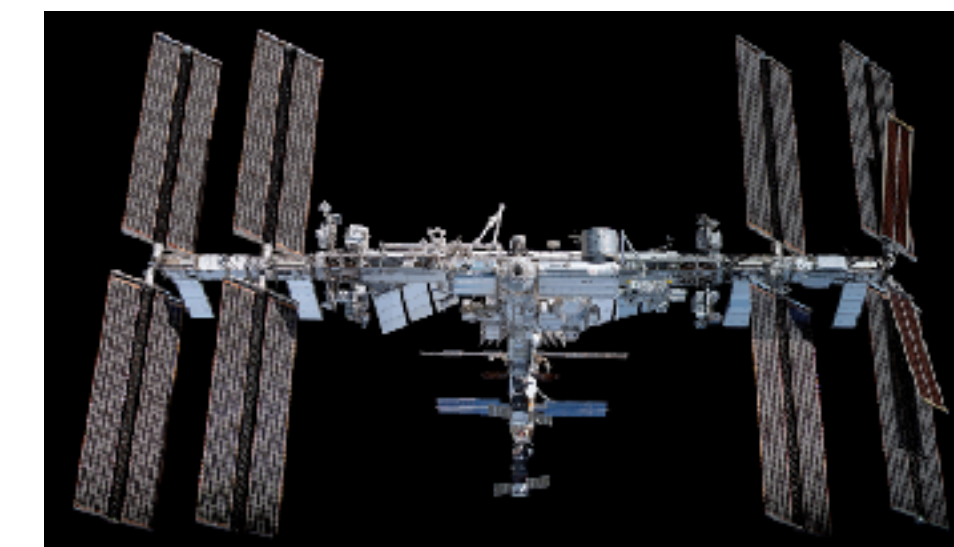
Cold Atom Clock Experiment in Space (CACES) 🇨🇳



Demonstrated stability in space $\sim 10^{-14}$

Goal for Tiangong space station $\sim 10^{-18}$

Atomic Clock Ensemble in Space (ACES) 🇪🇺



Launch date: 2025 (?)

Target stability $\sim 10^{-16}$

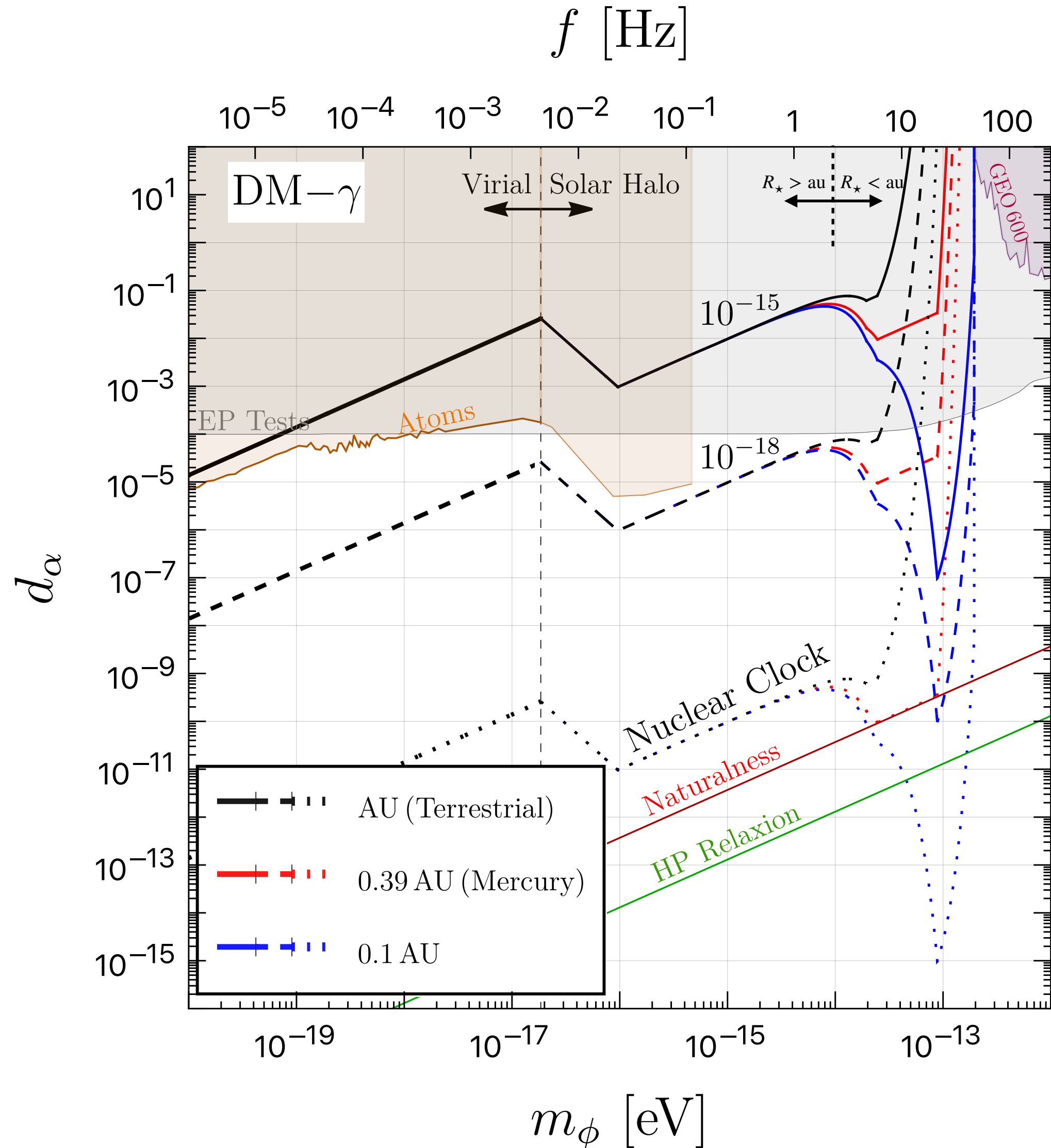
Motivations:

- Better time-keeping on ISS
- Improved GPS
- Comparison with ground clocks

- Spacecraft navigation
- Communication w/ moon / Mars / (?)
- **ULDM detection??**

Tsai, **JE**, Safronova
(2112.07674)

Probes in Space



Lots of work still needed to make this possible!

On the space probe side:

- Temperature variations?
 - $(-40 \text{ to } 40)^\circ\text{C}$ variation behind heat shield
 - Maybe cold is more dangerous than hot!
- Magnetic field variations?
 - $\pm 100 \text{ nT}$ variation on \sim sec timescales

discussions with
Parker Solar
Probe
scientists

On the clock side:

- Need clocks which are
 - Portable + Automated
 - Lightweight
 - Optimized for 'high frequency' run
- Ideally, nuclear + optical clocks to probe many couplings

discussions with
quantum
sensor / atomic
or nuclear clock
experts

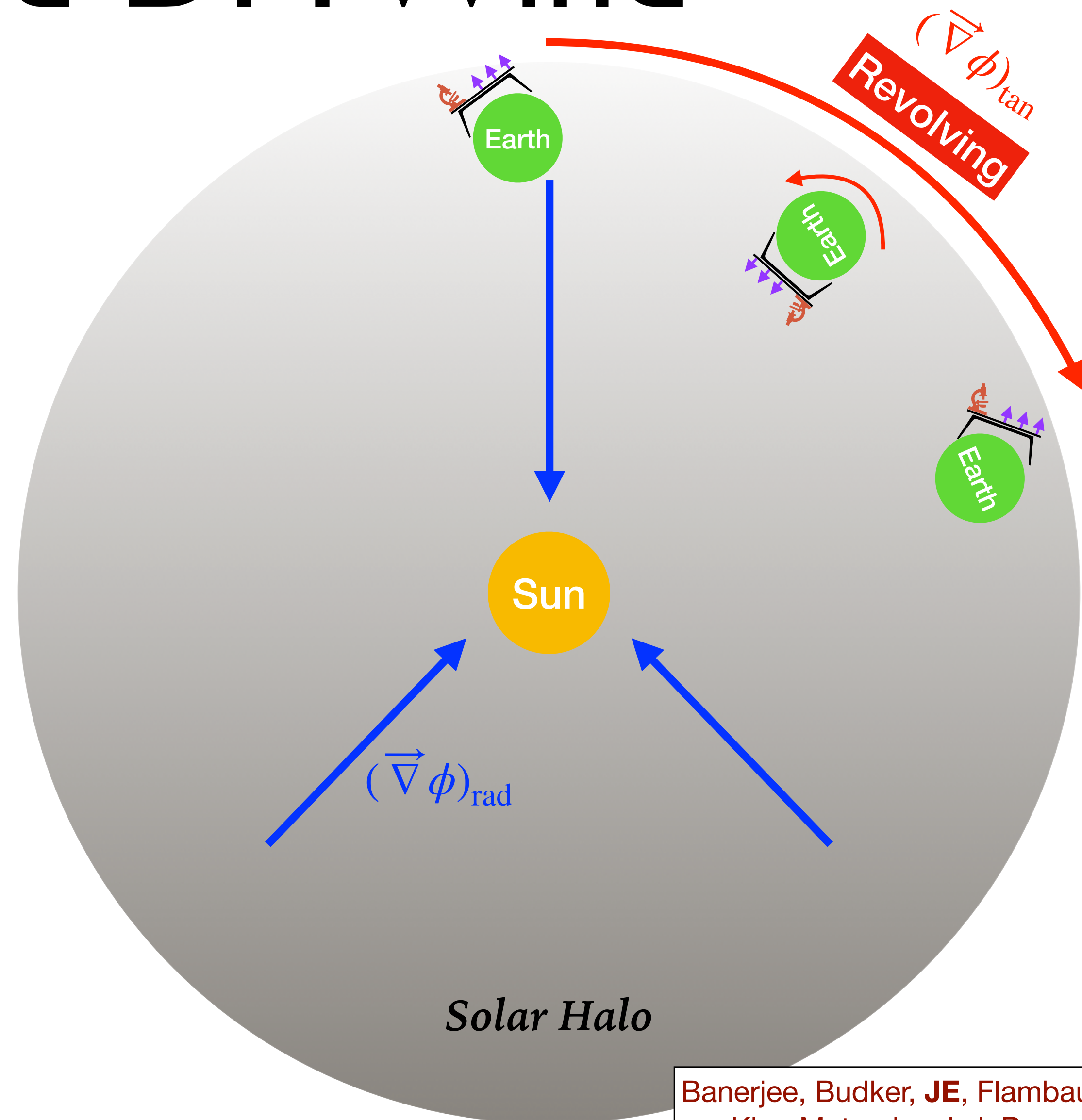
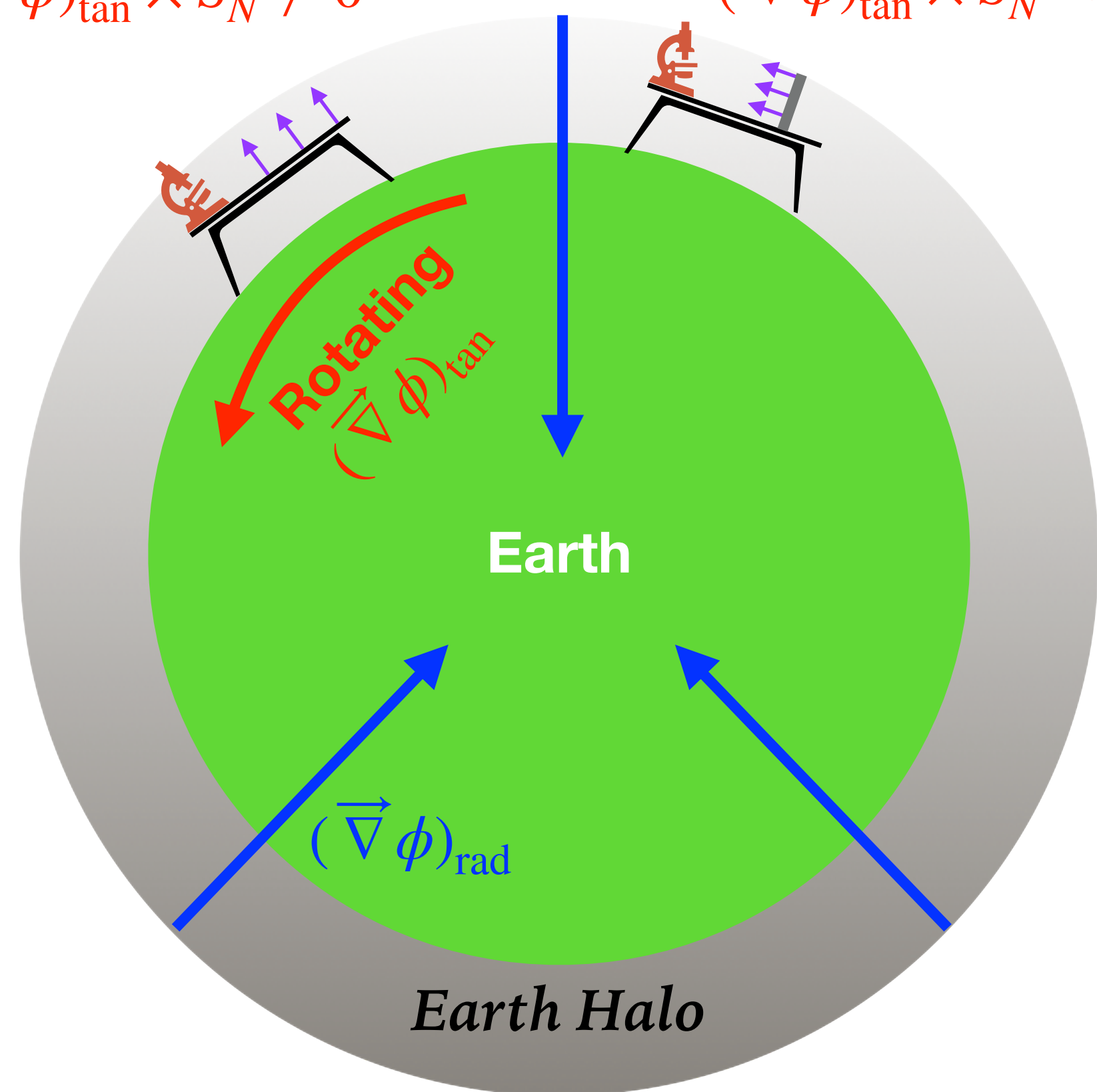
Orientation and DM Wind

$$(\vec{\nabla} \phi)_{\text{rad}} \times \vec{S}_N \rightarrow 0$$

$$(\vec{\nabla} \phi)_{\text{tan}} \times \vec{S}_N \neq 0$$

$$(\vec{\nabla} \phi)_{\text{rad}} \times \vec{S}_N \neq 0$$

$$(\vec{\nabla} \phi)_{\text{tan}} \times \vec{S}_N \rightarrow 0$$

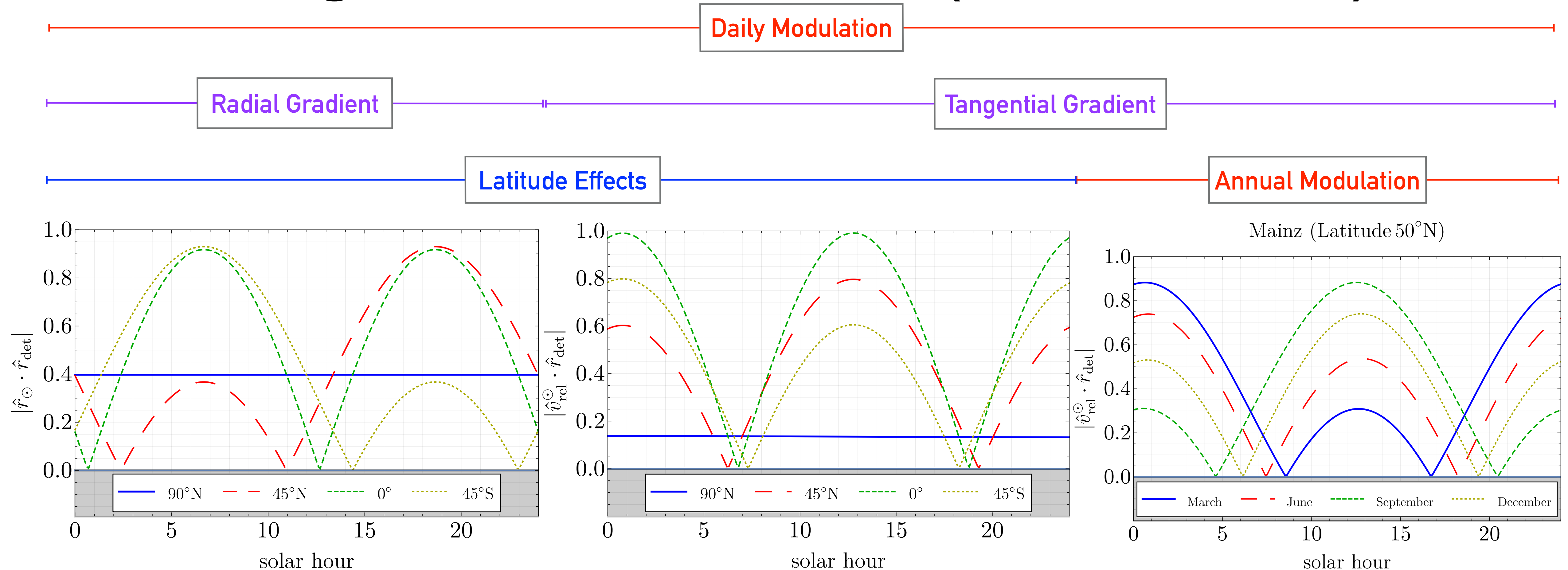


Signal depends both on detector orientation and latitude!

Daily and annual modulation of the signal!

Banerjee, Budker, **JE**, Flambaum, Kim, Matsedonskyi, Perez (1912.04295)

Signal Modulation (Solar Halo)



- Upshot: Sideband analysis in existing axion experiments can distinguish virialized ULDM from bound axion halos in our solar system
- Also motivates network searches (see e.g. [GNOME 2305.01785](#))

Banerjee, Budker, **JE**, Flambaum,
Kim, Matsedonskyi, Perez
(1912.04295)

Axion Stars (ask me later 😊)

