

Learning about dark matter from binary black hole mergers

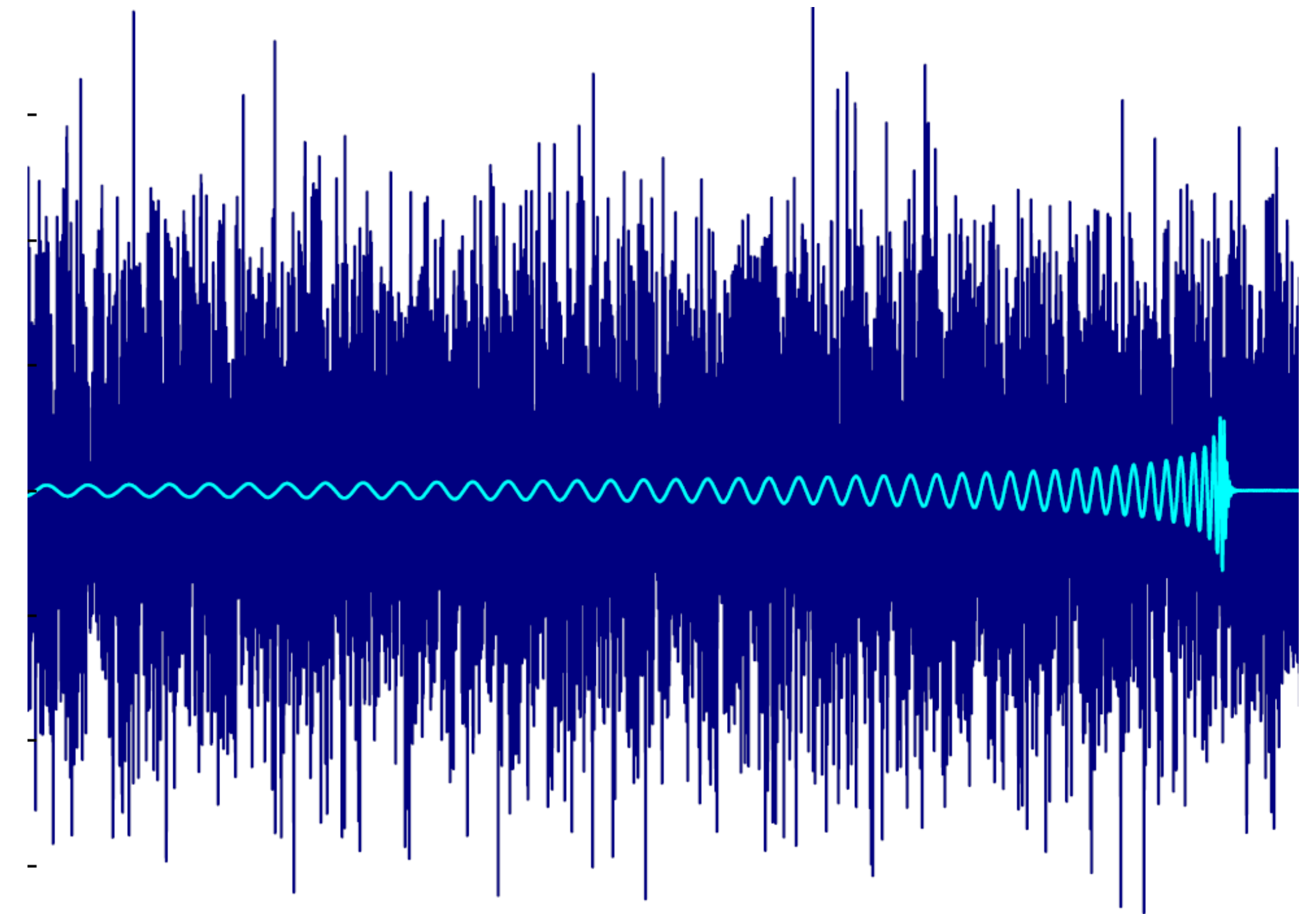
Philippa (Pippa) Cole, University of Amsterdam

**with Gianfranco Bertone, Adam Coogan, Daniele Gaggero, Bradley Kavanagh, Theophanes Karydas,
Thomas Spieksma and Giovanni Maria Tomaselli**

Based on [arXiv:2207.07576](https://arxiv.org/abs/2207.07576) and [arXiv:2211.01362](https://arxiv.org/abs/2211.01362)

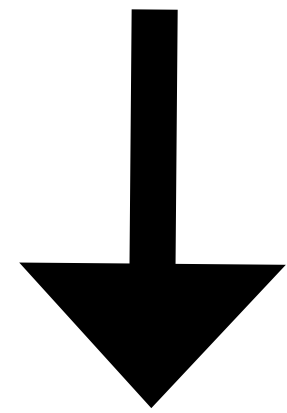
Vacuum or non-vacuum that is the question

- So far, all LIGO/Virgo/KAGRA binary black hole mergers have been detected and measured assuming that they occurred in vacuum
- OK for short duration signals, but looking towards future interferometers, long duration signals may be affected by their environment



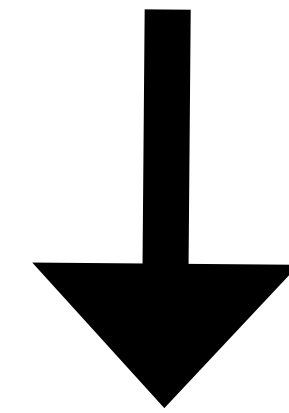
- Environmental effects can cause inspiral to either speed up or slow down with respect to vacuum case
- A dephasing to accumulate, which alters the gravitational waveform from the binary's inspiral

$$\dot{r} = \dot{r}_{\text{GW}} + \dot{r}_{\text{env}}$$

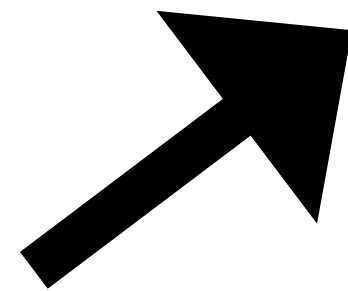


$$f(t) = \frac{1}{\pi} \sqrt{\frac{GM}{r(t)^3}}$$

$$\Phi(f) = \int_f^{f_{\text{ISCO}}} \frac{dt}{df'} f' df'$$



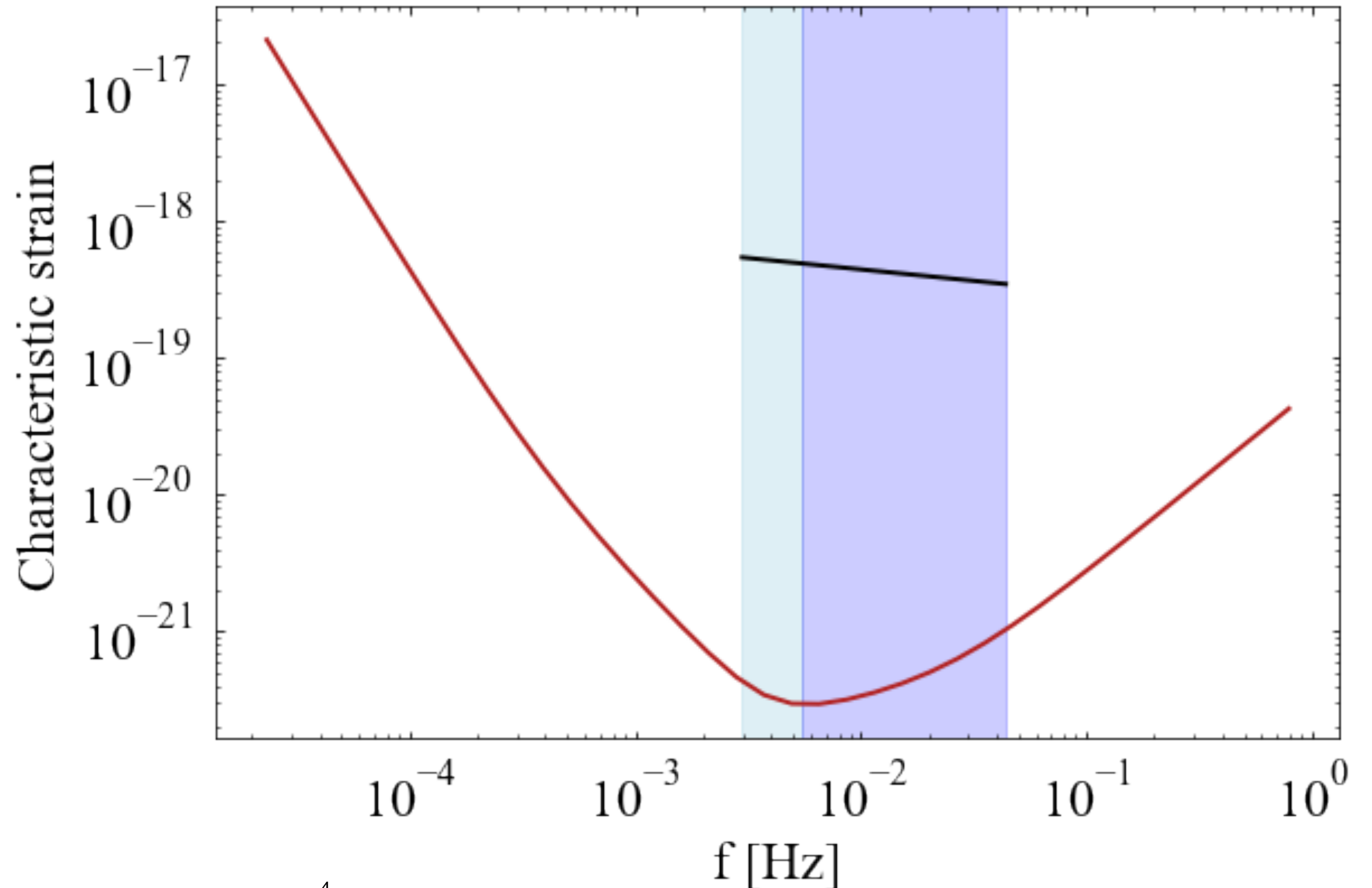
$$h_0(f) = \frac{1}{2} \frac{4\pi^{2/3} G_N^{5/3} \mathcal{M}^{5/3} f^{2/3}}{c^4} \sqrt{\frac{2\pi}{\ddot{\Phi}}}$$



Need to observe many cycles

$$m_1 = 10^5 M_\odot, \quad m_2 = 10 M_\odot$$

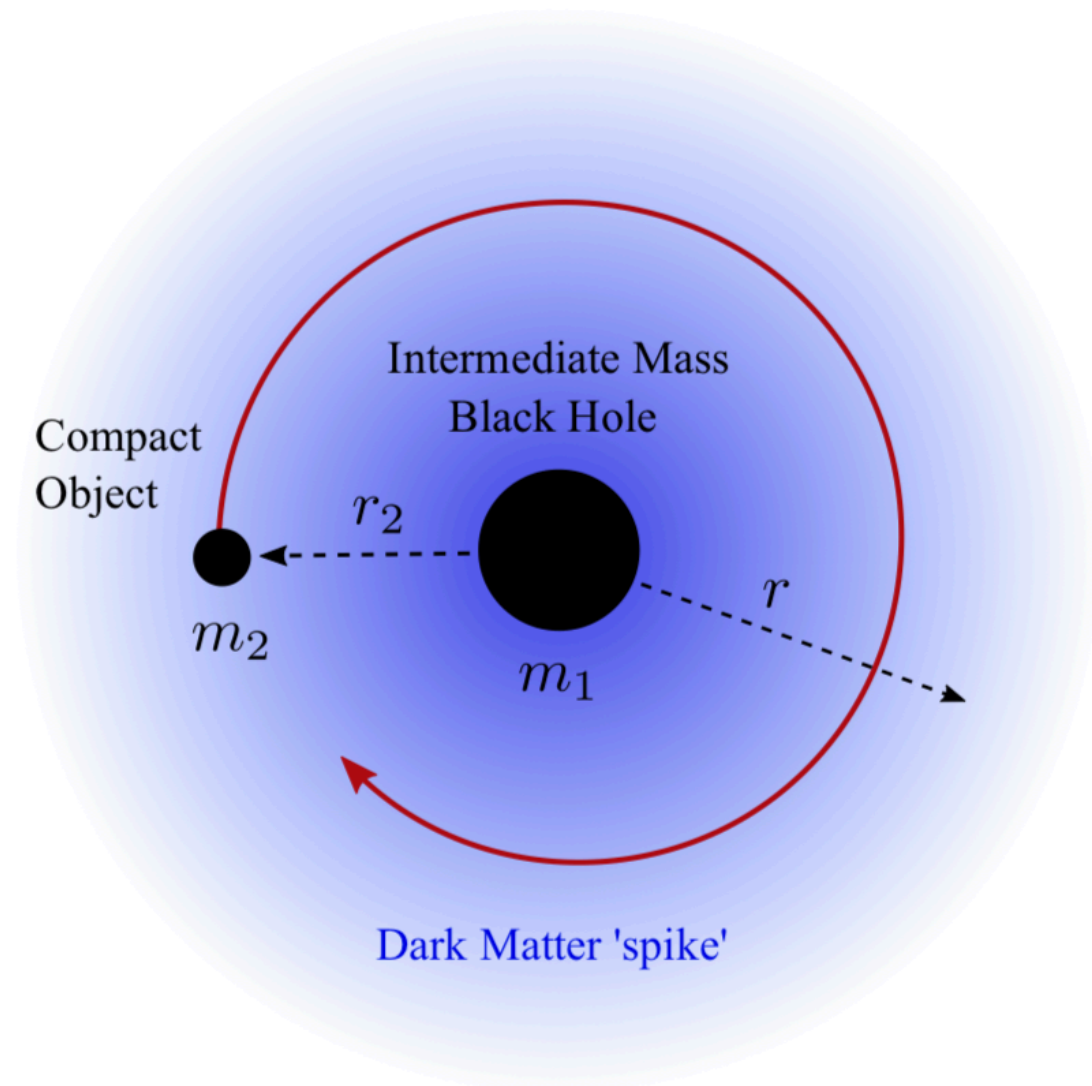
- dephasing accumulates over thousands or millions of cycles
- small mass ratio $q = \frac{m_2}{m_1} < 10^{-2.5}$ so that environment survives
- systems possible sources for LISA and Einstein Telescope/Cosmic Explorer



Why should we care about environmental effects?

- If we can measure the parameters of the environment via the dephasing in the waveform, chance to learn about the environment
- If we search the data with the wrong ‘template’ we might miss the signal
- If we do parameter estimation with the ‘wrong’ parameters, we might come up with biased results

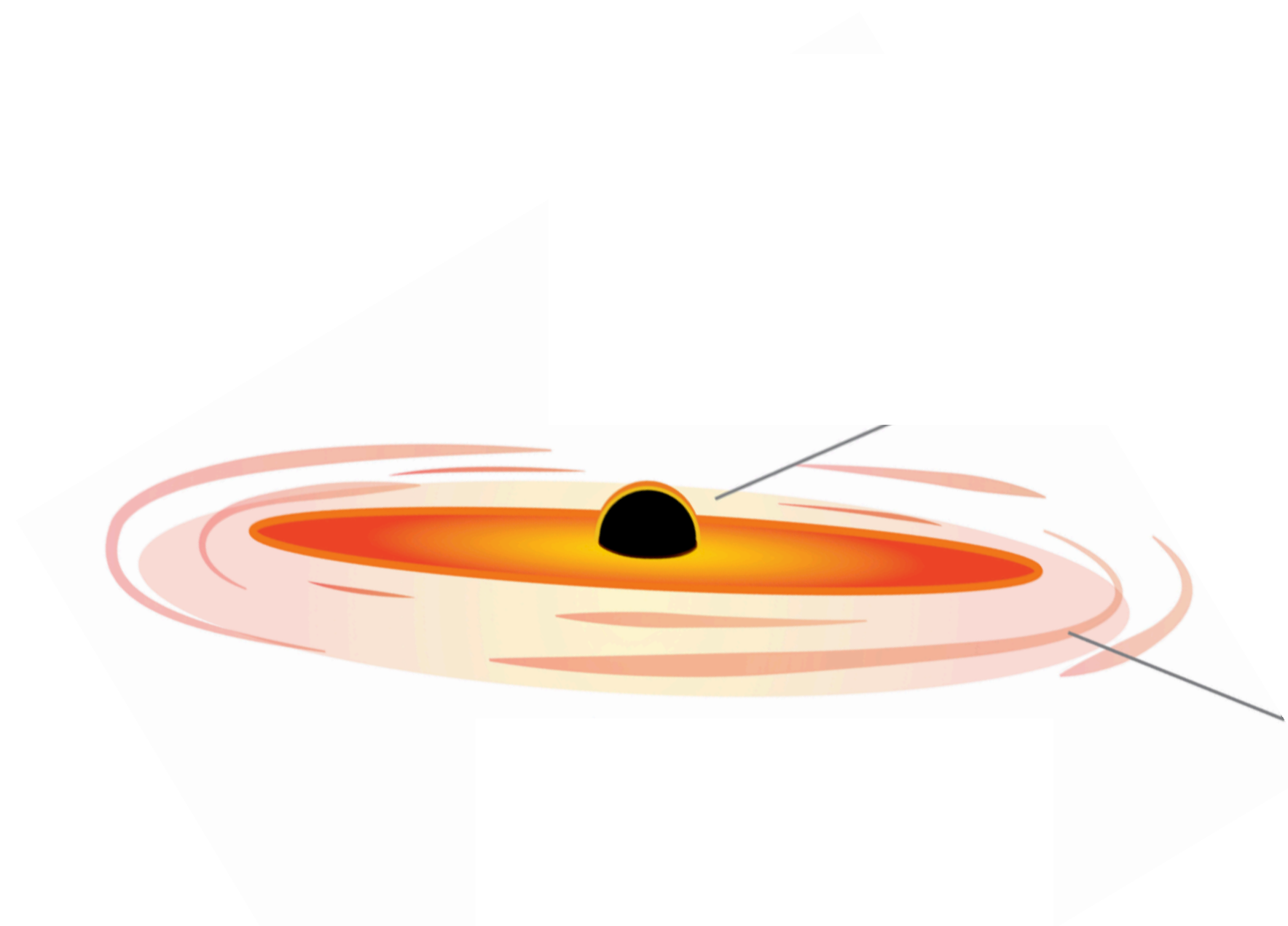
Dark dress



Cold, collisionless
dark matter

Eda et al. 2013, 2014
Gondolo, Silk 1999
Kavanagh et al. 2020
Coogan et al. 2021

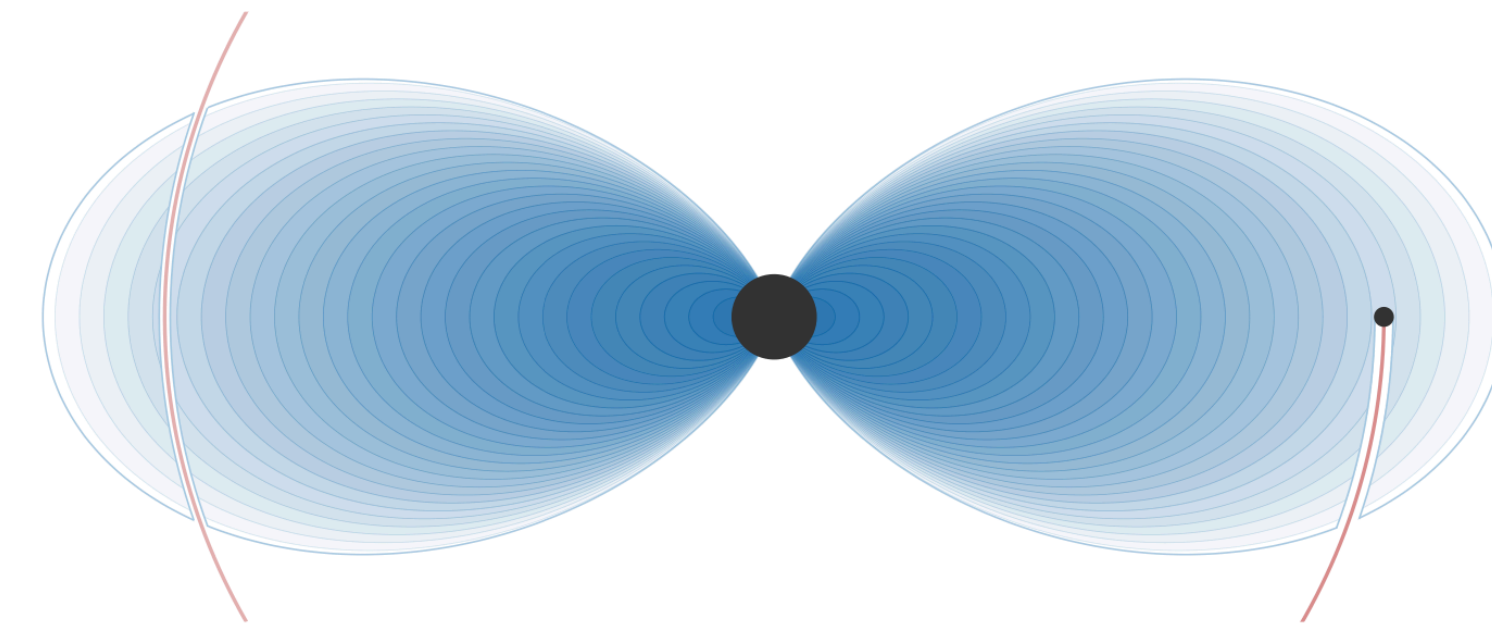
Accretion disc



Baryonic matter

Goldreich & Tremaine 1980
Tanaka 2002
Derdzinski et al. 2020

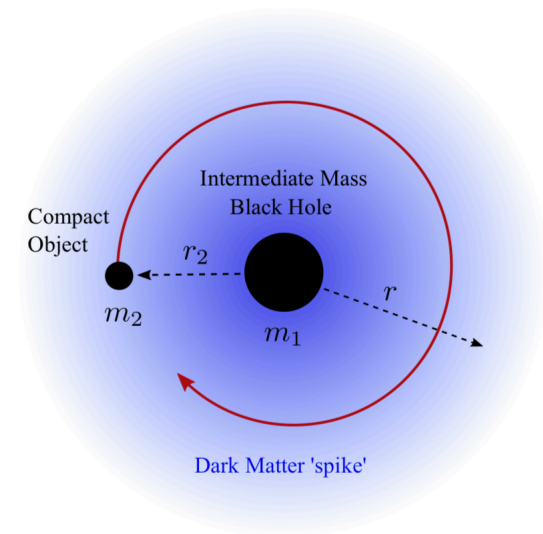
Gravitational atom



Ultra-light bosons

Baumann et al. 2019
Nielsen 2019
Bauman et al. 2021, 2022

Dark dress

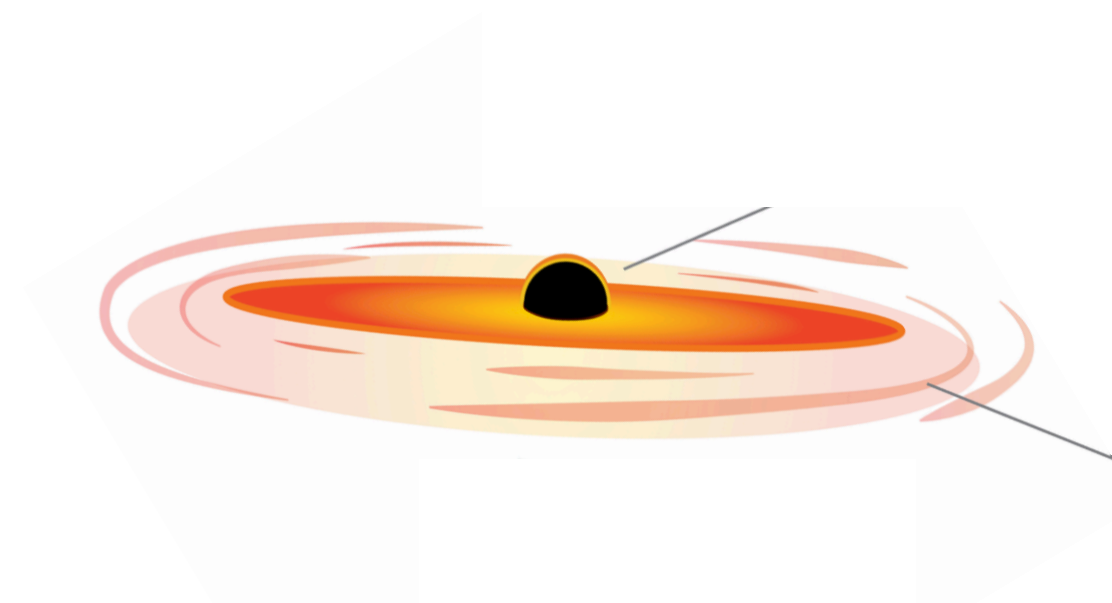


$$\rho(r) = \rho_6 \left(\frac{r_6}{r} \right)^{\gamma_s}$$

Spike density normalisation

Spike power law slope

Accretion disc



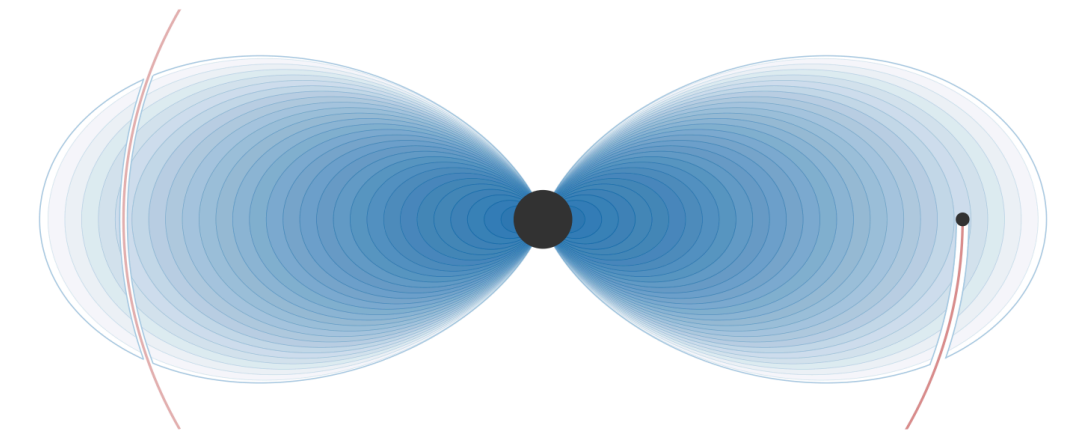
$$\Sigma(r) = \Sigma_0 \left(\frac{r}{r_0} \right)^{-1/2}$$

$$M = r/h$$

Surface density normalisation

Mach number

Gravitational atom



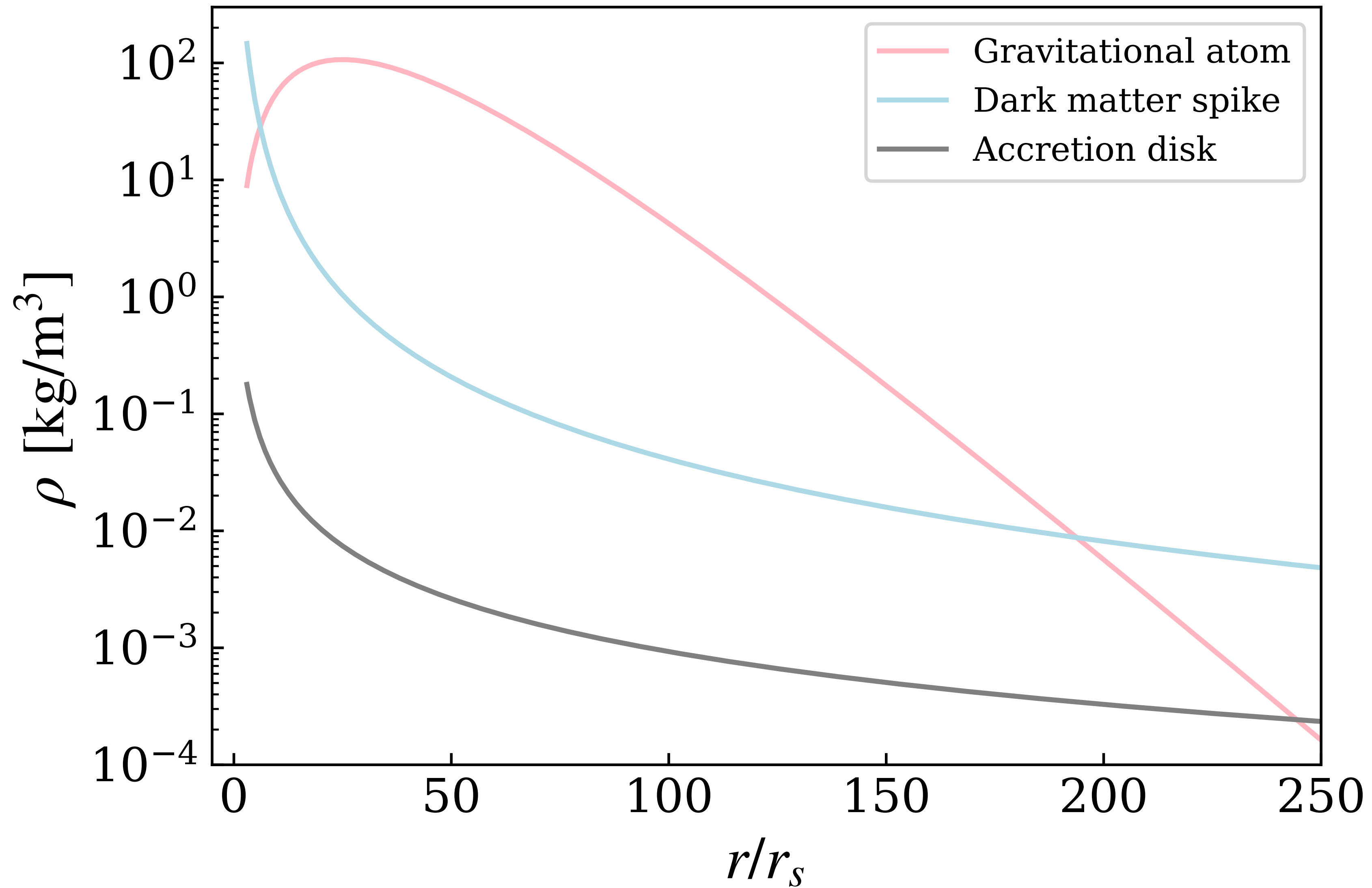
$$\rho(\vec{r}) = M_c |\psi(\vec{r})|^2$$

$$\alpha \equiv Gm_1\mu \ll 1$$

Mass of cloud

Mass of light scalar field
($10^{-10} - 10^{-20}$ eV)

What kind of densities?



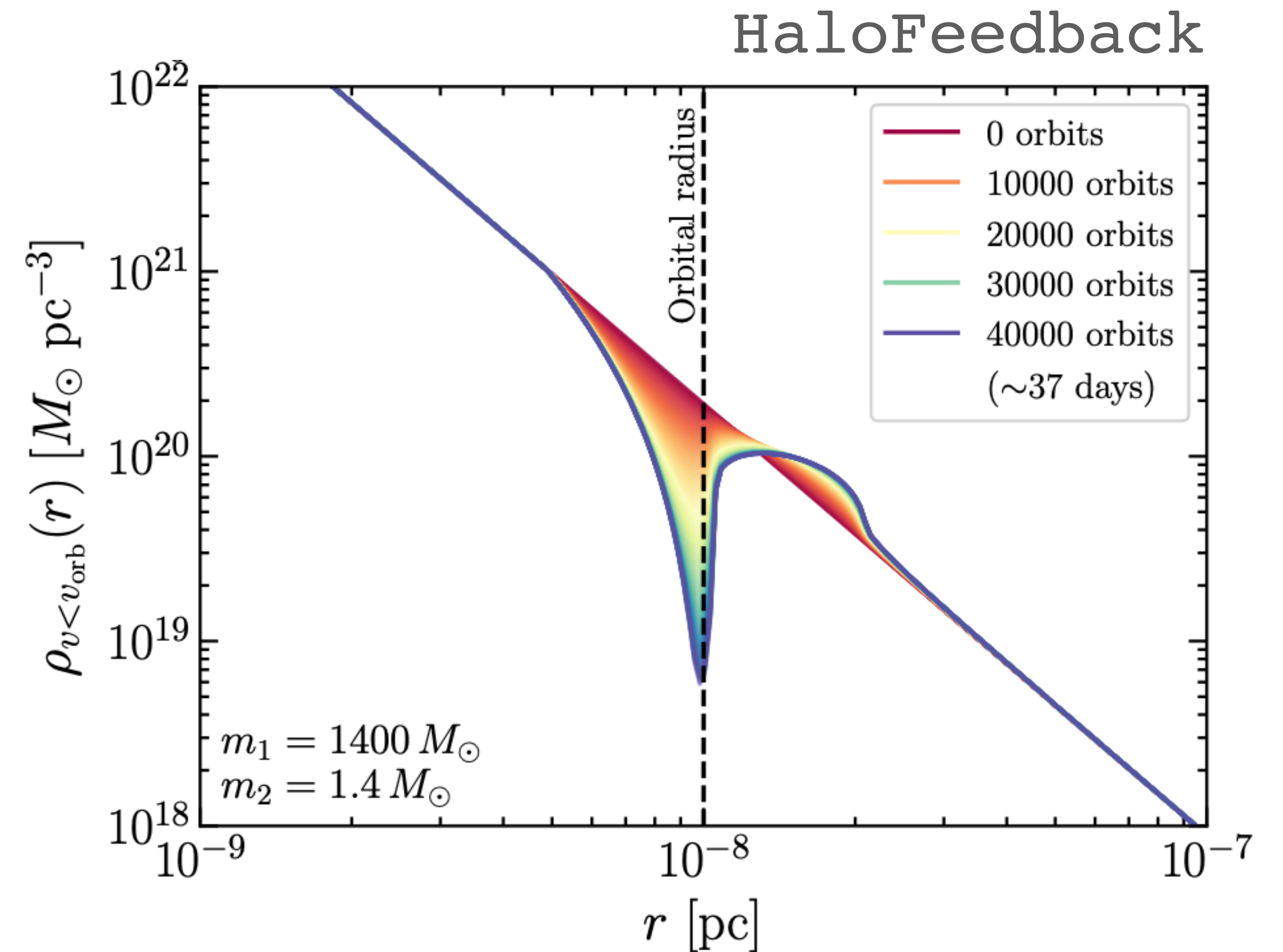
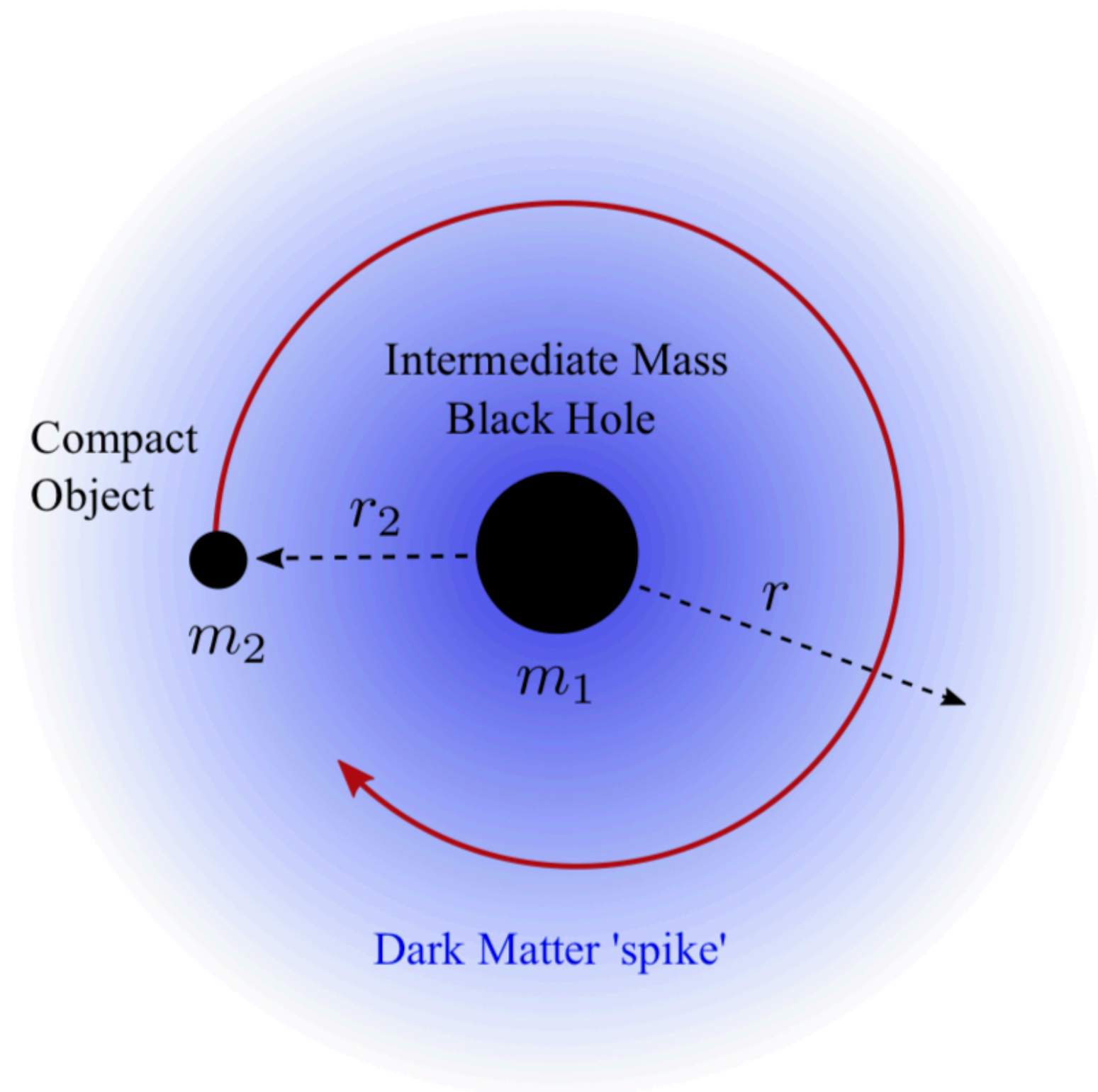
Form of energy losses

$$\dot{r} = \dot{r}_{\text{GW}} + \dot{r}_{\text{env}}$$

Dark dress	Accretion disk	Gravitational atom
$\dot{r}_{\text{env}} = \dot{r}_{\text{DF}}$ <p>Dynamical friction according to Chandrasekhar formula plus feedback on spike with HaloFeedback (Kavanagh et al. 2020)</p>	$\dot{r}_{\text{env}} = \dot{r}_{\text{gas}}$ <p>Gas torques according to Type I migration, analytic prescription including Lindblad and corotation torques</p>	$\dot{r}_{\text{env}} = \dot{r}_{\text{ion}} + \dot{r}_{\text{acc}}$ <p>Ionization (dynamical friction-like) and accretion of scalar field onto companion object</p>

Dynamical friction

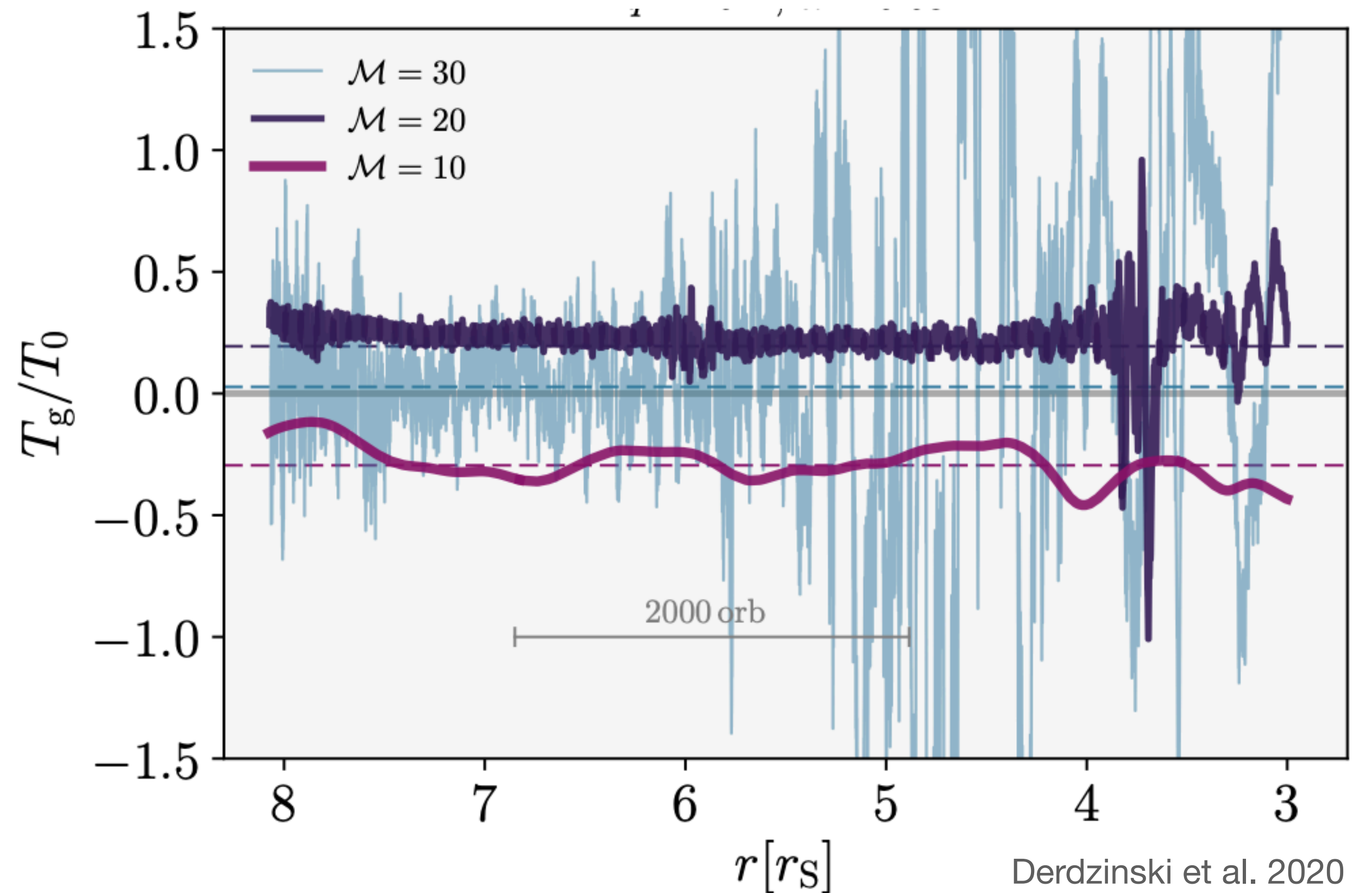
$$\dot{r}_{\text{DF}} = - \frac{8\pi G_N^{1/2} m_2 \log \Lambda r_2^{5/2} \rho_{\text{DM}}(r_2, t) \xi(r_2, t)}{\sqrt{M} m_1}$$



Gas torques

$$\dot{r}_{\text{gas}} = \frac{\dot{L}_{\text{gas}} r^{1/2}}{2\sqrt{G(m_1 + m_2)m_2}}$$

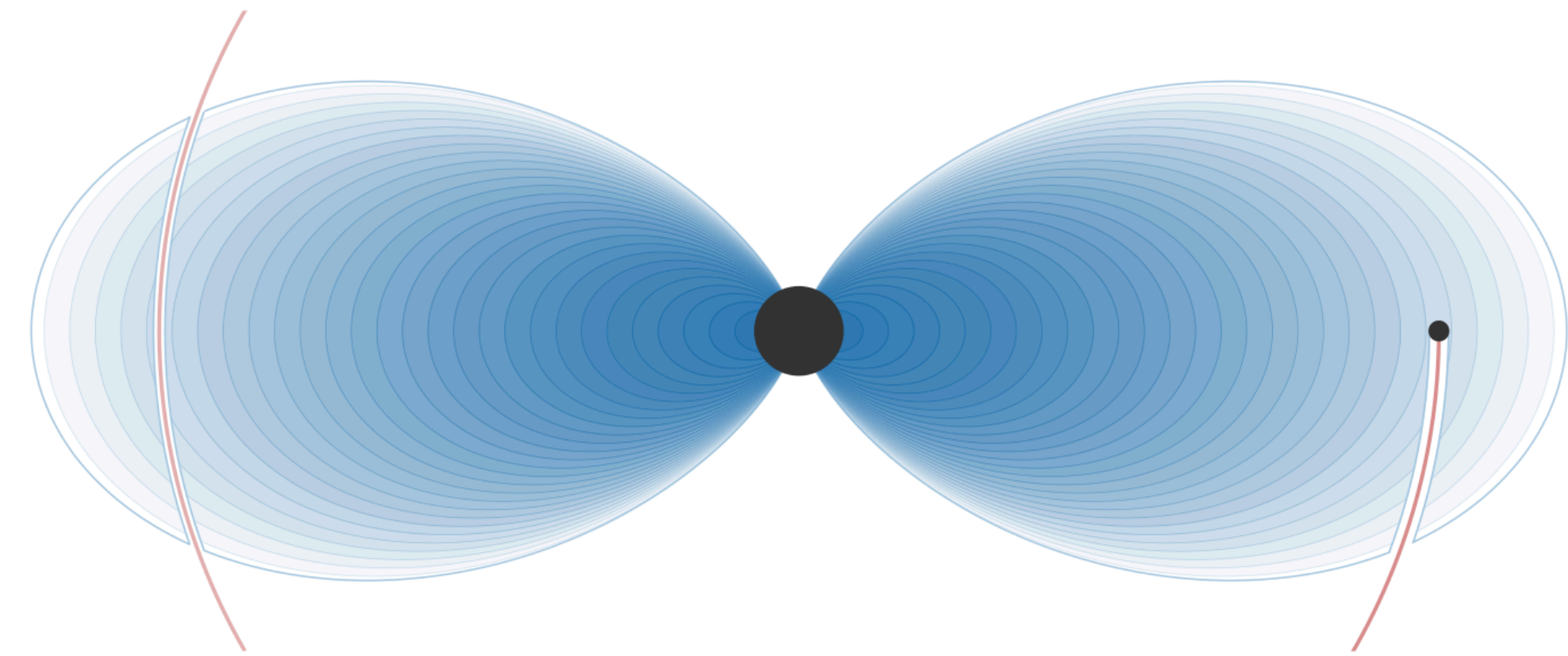
$$\dot{L}_{\text{gas}} = T_{\text{gas}} = \pm \Sigma(r)r^4\Omega^2q^2M^2$$



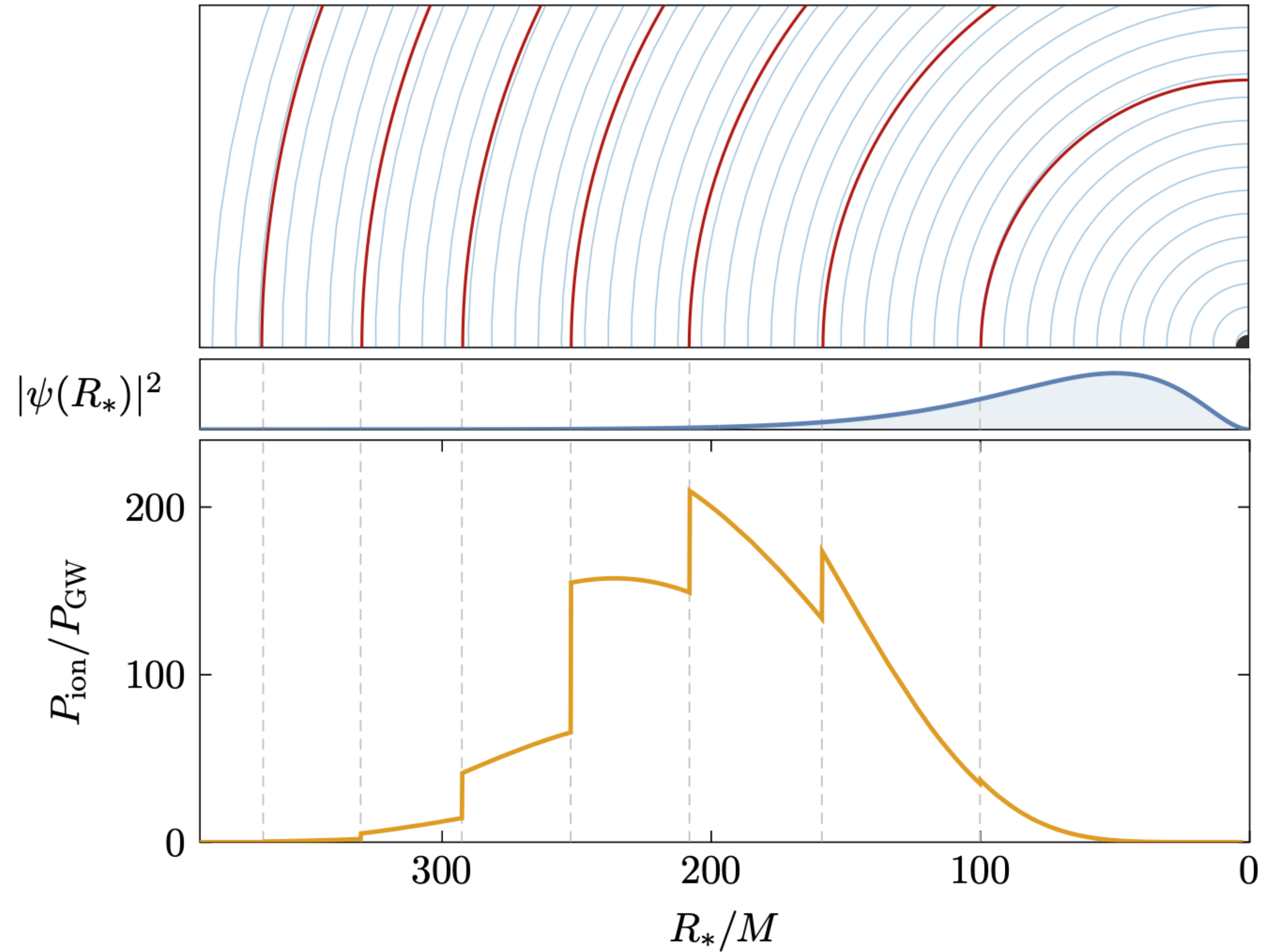
Assume gas in the disc is corotating with the companion object, which is orbiting in the plane of the disc.

Assume Mach number is locally constant, independent of r , i.e. locally isothermal.

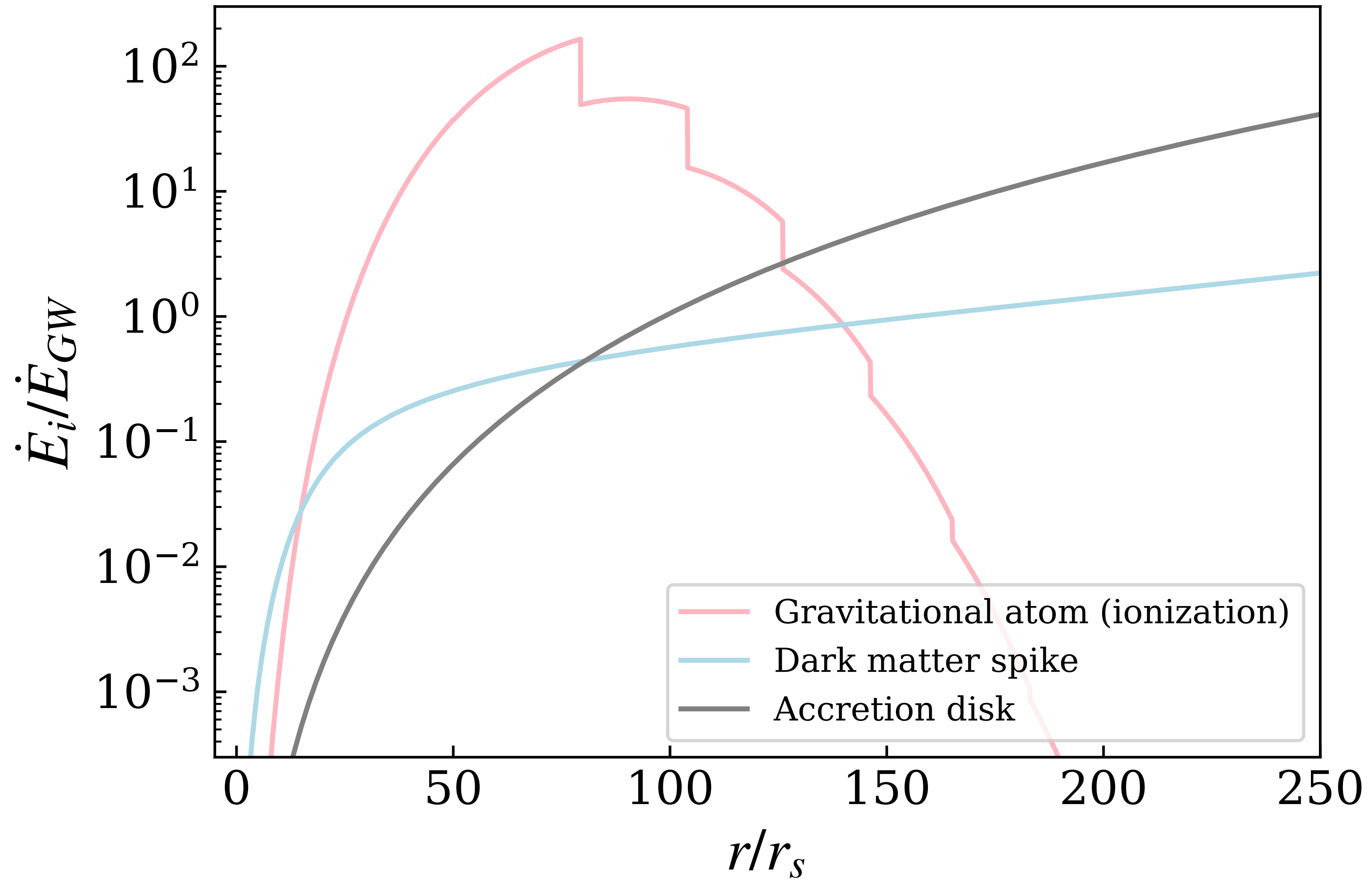
Ionization



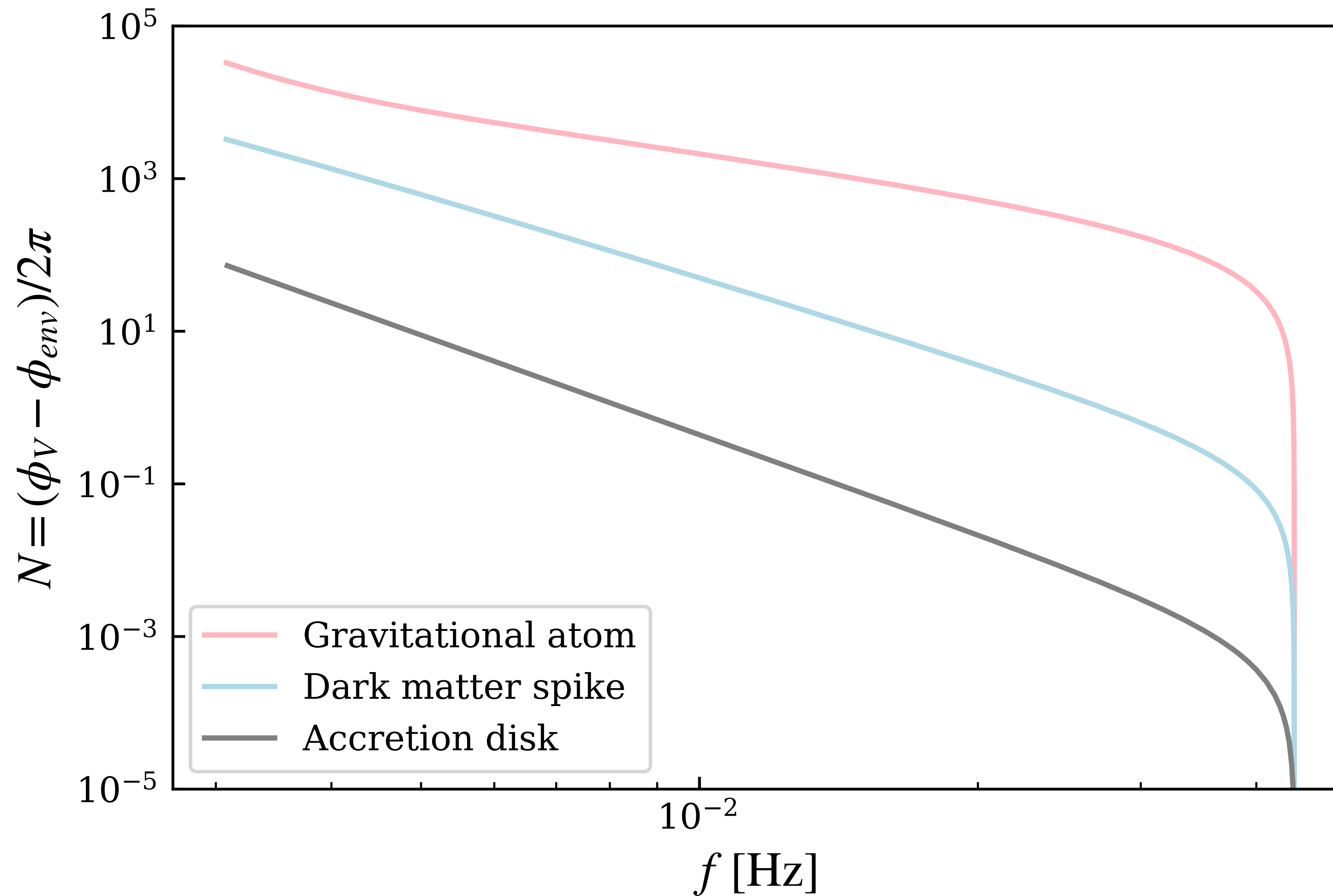
Perturber excites resonances in the cloud and it transitions from bound states to unbound states as the orbital frequency of the perturber hits the frequency of the energy difference between states



Energy losses

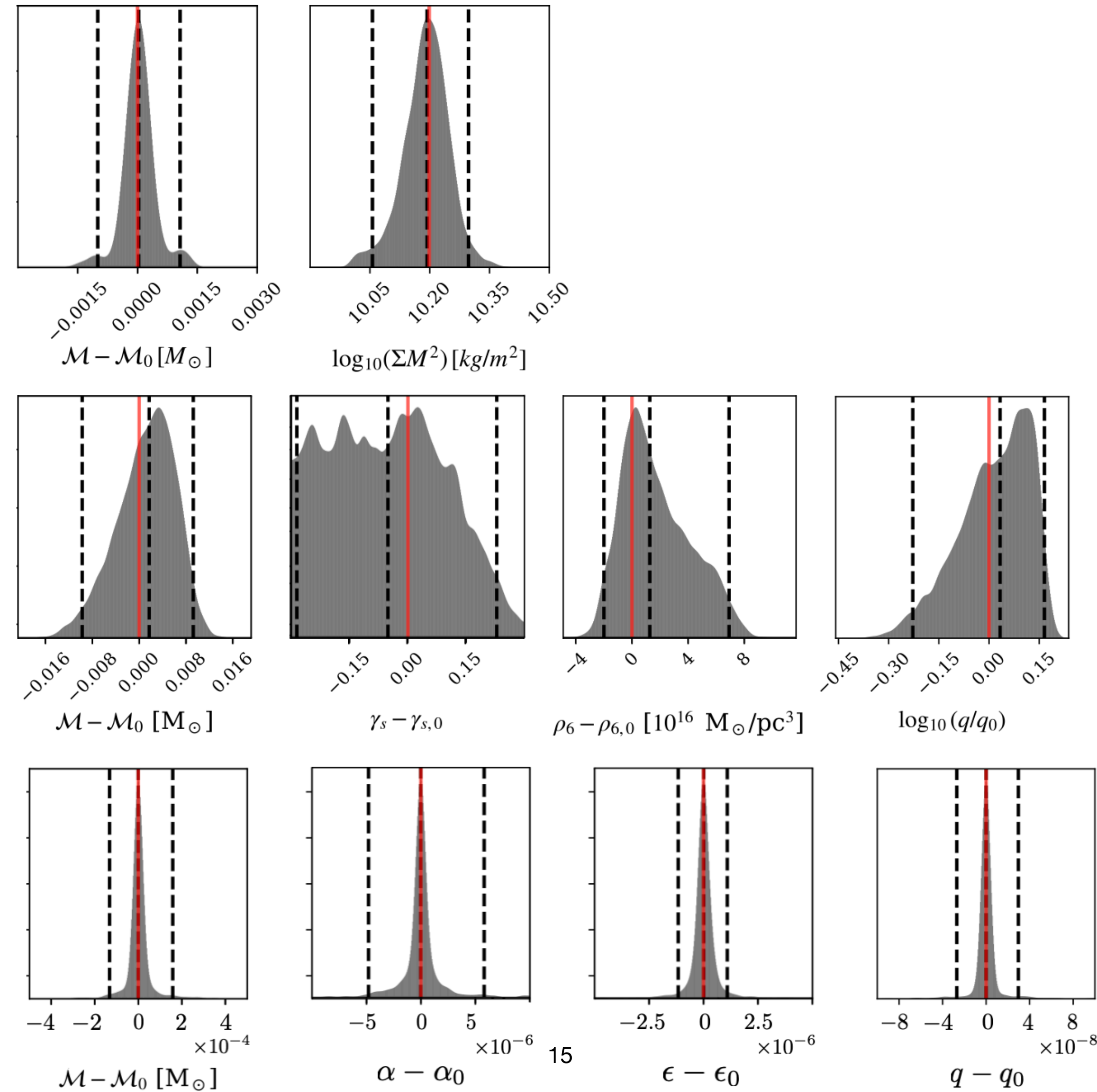


Dephasing

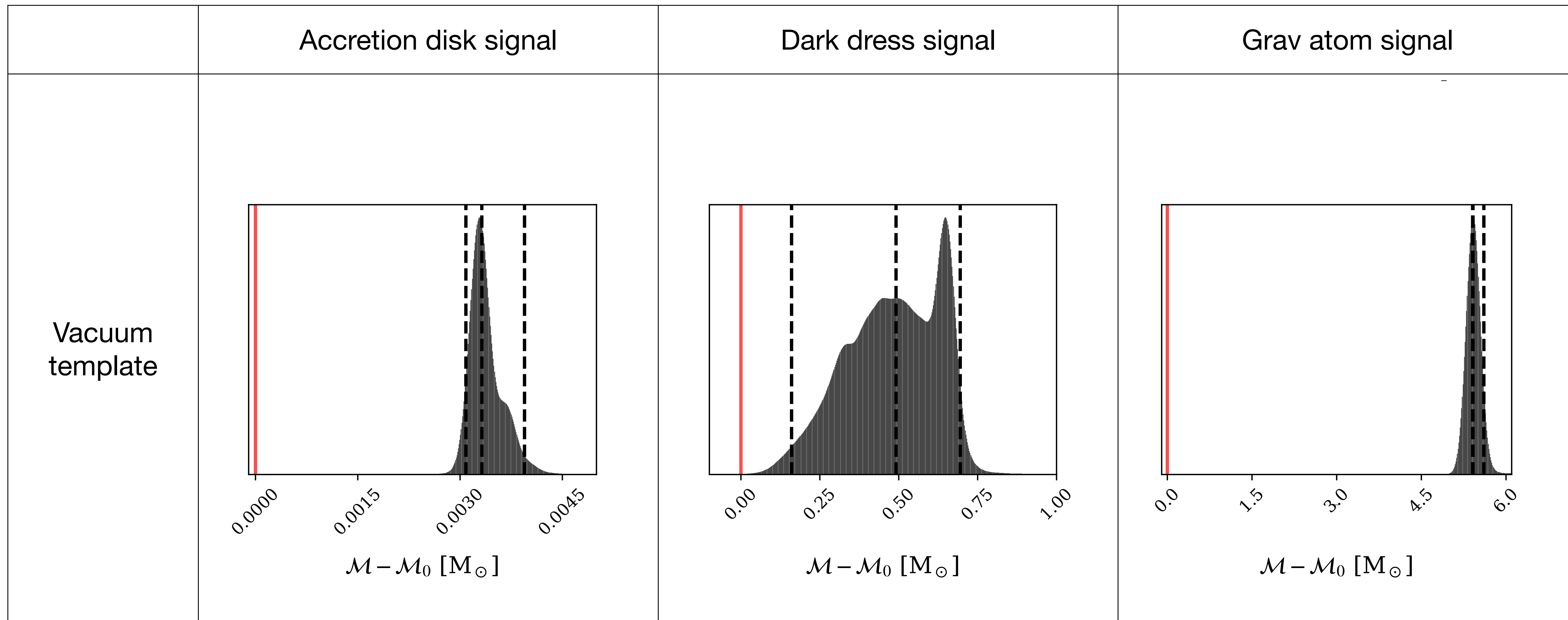


Assuming we've detected a signal, can we measure the parameters?

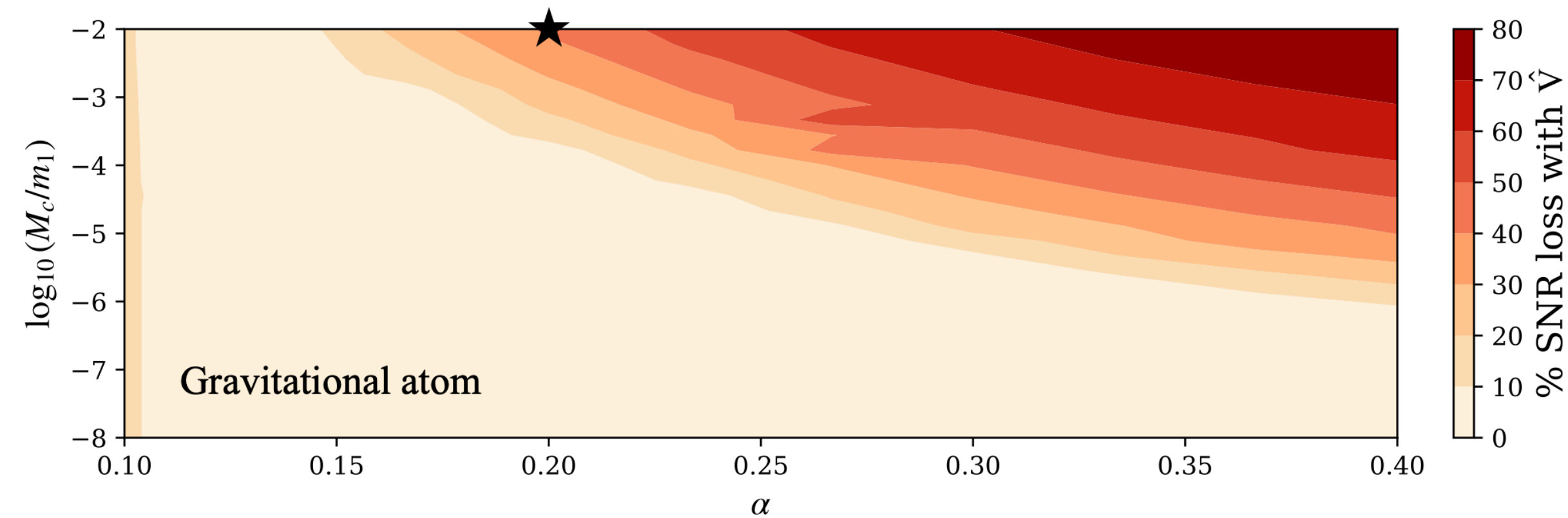
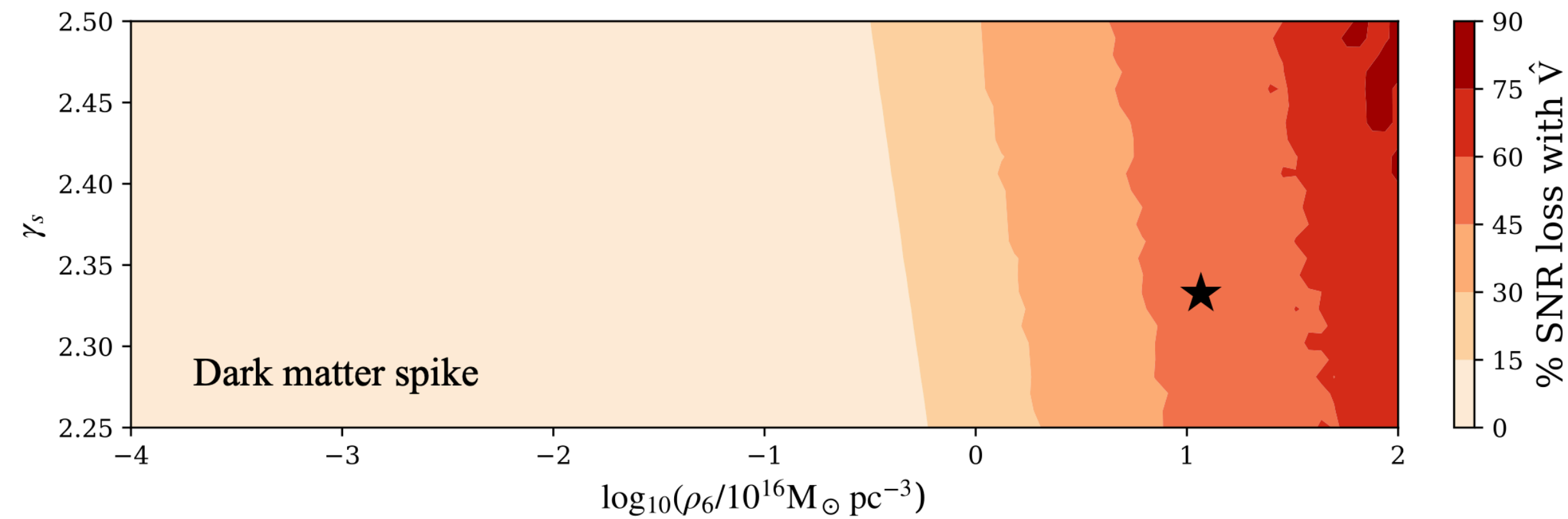
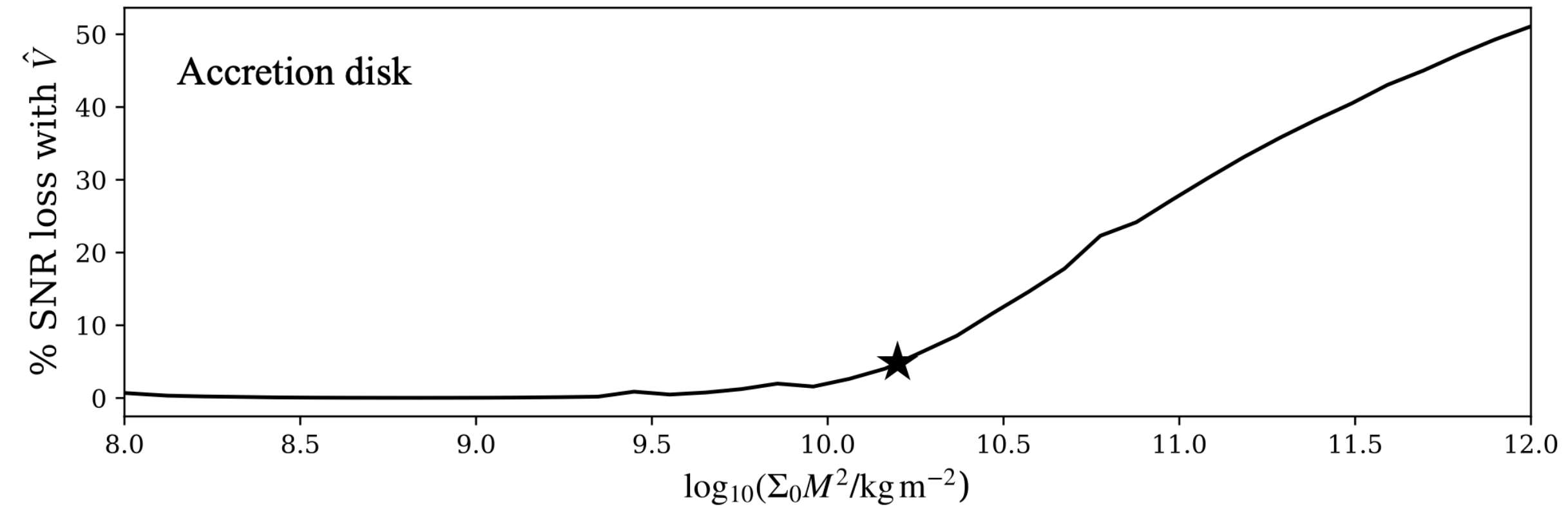
Parameter estimation with correct model



Parameter estimation with vacuum waveform



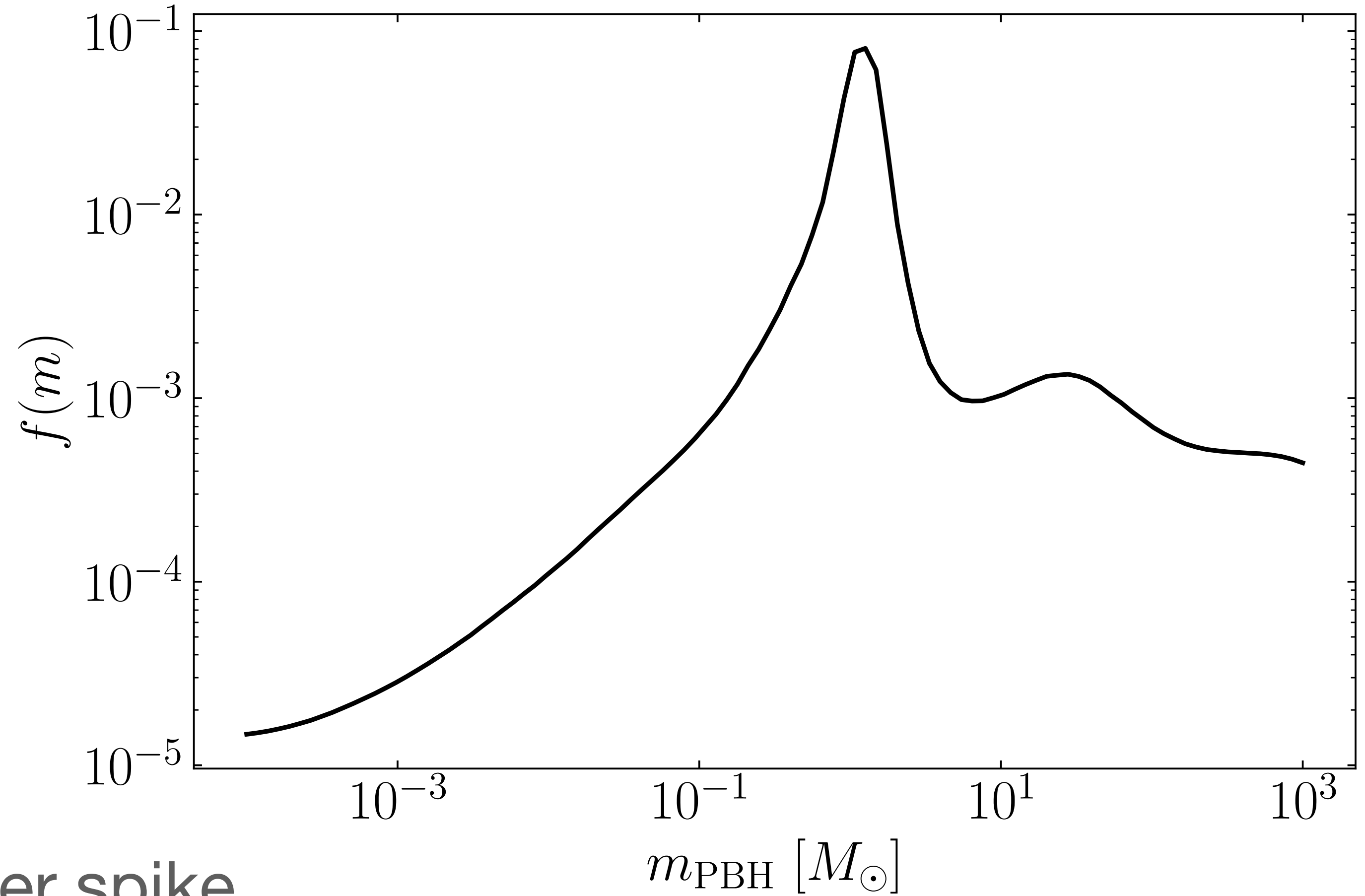
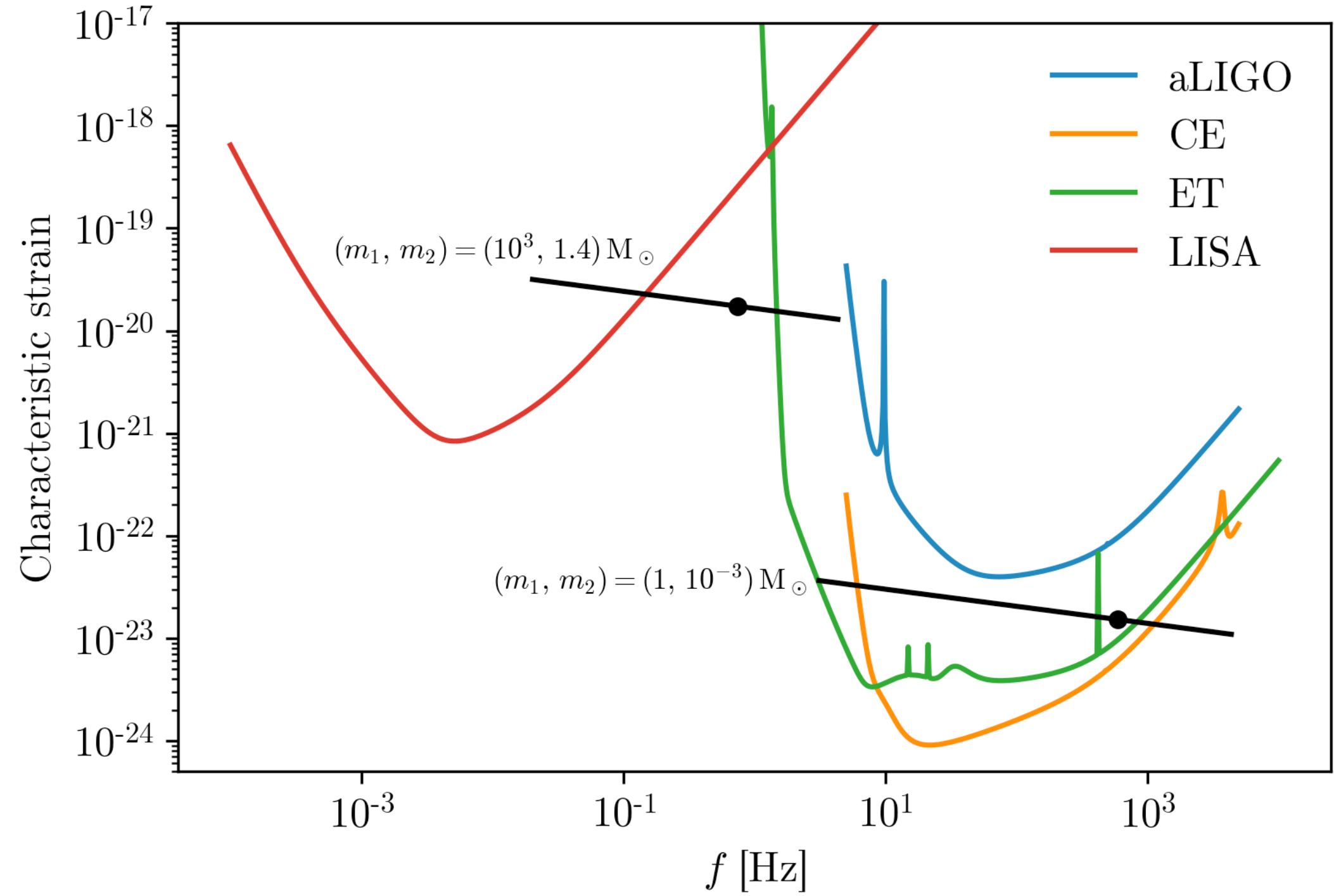
SNR loss: biased PE or miss signal entirely



Bayesian model comparison shows confident preference for correct model over any other environment

$\log_{10} \mathcal{B}$	Dark dress signal	Accretion disk signal	Gravitational atom signal
Vacuum template	34	6	39
Dark dress template	-	3	39
Accretion disk template	17	-	33
Gravitational atom template	24	6	-

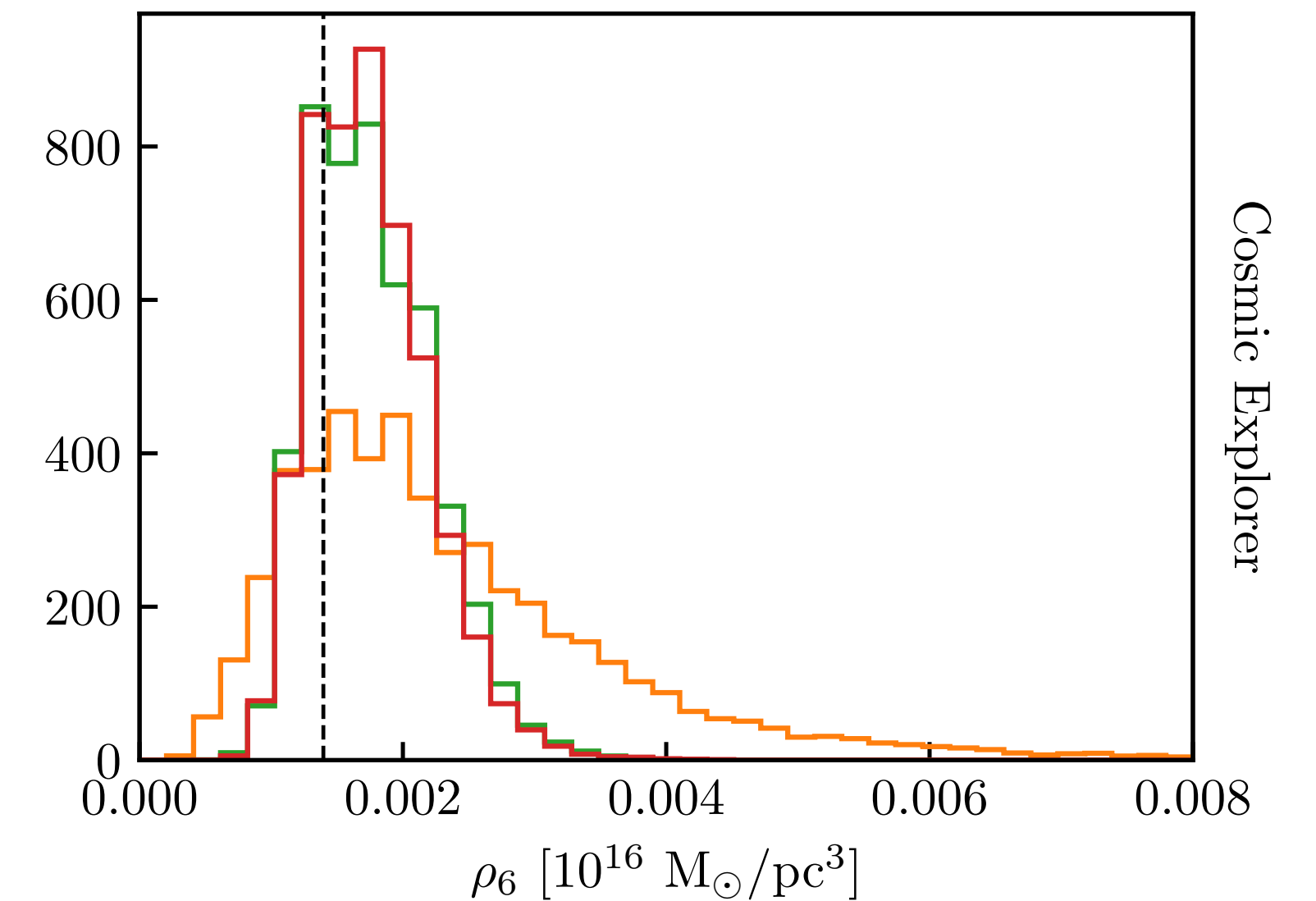
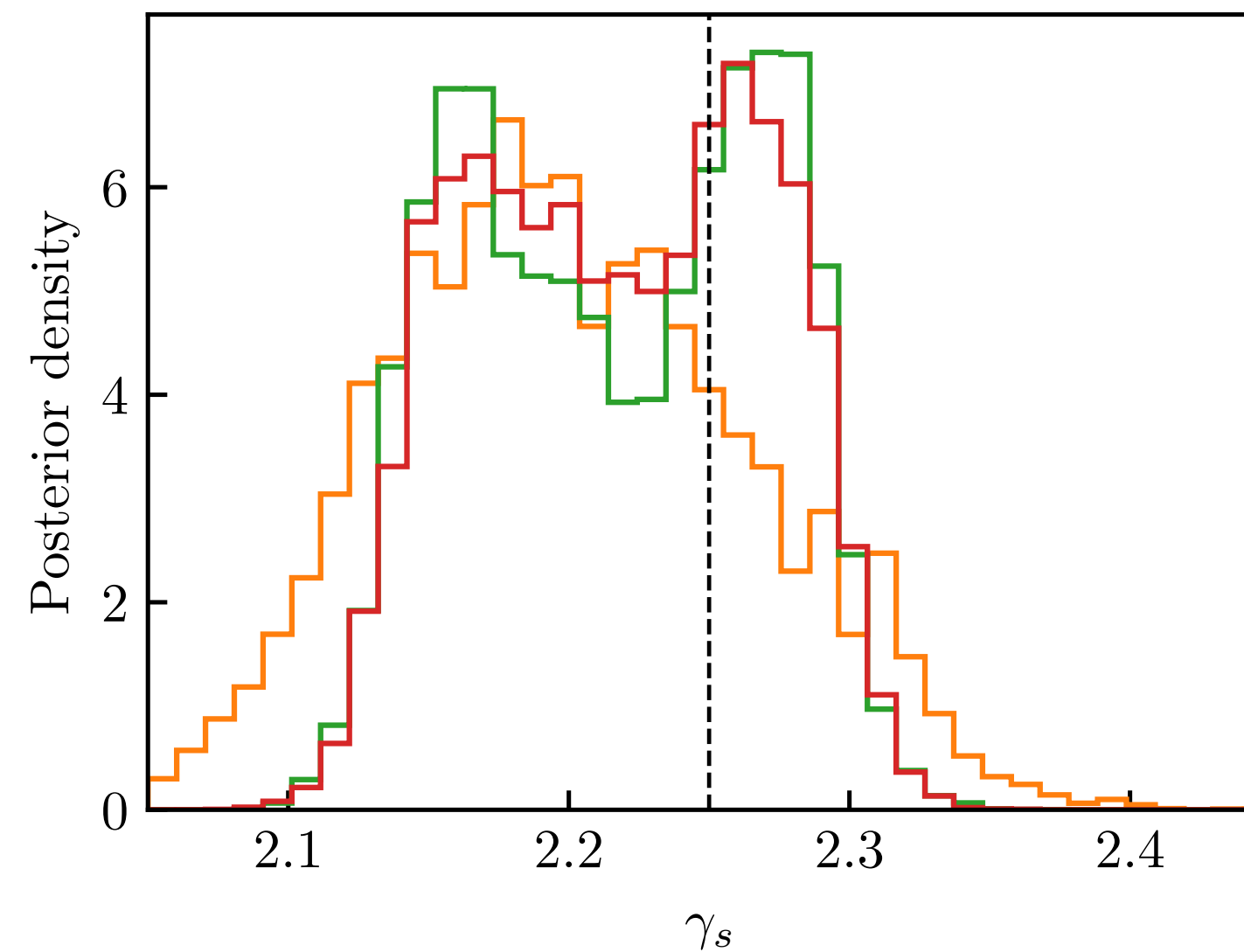
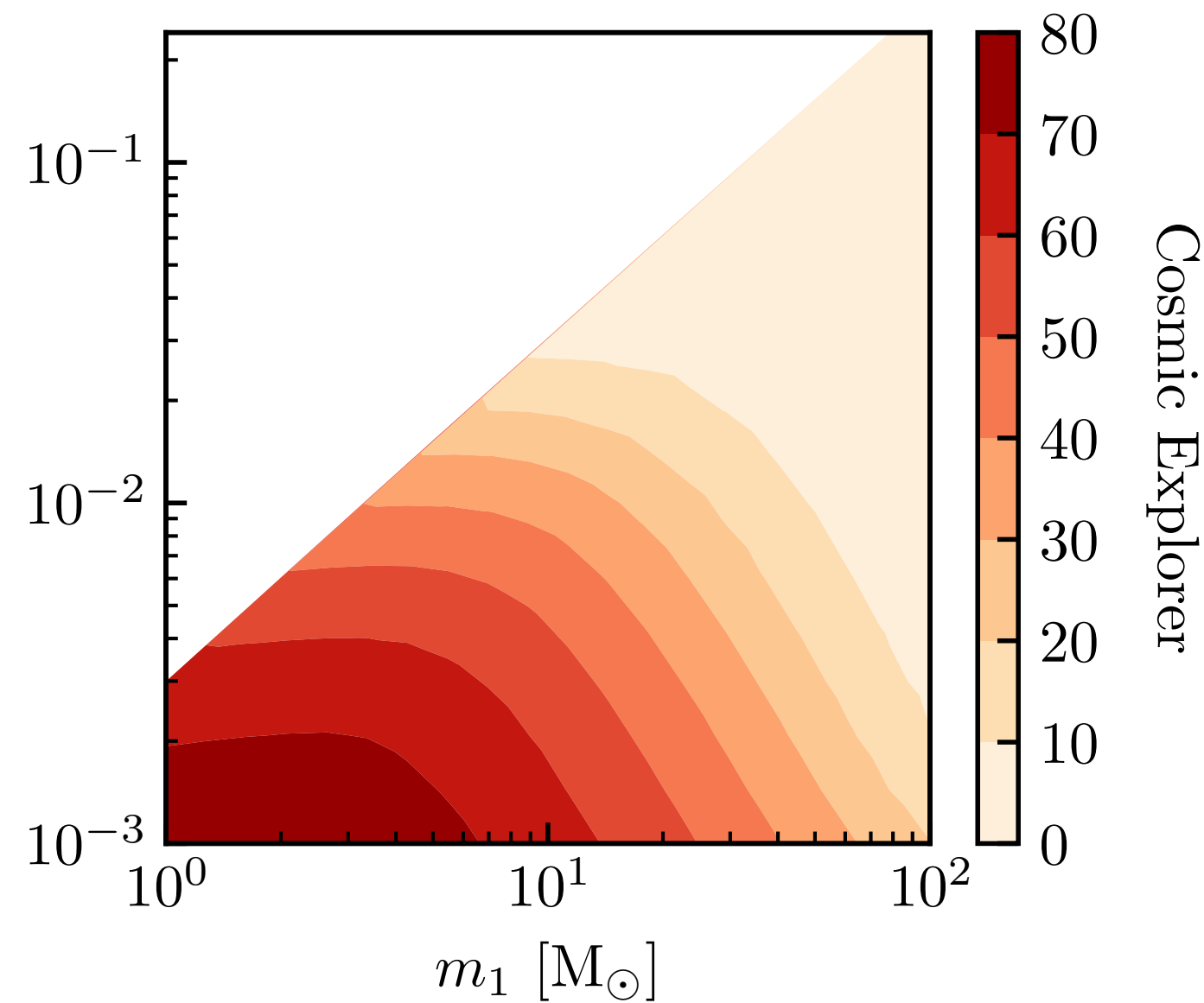
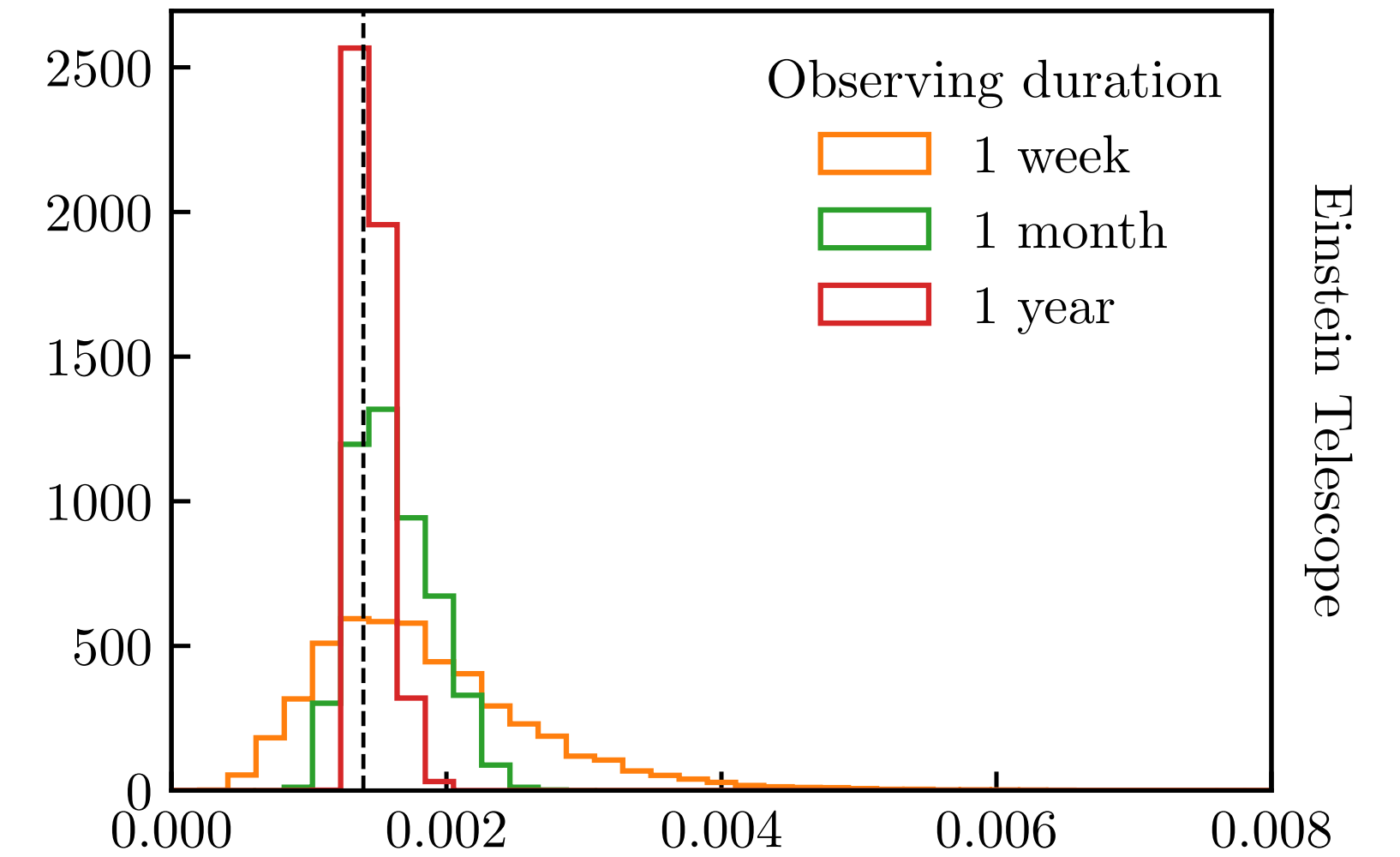
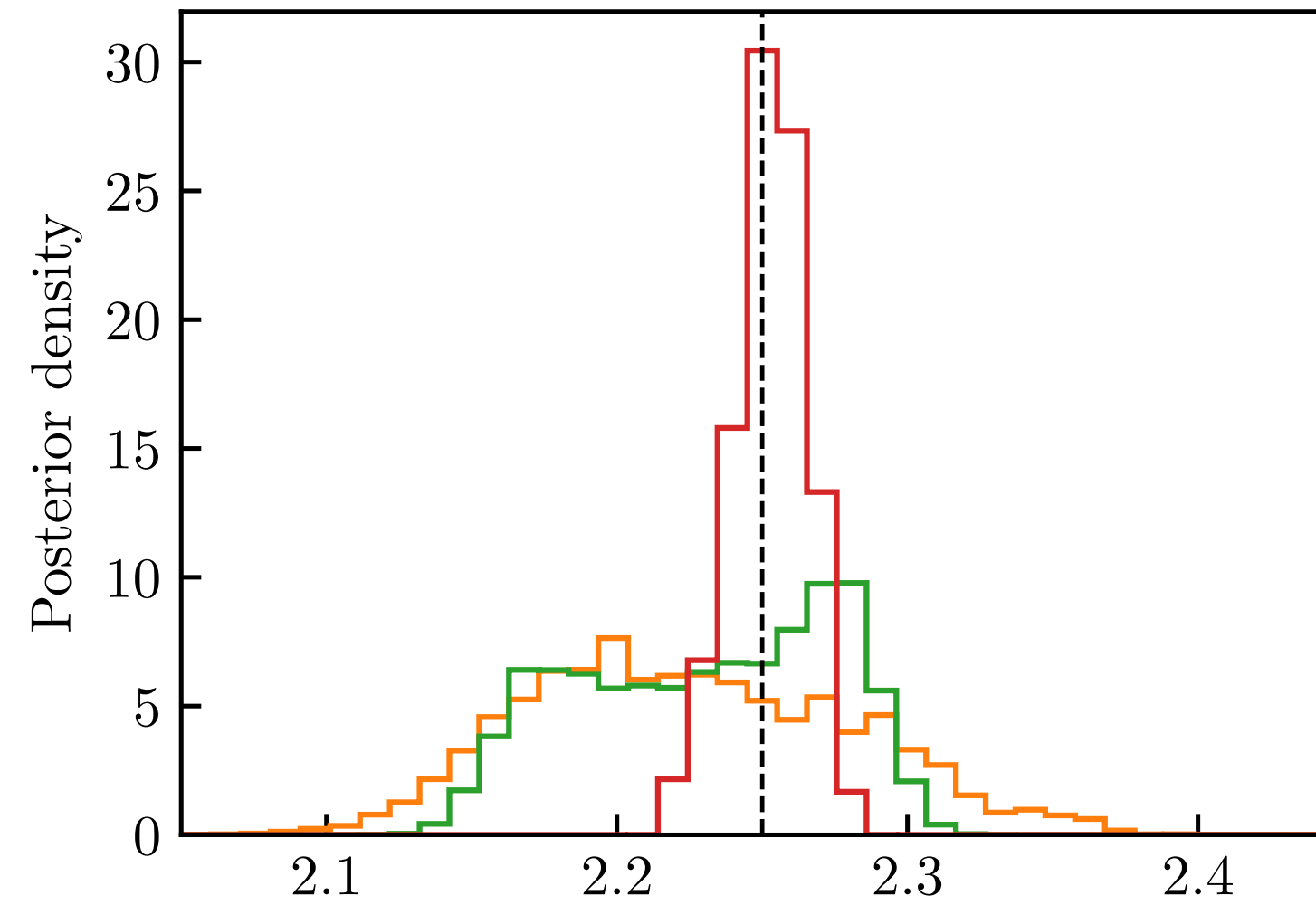
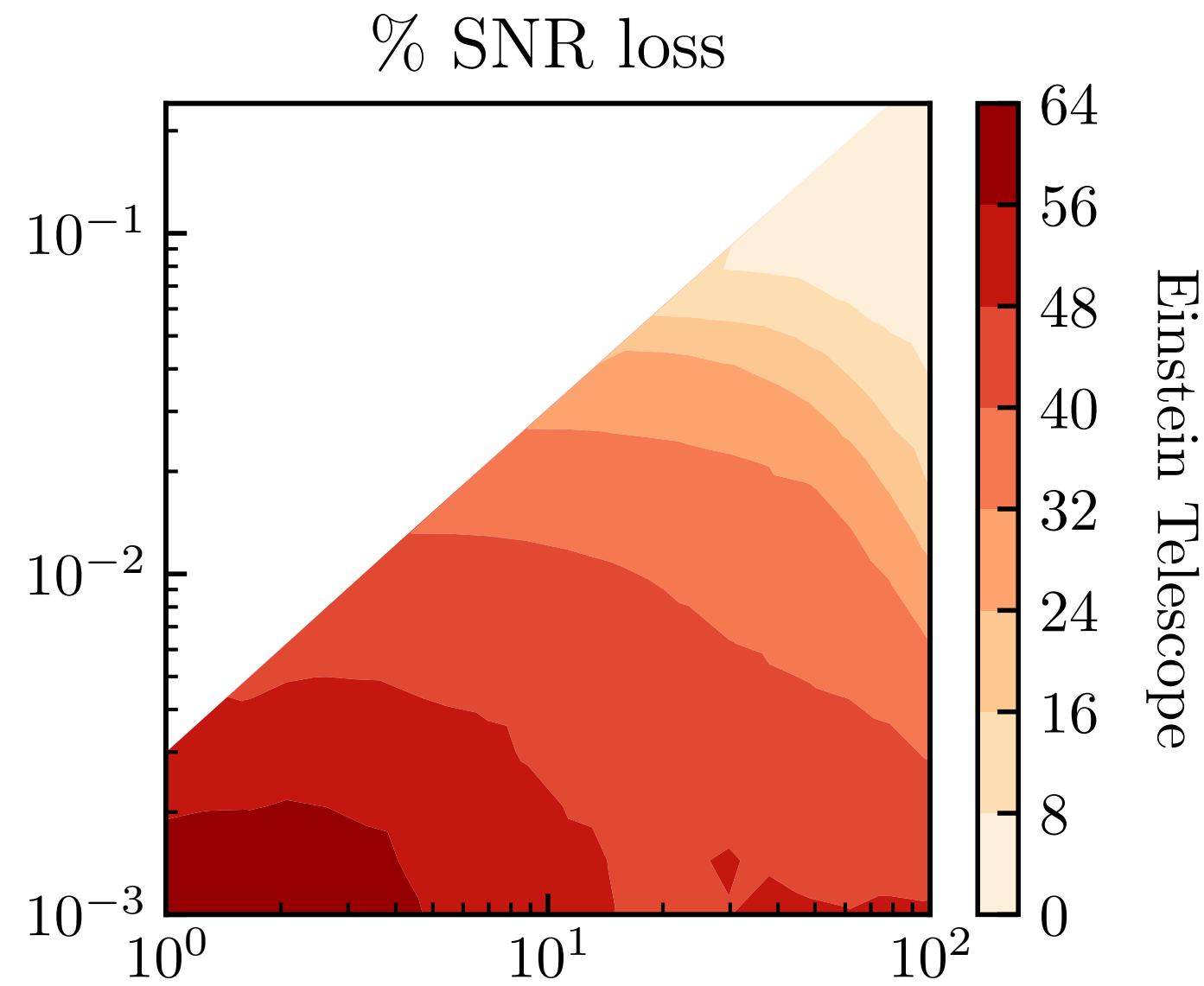
What about future ground-based detectors?



IMRI PBHs must have a dark matter spike

What about future ground-based detectors?

1 week should be enough!



Conclusions

- We can measure the properties of environments around binaries with future GW detectors
- We have an opportunity to learn about the nature of dark matter from IMRI gravitational waveforms
- We can distinguish between environments and avoid confusion with, for example, accretion disks
- Biased parameter reconstruction is possible if the wrong model is used

Future work:

- More accurate waveforms required
- Include for example eccentricity, spins...
- Go to higher dimensions in parameter estimation to check for degeneracies with extrinsic parameters

Thank you for listening!