# Tests of general relativity from the working group perspective 

Alejandro Cárdenas-Avendaño<br>Princeton Gravity Initiative, Princeton University

## Fundamental physics in extreme gravitational regimes

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There is no universal criteria for when non-linear effects become significant enough to qualitatively change the solutions.

What can we learn from LISA that we cannot using other instruments?

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$d f$
$d t$

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$$
\frac{d f}{d t}=\underbrace{\left(\frac{d E}{d f}\right)^{-1}}_{\text {conservative }}
$$

Sectors for tests of general relativity

dissipative
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Propagation effects!
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Energy flux density carried by the GWs:

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\frac{d E}{d A d t}=\frac{c^{3}}{16 \pi G}\left\langle\left(\frac{d h_{+}}{d t}\right)^{2}+\left(\frac{d h_{\mathrm{x}}}{d t}\right)^{2}\right\rangle
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## Tests of general relativity



Tests of general relativity


Theory-Specific:

## Tests of general relativity

e.g., Berti, et al., CQG Topical Review (2015) Yagi \& Stein, CQG Focus Issue (2016)


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No extra fields*
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Diffeomorphism invariance


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e.g., Berti, Cardoso \& Starinets (0905.2975); Cardoso et al., (1901.01265)

Parametrized ringdown (вн Spectroscopy)











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\text { Constant radial profile: } \quad \rho_{\ell m}=\beta_{\ell} \delta_{m, 0} \frac{M}{R^{3}}\left(\frac{a}{R}\right)^{\ell}
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## Black holes beyond general relativity

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## Metric tenso <br> Matter fields <br> e.g., Maselli, Franchini, Gualtieri \& Sotiriou, (2004.11895) <br> $S[\mathbf{g}, \varphi, \Psi]=S_{0}[\mathbf{g}, \varphi]+\alpha S_{c}[\mathbf{g}, \varphi]+S_{\mathrm{m}}[\mathbf{g}, \varphi, \Psi]$ <br> Scalar field

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Metric tensor Matter fields

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Scalar field

$$
S_{0} \propto \int \sqrt{-g}\left(R-\partial_{\mu} \varphi \partial^{\mu} \varphi-\mu_{\mathrm{s}}^{2} \varphi^{2}\right)
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## Black holes beyond general relativity



Shift symmetric Gauss-Bonnet gravity
e.g., Kanti, Mavromatos, Rizos, et al. (9511071)

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Massless $\left(\mu_{\mathrm{s}}=0\right)$ scalar field

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\mathscr{G}=R^{\mu \nu \rho \sigma} R_{\mu \nu \rho \sigma}-4 R^{\mu \nu} R_{\mu \nu}+R^{2}
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$S_{c}[\mathbf{g}, \varphi] \propto \alpha_{\mathrm{SGB}} \int \sqrt{-\mathrm{g}} \varphi \mathscr{G}$

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Thus, these "large" BHs are effectively Kerr BHs.

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EMRIs with LISA can potentially probe the charges (up to $d \sim 10^{-2}$ ) carried by the secondary!

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The resonance condition: $m \Omega_{r}+n \Omega_{\theta}+l \Omega_{\phi}=0$

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$\Omega_{r} ; \Omega_{\theta} ; \Omega_{\phi}$

## EMRIs: Geodesics and Orbital Resonances

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Carter, Phys. Rev. 174 (1968); Schmidt (0202090)
$\Omega_{r} ; \Omega_{\theta} ; \Omega_{\phi}$


Under a perturbation:
(e.g., accretion disk, third body, non-GR)

$$
\left|m \Omega_{r}+n \Omega_{\theta}\right|>\frac{K(\epsilon)}{(|m|+|n|)^{3}}
$$

e.g., Arnold, Kozlov, \& Neishtadt, Mathematical Aspects of Classical

A possible observational detection of non-Kerr black holes
From the trajectory


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Linearized Einsteinian adiabatic fluxes:

$$
E(t)=E(0)+\left.\frac{d E}{d t}\right|_{0} ^{t} \quad L_{z}(t)=L_{z}(0)+\left.\frac{d L_{z}}{d t}\right|_{0} t
$$

## Can we model this behavior by just changing the fluxes?

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$$
d s^{2}=d s_{\text {Kerr }}^{2} \mathcal{O}\left(a^{2}\right)
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Now, let us add a "kick" to the flux!

## The impact of resonances (very brief description)

e.g., Flanagan \& Hinderer (0805.3337) The (simplified) GR two body problem in the EMRI case in terms of action-angle variables $\mathbf{J}=\left(J_{r}, J_{\theta}, J_{\phi}\right)$

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For a slowly evolving system, each component of the self force can be written as:

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G_{\mathrm{SF}}^{(1)}\left(j_{r}, j_{\theta}, \mathbf{J}\right)=\sum_{m n} G_{\mathrm{SF}, \mathrm{mn}}^{(1)}(\mathbf{J}) e^{i\left(m j_{r}+n j_{\theta}\right)}
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## Why a kick? Resonant effects within GR

Self-force resonances
Flanagan \& Hinderer (1009.4923)

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## Outline: Three short stories

New physics $\rightarrow$ New fields: scalar fields and BHs
A scalar (and potentially other) charge on the secondary will affect the waveform.

Beyond the adiabatic approximation:
Development of full usable waveforms in beyond-GR theories or environments

Multiband or multi-messenger prospects:
Exploit the fundamental role played by different detectors across the
gravitational and electromagnetic spectra

## Multi-band gravitational wave tests of general relativity



## Multi-band gravitational wave tests of general relativity



Yunes \& Pretorius (0909.3328)
Considering the parameterized post-Einsteinian (ppE) formalism:

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# Multi-band gravitational wave tests of general relativity: dCS as an example 

Analytical solutions valid only in the small coupling approximation

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Silva, Holgado, ACA \& Yunes (2004.01253)
I+Love+Q with NICER \& LVC data: $\quad \alpha_{d C S}^{1 / 2} \leq 8.5 \mathrm{~km}$

## Multi-band EM signature in massive black hole

binaries: strong thermal X-ray emission until 1-2 days prior to the merger

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 binaries: strong thermal X-ray emission until 1-2 days prior to the mergerMajor Krauth, Davelaar, Haiman, et. al., (2304.02575)



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## Discussion

We cannot make a list of the truly unexpected. However, there are sources that the community have speculated about that would be quite interesting and revolutionary, if discovered.

Breaking degeneracies with astrophysics, environmental effects, etc., also requires precise modeling. We can also think about synergies for multi band and multi-messenger observations

It is crucial to match the increased level of modeling precision with the expected level of observation precision.

LISA design is changing. Most predictions and tests might be
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Thank you!
cardenas-avendano@princeton.edu

