Tests of general relativity from the working group perspective

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What can we learn from LISA that we cannot using other instruments?









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Sectors for tests of general relativity

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df dt

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conservative

Sectors for tests of general relativity dissipative dE dt



conservative





Sectors for tests of general relativity dissipative Propagation effects!





conservative



Sectors for tests of general relativity dissipative dE dt



conservative

Propagation effects!



Sectors for tests of general relativity dissipative dE **M**t



conservative

Energy flux density carried by the GWs:

 $\frac{dE}{dAdt} = \frac{c^3}{16\pi G} \left\langle \left(\frac{dh_+}{dt}\right)^2 + \left(\frac{dh_\times}{dt}\right)^2 \right\rangle$

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Sectors for tests of general relativity dissipative dE At **Propagation effects!** conservative **Energy flux density** carried by the GWs:





 $\frac{c^3}{G} \sim 10^{36} J \cdot s/m^2$







Theory-Specific:





e.g., Berti, et al., CQG Topical Review (2015) Yagi & Stein, CQG Focus Issue (2016)



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No extra fields*

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e.g., Flanagan (0308111); Pani, Sotiriou & Vernieri (1306.1835); ...







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e.g., Berti, Cardoso & Starinets (0905.2975); Cardoso et al., (1901.01265); McManus, et al., (1906.05155); Baibhav, et al., (2302.03050)...

Parametrized ringdown (**BH Spectroscopy**)




























Modified from Arun et al., Living Reviews in Relativity (2022) 25:4





New physics → New fields: scalar fields and BHs



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Exploit the fundamental role played by different detectors across the gravitational and electromagnetic spectra





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Newtonian star with a **density profile**:

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Bonga & Yang (2106.08342)

 $\rho_{\ell 0}(r=R)Y_{\ell 0}(\theta=0,\phi) = i^{\ell}\frac{2\ell+1}{4\pi}\left(\frac{a}{R}\right)^{\ell}$ Thin shell:

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Decaying radial profile:

$$\rho_{\ell m} = \gamma_{\ell} \delta_{m,0} \frac{M}{R^3} \left(\frac{a}{R}\right)^{\ell} \cos \theta_{\ell m}$$







e.g., Maselli, Franchini, Gualtieri & Sotiriou, (2004.11895)

$S[\mathbf{g}, \varphi, \Psi] = S_0[\mathbf{g}, \varphi] + \alpha S_c[\mathbf{g}, \varphi] + S_m[\mathbf{g}, \varphi, \Psi]$



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Non-minimal couplings





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e.g., Kanti, Mavromatos, Rizos, et al. (9511071)



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Massless ($\mu_s = 0$) scalar field $\mathscr{G} = R^{\mu\nu\rho\sigma}R_{\mu\nu\rho\sigma} - 4R^{\mu\nu}R_{\mu\nu} + R^2$ $S_c[\mathbf{g},\varphi] \propto \alpha_{\mathrm{sGB}} \left[\sqrt{-g} \varphi \mathcal{G} \right]$



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$$R^{2} \qquad \text{Massless } (\mu_{s} = 0) \text{ scalar field} \qquad \mathcal{P} = *R^{\mu\nu\rho\sigma}R_{\mu\nu}$$
$$\int S_{c}[\mathbf{g}, \varphi] \propto \alpha_{dCS} \int \sqrt{-g}\varphi \mathcal{P}$$



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If *M* is the **only relevant scale** for the BH:

Thus, these "large" BHs are effectively Kerr BHs.



 $\mathcal{G} = \nabla_a \mathcal{G}^a$

 $\alpha \ll M^2 \longrightarrow \frac{C}{M} \ll 1$




The field equations become





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$$G^{\alpha\beta} = 8\pi m_p \int \frac{\delta^{(4)} \left[x - y_p(\lambda) \right]}{\sqrt{-g}} \frac{dy_p^{\alpha}}{d\lambda} \frac{dy_p^{\beta}}{d\lambda} d\lambda$$





Difference in the GW phase evolution of EMRIs with and without scalar charge Maselli, Franchini, Gualtieri, et al., (2106.11325)

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EMRIs with LISA can potentially probe the charges (up to $d \sim 10^{-2}$) carried by the secondary!







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$$d_{\rm sGB} = \frac{2\alpha_{\rm sGB}}{m_p^2} + \frac{73}{60} \frac{\alpha_{\rm sGB}^3}{m_p^6}$$

Julie & Berti (1909.05258)





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Outline: Three short stories

New physics \rightarrow New fields: scalar fields and BHs

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Beyond the adiabatic approximation:

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$\Omega_r; \ \Omega_{\theta}; \ \Omega_{\phi}$



$\Omega_r; \ \Omega_{\theta}; \ \Omega_{\phi}$



EMRIs: Geodesics and Orbital Resonances The **resonance** condition: $m\Omega_r + n\Omega_\theta + l\Omega_\phi = 0$

Carter, Phys. Rev. 174 (1968); Schmidt (0202090)

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Under a perturbation:

(e.g., accretion disk, third body, non-GR)

e.g., Arnold, Kozlov, & Neishtadt, Mathematical Aspects of Classical and Celestial Mechanics (1997)

 $\left| m\Omega_r + n\Omega_\theta \right| > \frac{K(\epsilon)}{\left(\left| m \right| + \left| n \right| \right)^3}$







From the trajectory





From the trajectory





Lukes-Gerakopoulos, Apostolatos & Contopoulos (0906.0093)





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$$E(t) = E(0) + \frac{dE}{dt} \bigg|_{0} t \qquad L_{z}(t) = L_{z}(0) + \frac{dL_{z}}{dt} \bigg|_{0} t$$





 $ds^2 = ds_{\text{Kerr}}^2 \mathcal{O}\left(a^2\right)$









IC: a=0.8, p=4.66 e=0.7 m/M=10⁻⁶

$$FF\left(h_{1},h_{2}
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see Pan, Yang, Bernard, et al., (2306.06576) for an introduction of a resonance effective Hamiltonian

e.g., Flanagan & Hinderer (0805.3337)

 $\frac{dJ_i}{dI_i} = 0$ Action-angles $d\lambda$

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 Action-angles

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$$\frac{dj_i}{d\lambda} = \omega_i(\mathbf{J}) + qg_{i\mathsf{SF}}^{(1)}\left(j_r, j_\theta, \mathbf{J}\right) + \mathcal{O}\left(q^2\right) \qquad \qquad \frac{dJ_i}{d\lambda} = qG_{i\mathsf{SF}}^{(1)}\left(j_r, j_\theta, \mathbf{J}\right) + \mathcal{O}\left(q^2\right)$$

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 Action-angles

But, the motion is not geodesic, so gravitational radiation changes that description

$$\frac{dj_i}{d\lambda} = \omega_i(\mathbf{J}) + qg_{i\mathsf{SF}}^{(1)}\left(j_r, j_\theta, \mathbf{J}\right) + \mathcal{O}\left(q^2\right) \qquad \qquad \frac{dJ_i}{d\lambda} = qG_{i\mathsf{SF}}^{(1)}\left(j_r, j_\theta, \mathbf{J}\right) + \mathcal{O}\left(q^2\right)$$

The adiabatic approximation allows to write:

$$\frac{dj_i}{d\lambda} \approx \omega_i(\mathbf{J})$$
$$\frac{dJ_i}{d\lambda} \approx q \left\langle G_{i}^{(1)}\left(j_r, j_{\theta}, \mathbf{J}\right) \right\rangle$$
$$\overset{\text{d}}{\text{iSF}} \left(j_r, j_{\theta}, \mathbf{J}\right) \right\rangle$$

e.g., Flanagan & Hinderer (0805.3337)



$$\frac{dJ_i}{d\lambda} = 0$$
 Action-angles

But, the motion is not geodesic, so gravitational radiation changes that description

$$\frac{dj_i}{d\lambda} = \omega_i(\mathbf{J}) + qg_{i\mathsf{SF}}^{(1)}\left(j_r, j_\theta, \mathbf{J}\right) + \mathcal{O}\left(q^2\right) \qquad \qquad \frac{dJ_i}{d\lambda} = qG_{i\mathsf{SF}}^{(1)}\left(j_r, j_\theta, \mathbf{J}\right) + \mathcal{O}\left(q^2\right)$$

The adiabatic approximation allows to write:

 $\frac{\omega_{J_i}}{d\lambda} \approx \omega_i(\mathbf{J})$

For a slowly evolving system, each component of the self force can be written as:

$$\frac{dJ_i}{d\lambda} \approx q \left\langle G_{i}^{(1)} \left(j_r, j_{\theta}, \mathbf{J} \right) \right\rangle$$

e.g., Flanagan & Hinderer (0805.3337)



$$G_{\mathsf{SF}}^{(1)}\left(j_{r}, j_{\theta}, \mathbf{J}\right) = \sum_{mn} G_{\mathsf{SF},\mathsf{mn}}^{(1)}\left(\mathbf{J}\right) e^{i\left(mj_{r} + nj_{\theta}\right)}$$

$$\frac{dJ_i}{d\lambda} = 0$$
 Action-angles

But, the motion is not geodesic, so gravitational radiation changes that description

$$\frac{dj_i}{d\lambda} = \omega_i(\mathbf{J}) + qg_{i\mathsf{SF}}^{(1)}\left(j_r, j_\theta, \mathbf{J}\right) + \mathcal{O}\left(q^2\right) \qquad \qquad \frac{dJ_i}{d\lambda} = qG_{i\mathsf{SF}}^{(1)}\left(j_r, j_\theta, \mathbf{J}\right) + \mathcal{O}\left(q^2\right)$$

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1.

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For a slowly evolving system, each component of the self force can be written as:

 $G_{\mathsf{SF}}^{(1)}\left(j_r, j_{\theta}, \mathbf{s}_{\theta}\right)$

e.g., Flanagan & Hinderer (0805.3337)



$$\mathbf{J} = \sum_{mn} G_{\text{SF,mn}}^{(1)} (\mathbf{J}) e^{i(mj_r + nj_\theta)} \rightarrow \left\langle G_{\text{SF}}^{(1)} (j_r, j_\theta, \mathbf{J}) \right\rangle = G_{\text{SF,mn}}^{(1)}$$



Self-force resonances

Self-force resonances



Self-force resonances



Self-force resonances



Self-force resonances

Flanagan & Hinderer (1009.4923)



While detection may only result in a small percentage of loss, it plays a crucial role in accurately estimating parameters!

Berry, Cole, Cañizares, et al., (1608.08951)

Self-force resonances



Self-force resonances

Flanagan & Hinderer (1009.4923)



We can use the same formalism to account for perturbations/modifications

Self-force resonances

Flanagan & Hinderer (1009.4923)



We can use the same formalism to account for perturbations/modifications



Self-force resonances

Flanagan & Hinderer (1009.4923)



We can use the same formalism to account for perturbations/modifications

$$\frac{dj_i}{d\lambda} = \omega_i(\mathbf{J}) + qg_{i\mathsf{SF}}^{(1)}\left(j_r, j_\theta, \mathbf{J}\right) + \epsilon g_{i\mathsf{Pert}}^{(1)}\left(j_r, j_\theta, j_\phi, \mathbf{J}\right) + \mathcal{O}$$

 $\frac{dJ_i}{d\lambda} \approx qG_{i}^{(1)}\left(j_r, j_{\theta}, \mathbf{J}\right) + \frac{\epsilon G_{i}^{(1)}\left(j_r, j_{\theta}, j_{\phi}, \mathbf{J}\right)}{i\text{Pert}}\left(j_r, j_{\theta}, j_{\phi}, \mathbf{J}\right)$ $(q^2, \epsilon^2, q\epsilon)$



Self-force resonances

Flanagan & Hinderer (1009.4923)



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Tidal resonances





Self-force resonances

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Outline: Three short stories

New physics \rightarrow New fields: scalar fields and BHs

Beyond the adiabatic approximation:

Multiband or multi-messenger prospects:

A scalar (and potentially other) charge on the secondary will affect the waveform.

Development of full usable waveforms in beyond-GR theories or environments

Exploit the fundamental role played by different detectors across the gravitational and electromagnetic spectra





Multi-band gravitational wave tests of general relativity



Multi-band gravitational wave tests of general relativity



Yunes & Pretorius (0909.3328)

Considering the parameterized post-Einsteinian (ppE) formalism:

$$\Psi\left(f\right) = \Psi_{\mathsf{GR}}\left(f\right)\left(1 + \beta u^{2n-5}\right)$$

Multi-band gravitational wave tests of general relativity



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Carson & Yagi (1905.13155)



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see, also, Perkins, Yunes & Berti (2010.09010) for a comprehensive study and scenarios



Analytical solutions valid only in the small coupling approximation

Analytical solutions valid only in the small coupling approximation

$$\zeta \equiv \frac{16\pi\alpha^2}{M^4} \quad \rightarrow \quad \zeta \ll 1$$

Analytical solutions valid only in the small coupling approximation



Analytical solutions valid only in the small coupling approximation



Carson & Yagi (1905.13155), see also, e.g., Gnocchi, Maselli, Abdelsalhin, et al., (1905.13460)





Analytical solutions valid only in the small coupling approximation



Carson & Yagi (1905.13155), see also, e.g., Gnocchi, Maselli, Abdelsalhin, et al., (1905.13460)



Silva, Holgado, ACA & Yunes (2004.01253)

I+Love+Q with NICER & LVC data:





Multi-band EM signature in massive black hole **binaries:** strong thermal X-ray emission until 1-2 days prior to the merger



Multi-band EM signature in massive black hole **binaries:** strong thermal X-ray emission until 1-2 days prior to the merger

Major Krauth, Davelaar, Haiman, et. al., (2304.02575)





Multi-band EM signature in massive black hole **binaries:** strong thermal X-ray emission until 1-2 days prior to the merger



Major Krauth, Davelaar, Haiman, et. al., (2304.02575)





Discussion

We cannot make a list of the truly unexpected. However, there are sources that the community have speculated about that would be quite interesting and revolutionary, if discovered.

Breaking degeneracies with astrophysics, environmental effects, etc., also requires precise modeling. We can also think about synergies for multi band and multi-messenger observations

It is crucial to match the increased level of modeling precision with the expected level of observation precision.

LISA design is changing. Most predictions and tests might be revisited once we know the configuration LISA will fly with.



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I particularly thank T. Baker, E. Berti, R. Brito, V. Cardoso, P. Pani, C. Sopuerta & T. Sotiriou for their input to write this talk.



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> **Thank you!** cardenas-avendano@princeton.edu

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